

NIST-JANAF Thermochemical Tables. I. Ten Organic Molecules Related to Atmospheric Chemistry

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The structural, spectroscopic, and thermodynamic properties of 10 gas phase organic molecules related to atmospheric chemistry, including three peroxides and four carboxylic acids, are reviewed. The calculation of the thermochemical tables involved the critical evaluation of new spectroscopic data, enthalpy of formation determinations, and the use of recent internal rotation data. Since insufficient information to characterize all 10 molecules exists, estimation schemes were used to provide the missing experimental and theoretical data. © 2001 by the U.S. Secretary of Commerce on behalf of the United States. All rights reserved.

Key words: critical evaluation; enthalpy of formation; group additivity method; heat capacity; internal rotation; molecular structure; normal coordinates; spectroscopic properties; thermodynamic properties; transferable force fields.

Contents

| | | | |
|--|-----|---|-----|
| 1. Introduction..... | 477 | 6.1. Enthalpy of Formation..... | 493 |
| 2. Bromoacetic Acid, $\text{CH}_2\text{Br}-\text{COOH}$ | 478 | 6.2. Heat Capacity and Entropy..... | 493 |
| 2.1. Enthalpy of Formation..... | 478 | 6.3. References..... | 495 |
| 2.2. Heat Capacity and Entropy..... | 479 | 7. Cyanooxomethyl Radical, OCCN | 496 |
| 2.3. References..... | 480 | 7.1. Enthalpy of Formation..... | 496 |
| 3. Chloroacetic Acid, $\text{CH}_2\text{Cl}-\text{COOH}$ | 481 | 7.2. Heat Capacity and Entropy..... | 496 |
| 3.1. Enthalpy of Formation..... | 482 | 7.3. References..... | 497 |
| 3.2. Heat Capacity and Entropy..... | 483 | 8. Oxalic Acid, $\text{HO}-\text{CO}-\text{CO}-\text{OH}$ | 498 |
| 3.3. References..... | 484 | 8.1. Enthalpy of Formation..... | 498 |
| 4. Oxopropanedinitrile, $\text{NC}-\text{CO}-\text{CN}$ | 485 | 8.2. Heat Capacity and Entropy..... | 498 |
| 4.1. Enthalpy of Formation..... | 485 | 8.3. References..... | 499 |
| 4.2. Heat Capacity and Entropy..... | 485 | 9. Methyl Hydroperoxide, $\text{CH}_3-\text{O}-\text{O}-\text{H}$ | 500 |
| 4.3. References..... | 486 | 9.1. Enthalpy of Formation..... | 501 |
| 5. Glycolic Acid, $\text{HO}-\text{CH}_2-\text{COOH}$ | 487 | 9.2. Heat Capacity and Entropy..... | 502 |
| 5.1. Enthalpy of Formation..... | 488 | 9.3. References..... | 502 |
| 5.2. Heat Capacity and Entropy..... | 489 | 10. Dimethyl Peroxide, $\text{CH}_3-\text{O}-\text{O}-\text{CH}_3$ | 504 |
| 5.3. References..... | 490 | 10.1. Enthalpy of Formation..... | 504 |
| 6. Glyoxal, $\text{O}=\text{CH}-\text{CH}=\text{O}$ | 492 | 10.2. Heat Capacity and Entropy..... | 505 |
| | | 10.3. References..... | 505 |
| | | 11. Diacetyl Peroxide, $\text{CH}_3-\text{CO}-\text{O}-\text{O}-\text{CO}-\text{CH}_3$ | 507 |
| | | 11.1. Enthalpy of Formation..... | 508 |
| | | 11.2. Heat Capacity and Entropy..... | 508 |
| | | 11.3. References..... | 509 |
| | | 12. Conclusions..... | 510 |
| | | 13. Acknowledgments..... | 510 |
| | | 14. Extended Bibliographies..... | 510 |

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| | |
|---|-----|
| 14.1. Extended Bibliography for Bromoacetic Acid, C ₂ H ₃ BrO ₂ | 510 |
| 14.2. Extended Bibliography for Chloroacetic Acid, C ₂ H ₃ ClO ₂ | 510 |
| 14.3. Extended Bibliography for Oxopropanedinitrile, C ₃ N ₂ O. | 511 |
| 14.4. Extended Bibliography for Glycolic Acid, C ₂ H ₄ O ₃ | 511 |
| 14.5. Extended Bibliography for Glyoxal, C ₂ H ₂ O ₂ | 511 |
| 14.6. Extended Bibliography for Cyanooxomethyl Radical, NC ₂ O. | 512 |
| 14.7. Extended Bibliography for Oxalic Acid, C ₂ H ₂ O ₄ | 513 |
| 14.8. Extended Bibliography for Methyl Hydroperoxide, CH ₄ O ₂ | 513 |
| 14.9. Extended Bibliography for Dimethyl Peroxide, C ₂ H ₆ O ₂ | 513 |
| 14.10. Extended Bibliography for Diacetyl Peroxide, C ₄ H ₆ O ₄ | 513 |

List of Tables

| | |
|--|-----|
| 1. Ideal gas thermodynamic properties of bromoacetic acid C ₂ H ₃ BrO ₂ (g) at the standard state pressure, $p^\circ = 0.1$ MPa ($T_r = 298.15$ K). | 481 |
| 2. Ideal gas thermodynamic properties of chloroacetic acid C ₂ H ₃ ClO ₂ (g) at the standard state pressure, $p^\circ = 0.1$ MPa ($T_r = 298.15$ K). | 483 |
| 3. Ideal gas thermodynamic properties of oxopropanedinitrile C ₃ N ₂ O(g) at the standard state pressure, $p^\circ = 0.1$ MPa ($T_r = 298.15$ K). | 487 |
| 4. Ideal gas thermodynamic properties of glycolic acid C ₂ H ₄ O ₃ (g) at the standard state pressure, $p^\circ = 0.1$ MPa ($T_r = 298.15$ K). | 491 |
| 5. Ideal gas thermodynamic properties of glyoxal C ₂ H ₂ O ₂ (g) at the standard state pressure, $p^\circ = 0.1$ MPa ($T_r = 298.15$ K). | 494 |
| 6. Ideal gas thermodynamic properties of the cyanooxomethyl radical C ₂ NO(g) at the standard state pressure, $p^\circ = 0.1$ MPa ($T_r = 298.15$ K). | 497 |
| 7. Ideal gas thermodynamic properties of oxalic acid C ₂ H ₂ O ₄ (g) at the standard state pressure, $p^\circ = 0.1$ MPa ($T_r = 298.15$ K). | 500 |
| 8. Ideal gas thermodynamic properties of methyl hydroperoxide CH ₄ O ₂ (g) at the standard state pressure, $p^\circ = 0.1$ MPa ($T_r = 298.15$ K). | 503 |
| 9. Ideal gas thermodynamic properties of dimethyl peroxide C ₂ H ₆ O ₂ (g) at the standard state pressure, $p^\circ = 0.1$ MPa ($T_r = 298.15$ K). | 506 |
| 10. Ideal gas thermodynamic properties of diacetyl peroxide, C ₄ H ₆ O ₄ (g) at the standard state pressure, $p^\circ = 0.1$ MPa ($T_r = 298.15$ K). | 509 |
| 11. Summary of the thermodynamic properties at 298.15 K and the standard state pressure, $p^\circ = 0.1$ MPa. | 510 |

1. Introduction

A large number of biogenic and industrial pollutant species play a direct or indirect role in tropospheric smog chemistry. Modeling of the kinetics of tropospheric chemical reaction processes often requires thermodynamic data. In the following, evaluated thermodynamic data for a few smaller organic species including some peroxides relevant to smog chemistry and atmospheric chemistry in general are presented.

The ideal gas thermodynamic properties of the polyatomic molecules were calculated by standard statistical mechanical methods in which a rigid-rotor harmonic-oscillator model, modified where appropriate for internal rotations, was assumed for each compound. The statistical formulas for thermodynamic functions are discussed in several textbooks, review articles, and reference books.¹⁻⁶ Molecular and spectroscopic constants needed for the calculations were selected from the literature. In a few cases missing data were estimated by analogy to related compounds. For some molecules, the fundamental frequencies were estimated by normal coordinate calculations using force constants transferred from related molecules and the program NCA written by Novikov and Malyshev.⁷

To evaluate the internal rotational contributions to the thermodynamic functions, the internal rotational partition function was formed by the summation of internal rotational energy levels for each rotor. These energy levels were obtained by the diagonalization of the one dimensional Hamiltonian using a potential function of the form

$$V(\varphi) = \frac{1}{2} \sum_n V_n (1 - \cos n\varphi), \quad (1)$$

where φ is the internal rotational angle. The method of generating the internal rotation energy levels has been described by Lewis *et al.*^{8,9} The constant required to generate the internal rotational energy levels for each rotor is the internal rotational constant (F) or reduced moment of inertia of the rotating group (I_r). Where available, the V_n terms and internal rotational constant were taken from spectroscopic data. If the F value was unavailable, it was calculated from the reduced moment of inertia with the relationship

$$F = h/8\pi^2 c I_r. \quad (2)$$

The value of I_r was calculated using molecular structural parameters with a computer program based on a method of calculating the reduced moments of inertia developed by Pitzer and Gwinn.^{10,11}

A molecular model of an equilibrium mixture of *trans* and *cis* isomers was employed for calculating the thermodynamic functions of glyoxal.^{4,12,13} This method uses the enthalpy difference between the two conformers to calculate the equilibrium mole fraction of each species. From these data and thermodynamic functions of two conformers, values of thermodynamic functions were calculated, allowing for the mixing of two conformers.

The sources of uncertainties in the calculated thermodynamic functions arise from uncertainties in the molecular

constants used in the calculations as well as deviations from the rigid-rotor harmonic-oscillator model. In this work the uncertainties in the thermodynamic functions were estimated by the procedure developed by Gurvich *et al.*⁶ This approach predicts the uncertainties in the thermodynamic functions $S^\circ(T)$ and $C_p^\circ(T)$ for simple molecules such as C_3N_2O reasonably well. For molecules with one or more internal rotations, the additional uncertainties due to deviations from the rigid-rotor harmonic-oscillator model are difficult to assess. The largest uncertainty probably arises from the anharmonicity of the asymmetric torsion. This will have little effect at room temperature but may be significant at the higher temperatures. The total estimated uncertainties in the thermodynamic functions $S^\circ(T)$ and $C_p^\circ(T)$ in the range between 298.15 and 2000 K are given in the discussions for each molecule.

Based on the selected values of the molecular constants, the ideal gas thermodynamic functions, heat capacity, $C_p^\circ(T)$, entropy, $S^\circ(T)$, enthalpy [$H^\circ(T) - H^\circ(298.15\text{ K})$], and the Gibbs energy function $\{-[G^\circ(T) - H^\circ(298.15\text{ K})]/T\}$, have been calculated for selected temperatures up to 2000 K at the standard state pressure, $p^\circ = 0.1\text{ MPa}$. (In the tables that follow in a few cases excited electronic states have been factored into the calculations; the energy of an electronic state relative to the ground electronic state is given as ϵ_Γ ; the degeneracy of electronic states are referred to in these tables as the "quantum weight," g_Γ .) The enthalpy of formation values [$\Delta_f H^\circ(298.15\text{ K})$] were selected by analyzing experimental studies which may result in the enthalpies of formation determination. In the absence of experimental data, the $\Delta_f H^\circ(298.15\text{ K})$ values were estimated by approximate methods accepted as standard for organic molecules and radicals.^{14,15}

The calculated values of the enthalpy difference, [$H^\circ(T) - H^\circ(298.15\text{ K})$], and entropy, $S^\circ(T)$, of the ideal gas were

combined with values of the enthalpies and entropies of the elements in their reference states to derive values of enthalpy of formation ($\Delta_f H^\circ$), Gibbs energy of formation ($\Delta_f G^\circ$), and the logarithm of the equilibrium constant of formation ($\log K_f^\circ$) of the substances as a function of temperature over the range of 0–2000 K.

Values used here of [$H^\circ(T) - H^\circ(298.15\text{ K})$] and $S^\circ(T)$ for the elements in their reference states [$H_2(g)$, C(cr, graphite), $O_2(g)$, $N_2(g)$, $F_2(g)$, $Cl_2(g)$, and $Br_2(cr, liq, T < 332.503\text{ K})$ and $Br_2(g, T > 332.503\text{ K})$] are those given in the JANAF Thermochemical Tables.⁵

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2. Bromoacetic Acid, CH₂Br—COOHBromoacetic acid (C₂H₃BrO₂)

Ideal gas

 $M_r = 138.9485$ $S^\circ(298.15\text{ K}) = 337.0 \pm 5.0\text{ J K}^{-1}\text{ mol}^{-1}$ $\Delta_f H^\circ(0\text{ K}) = -364.6 \pm 3.1\text{ kJ mol}^{-1}$
 $\Delta_f H^\circ(298.15\text{ K}) = -383.5 \pm 3.1\text{ kJ mol}^{-1}$

Molecular constants

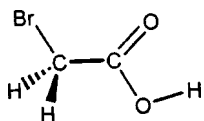
Point group: C_s Symmetry number: $\sigma = 1$ Ground electronic state: \tilde{X}^1A Energy: $\epsilon_X = 0\text{ cm}^{-1}$ Quantum weight: $g_X = 1$

| Vibrational frequencies, ν_i , and degeneracies g_i | | | | | | | |
|---|----------------------------|-------|----------|----------------------------|------------|------------------|---|
| Symmetry | ν_i , cm^{-1} | g_i | Symmetry | ν_i , cm^{-1} | g_i | | |
| A' | ν_1 | 3566 | 1 | A'' | ν_{10} | 589 | |
| | ν_2 | 3037 | 1 | | ν_{11} | 384 | 1 |
| | ν_3 | 1808 | 1 | | ν_{12} | 180 | 1 |
| | ν_4 | 1449 | 1 | | ν_{13} | 3076 | 1 |
| | ν_5 | 1325 | 1 | | ν_{14} | 1243 | 1 |
| | ν_6 | 1208 | 1 | | ν_{15} | 806 | 1 |
| | ν_7 | 1047 | 1 | | ν_{16} | 611 | 1 |
| | ν_8 | 908 | 1 | | ν_{17} | 489 | 1 |
| | ν_9 | 747 | 1 | | ν_{18} | ... ^a | 1 |

^aInstead of torsional mode $\nu_{18} = 47\text{ cm}^{-1}$, the contributions due to the internal rotation about C—C bond were calculated from the potential: $V(\varphi) = \frac{1}{2}V_3(1 - \cos 3\varphi)$, where φ is the torsional angle, $V_3 = 450\text{ cm}^{-1}$.

CH₂Br top: Reduced moment of inertia, $I_r = 2.8300 \times 10^{-39}\text{ g cm}^2$, Symmetry number, $\sigma_m = 1$.

Geometry



$r(\text{C—C}) = 1.51 \pm 0.01\text{ \AA}$
 $r(\text{C=O}) = 1.21 \pm 0.01\text{ \AA}$
 $r(\text{C—O}) = 1.36 \pm 0.01\text{ \AA}$
 $r(\text{C—Br}) = 1.95 \pm 0.02\text{ \AA}$

$r(\text{C—H}) = 1.09 \pm 0.02\text{ \AA}$
 $r(\text{O—H}) = 0.97 \pm 0.02\text{ \AA}$
 $\angle \text{C—C=O} = 126 \pm 1^\circ$
 $\angle \text{C—C—O} = 111 \pm 1.5^\circ$
 $\angle \text{C—C—Br} = 112.5 \pm 1.5^\circ$
 $\angle \text{C—C—H} = 112 \pm 2^\circ$
 $\angle \text{H—C—H} = 109.5 \pm 1^\circ$
 $\angle \text{C—O—H} = 106 \pm 1^\circ$
 $\varphi(\text{O=C—C—Br}) = 0.0^\circ$
 $\varphi(\text{O=C—O—H}) = 0.0^\circ$

Rotational constants in cm^{-1} :

CH₂⁷⁹BrCOOH $A_0 = 0.348\,921$ $B_0 = 0.050\,528$ $C_0 = 0.044\,536$
 CH₂⁸¹BrCOOH $A_0 = 0.348\,916$ $B_0 = 0.050\,066$ $C_0 = 0.044\,176$

Product of moments of inertia: $I_A I_B I_C = 28\,178 \times 10^{-117}\text{ g}^3\text{ cm}^6$.

2.1. Enthalpy of Formation

The recommended value of enthalpy of formation of gaseous bromoacetic acid, $-(383.5 \pm 3.1)\text{ kJ mol}^{-1}$, was obtained by Lagoa *et al.*¹ from experimental measurements. This value is the sum of the enthalpy of formation of bromoacetic acid in the crystalline state, $\Delta_f H^\circ(\text{C}_2\text{H}_3\text{BrO}_2$,

cr I, 298.15 K) = $-(466.98 \pm 1.08)\text{ kJ mol}^{-1}$, determined from rotating bomb combustion calorimetry, and the enthalpy of sublimation, $\Delta_{\text{sub}} H^\circ(\text{C}_2\text{H}_3\text{BrO}_2) = (83.50 \pm 2.95)\text{ kJ mol}^{-1}$, determined from Knudsen effusion experiments. Note that the recommended value is close to the value predicted by the method of group equations:²

$$\begin{aligned} \Delta_f H^\circ(\text{CH}_2\text{Br}-\text{COOH}) \\ &= \Delta_f H^\circ(\text{CH}_2\text{Cl}-\text{COOH}) + \Delta_f H^\circ(\text{CH}_2\text{Br}-\text{CH}_3) \\ &\quad - \Delta_f H^\circ(\text{CH}_2\text{Cl}-\text{CH}_3) = (-427.6) + (-61.9) \\ &\quad - (-112.1) = -377.4 \text{ kJ mol}^{-1} \end{aligned}$$

(the $\Delta_f H^\circ$ values for CH_2BrCH_3 , and CH_2ClCH_3 were taken from compilation by Pedley,³ for CH_2ClCOOH —from Ref. 1). A somewhat lower estimate of the enthalpy of formation of bromoacetic acid, $-(395 \pm 6) \text{ kJ mol}^{-1}$ was given by Lias *et al.*⁴

2.2. Heat Capacity and Entropy

From their microwave study, van Eijck *et al.*⁵ determined the rotational constants of three isotopic species ($\text{CH}_2^{79}\text{BrCOOH}$, $\text{CH}_2^{81}\text{BrCOOH}$, and $\text{CH}_2^{79}\text{BrCOOD}$). Although no complete structure could be evaluated from the available data, the substitution coordinates of the Br atom and the carboxyl H atom were consistent only with the *trans* structure with respect to the atoms Br—C—C—O—H. This conformation is identical to lowest-energy form of chloroacetic acid named as *cis-syn* because of *cis* configuration for O=C—O—H and Cl—C—C=O groups. The *cis* structure with respect to the atoms Br—C—C—O—H was obtained by Chen *et al.*⁶ from *ab initio* calculation. However, the geometry of chloroacetic acid calculated by Chen *et al.*⁶ was also not consistent with that determined from experimental studies. The *trans* structure with respect to the atoms Br—C—C—O—H (C_s symmetry) is accepted in this work for lowest-energy conformer of bromoacetic acid in accord with the microwave data.⁵ The isotope-weighted value of the product of the principal moments of inertia of bromoacetic acid is calculated in this work from the rotational constants for $\text{CH}_2^{79}\text{BrCOOH}$ and $\text{CH}_2^{81}\text{BrCOOH}$.⁵ Structural parameters given above are those estimated by comparison with structural parameters of CH_3COOH ,⁷ CH_2ClCH_3 ,⁸

CH_2ClCOOH ,⁹ and CH_2BrCH_3 .¹⁰ These structural parameters yield values for the rotational constants which are 1.2%–1.6% different from the observed values used in the calculations. The difference has a negligible effect on the thermodynamic functions.

There is no information on other stable conformers of bromoacetic acid arising from internal rotation around the C—C bond. van Eijck *et al.*⁵ could only conclude that the *gauche* conformation, if present, is not significantly higher in energy than the *trans* conformation. Their rough estimate of the torsional frequency, 47 cm^{-1} , may be compared with 62 cm^{-1} for chloroacetic acid.¹¹ As in the case of chloroacetic acid, the simple potential,

$$V(\varphi) = \frac{1}{2}V_3(1 - \cos 3\varphi),$$

where φ is the Br—C—C=O torsional angle, is used in this work to calculate the internal rotational contributions to the thermodynamic functions of bromoacetic acid. The barrier height for the rotation about the C—C bond is practically the same in CH_2ClCH_3 and CH_2BrCH_3 molecules.^{12,13} For that reason, the value of V_3 for bromoacetic acid was accepted to be the same as that for chloroacetic acid. The value of the reduced moment of inertia for the CH_2Br top was derived from structural parameters adopted in this work (see above).

Vibrational spectra of bromoacetic acid were investigated only for a solid phase.^{14–17} These vibrational assignments are incomplete and it may be expected that they are much different from a gaseous spectrum as in the case of chloroacetic acid. Fundamental frequencies of gaseous bromoacetic acid were estimated in this work by normal coordinate calculations using the force constants transferred from related compounds. Simplified force fields for CH_3COOH , CH_2ClCH_3 , CH_2BrCH_3 , and CH_2ClCOOH were determined using experimental vibrational assignments for these molecules.^{18–21} 37 force constants were used to calculate the vibrational frequencies of bromoacetic acid:

| | | | | | |
|---|--------|---|--------|---|--------|
| $f_{\text{O}-\text{H}}$ | 7.092 | $f_{\text{wag}(\text{C}=\text{O})}$ | 0.310 | $f_{\text{C}-\text{C}-\text{C}-\text{H}}$ | 0.199 |
| $f_{\text{C}-\text{H}}$ | 4.998 | $f_{\text{tors}(\text{C}-\text{O})}$ | 0.183 | $f_{\text{C}-\text{C}-\text{C}-\text{O}}$ | -0.425 |
| $f_{\text{C}=\text{O}}$ | 14.076 | $f_{\text{tors}(\text{C}-\text{C})}$ | 0.028 | $f_{\text{C}-\text{C}-\text{C}-\text{Br}}$ | -0.017 |
| $f_{\text{C}-\text{C}}$ | 4.101 | $f_{\text{C}-\text{H},\text{C}-\text{H}}$ | 0.061 | $f_{\text{C}-\text{C},\text{O}=\text{C}-\text{C}}=f_{\text{C}-\text{C},\text{O}=\text{C}-\text{O}}$ | 1.250 |
| $f_{\text{C}-\text{O}}$ | 3.891 | $f_{\text{C}-\text{H},\text{C}-\text{C}}$ | 0.332 | $f_{\text{C}-\text{O},\text{C}-\text{O}-\text{H}}$ | 0.041 |
| $f_{\text{C}-\text{Br}}$ | 3.531 | $f_{\text{C}-\text{H},\text{C}-\text{Br}}$ | -0.402 | $f_{\text{C}-\text{O},\text{C}-\text{C}-\text{O}}$ | 0.296 |
| $f_{\text{H}-\text{C}-\text{H}}$ | 0.274 | $f_{\text{C}-\text{C},\text{C}-\text{Br}}$ | 0.377 | $f_{\text{C}-\text{O},\text{O}=\text{C}-\text{C}}=f_{\text{C}-\text{O},\text{O}=\text{C}-\text{O}}$ | 0.605 |
| $f_{\text{C}-\text{C}-\text{H}}$ | 0.793 | $f_{\text{C}-\text{C},\text{C}-\text{O}}$ | 0.693 | $f_{\text{C}-\text{Br},\text{H}-\text{C}-\text{Br}}$ | 0.148 |
| $f_{\text{C}-\text{O}-\text{H}}$ | 0.424 | $f_{\text{C}=\text{O},\text{C}-\text{O}}$ | 2.221 | $f_{\text{C}-\text{Br},\text{C}-\text{C}-\text{Br}}$ | 0.044 |
| $f_{\text{C}-\text{C}-\text{O}}$ | 2.465 | $f_{\text{C}=\text{O},\text{C}-\text{C}}$ | 1.660 | $f_{\text{C}-\text{H},\text{H}-\text{C}-\text{Br}}$ | -0.533 |
| $f_{\text{H}-\text{C}-\text{Br}}$ | 0.547 | $f_{\text{C}=\text{O},\text{C}-\text{C}-\text{O}}$ | -1.365 | $f_{\text{C}-\text{H},\text{H}-\text{C}-\text{H}}$ | -0.184 |
| $f_{\text{C}-\text{C}-\text{Br}}$ | 1.042 | $f_{\text{C}=\text{O},\text{O}=\text{C}-\text{C}}=f_{\text{C}=\text{O},\text{O}=\text{C}-\text{O}}$ | 0.687 | $f_{\text{C}-\text{H},\text{C}-\text{C}-\text{H}}$ | 0.216 |
| $f_{\text{O}=\text{C}-\text{C}}=f_{\text{O}=\text{C}-\text{O}}$ | 2.312 | | | | |

(stretching and stretch–stretch interaction constants are in units of $\text{mdyn}/\text{\AA}$; bend, wagging, and torsion constants are in units of $\text{mdyn}/\text{\AA}$; stretch–bend interaction constants are in units of mdyn). These constants were transferred from CH_3COOH and CH_2BrCH_3 molecules with corrections made by analyzing the trends in force constants of molecules CH_3COOH , CH_2ClCH_3 , and CH_2ClCOOH .

The uncertainties in the calculated thermodynamic functions (Table 1) may reach $(3\text{--}6) \text{ J K}^{-1} \text{ mol}^{-1}$ for $C_p^\circ(T)$ and $(5\text{--}12) \text{ J K}^{-1} \text{ mol}^{-1}$ for $S^\circ(T)$. They are caused by the uncertainties in the adopted vibrational frequencies and the approximate treatment of internal rotation.

Ideal gas thermodynamic properties of bromoacetic acid have not been reported previously.

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TABLE 1. Ideal gas thermodynamic properties of bromoacetic acid $C_2H_3BrO_2(g)$ at the standard state pressure, $p^\circ = 0.1 \text{ MPa}$ ($T_r = 298.15 \text{ K}$)

| T (K) | C_p° ($J K^{-1} mol^{-1}$) | S° ($J K^{-1} mol^{-1}$) | $-[G^\circ - H^\circ(T_r)]/T$ ($J K^{-1} mol^{-1}$) | $H^\circ - H^\circ(T_r)$ ($kJ mol^{-1}$) | $\Delta_f H^\circ$ ($kJ mol^{-1}$) | $\Delta_f G^\circ$ ($kJ mol^{-1}$) | $\log K_f^\circ$ |
|------------|--|--------------------------------------|--|---|---|---|------------------|
| 0 | 0.000 | 0.000 | ∞ | -16.862 | -364.616 | -364.616 | ∞ |
| 25 | 36.587 | 211.762 | 852.075 | -16.008 | -365.365 | -365.150 | 762.919 |
| 50 | 41.371 | 238.698 | 539.288 | -15.029 | -366.712 | -364.466 | 380.746 |
| 75 | 45.202 | 256.212 | 442.164 | -13.946 | -368.184 | -363.016 | 252.820 |
| 100 | 48.832 | 269.713 | 397.423 | -12.771 | -369.448 | -361.099 | 188.614 |
| 150 | 56.415 | 290.927 | 358.541 | -10.142 | -371.701 | -356.428 | 124.116 |
| 180 | 61.214 | 301.635 | 348.180 | -8.378 | -372.997 | -353.251 | 102.508 |
| 190 | 62.843 | 304.988 | 345.818 | -7.758 | -373.427 | -352.142 | 96.808 |
| 200 | 64.482 | 308.253 | 343.859 | -7.121 | -373.854 | -351.012 | 91.672 |
| 210 | 66.128 | 311.439 | 342.239 | -6.468 | -374.281 | -349.858 | 87.020 |
| 220 | 67.778 | 314.553 | 340.910 | -5.799 | -374.706 | -348.686 | 82.786 |
| 230 | 69.431 | 317.602 | 339.831 | -5.113 | -375.130 | -347.493 | 78.916 |
| 240 | 71.083 | 320.592 | 338.967 | -4.410 | -375.553 | -346.282 | 75.364 |
| 250 | 72.731 | 323.527 | 338.291 | -3.691 | -375.975 | -345.055 | 72.094 |
| 260 | 74.374 | 326.412 | 337.779 | -2.955 | -376.398 | -343.809 | 69.070 |
| 270 | 76.007 | 329.250 | 337.411 | -2.204 | -382.140 | -342.466 | 66.252 |
| 280 | 77.630 | 332.043 | 337.169 | -1.435 | -382.632 | -340.987 | 63.610 |
| 290 | 79.240 | 334.795 | 337.040 | -0.651 | -383.115 | -339.491 | 61.147 |
| 298.15 | 80.540 | 337.010 | 337.010 | 0.000 | -383.500 | -338.262 | 59.261 |
| 300 | 80.834 | 337.509 | 337.011 | 0.149 | -383.587 | -337.981 | 58.846 |
| 350 | 88.515 | 350.553 | 338.023 | 4.385 | -400.269 | -329.432 | 49.164 |
| 400 | 95.600 | 362.842 | 340.365 | 8.991 | -401.362 | -319.235 | 41.687 |
| 450 | 102.021 | 374.480 | 343.514 | 13.934 | -402.305 | -308.911 | 35.857 |
| 500 | 107.785 | 385.532 | 347.168 | 19.182 | -403.114 | -298.489 | 31.182 |
| 600 | 117.569 | 406.081 | 355.302 | 30.468 | -404.395 | -277.439 | 24.153 |
| 700 | 125.486 | 424.820 | 363.915 | 42.634 | -405.305 | -256.204 | 19.118 |
| 800 | 132.017 | 442.016 | 372.617 | 55.519 | -405.904 | -234.859 | 15.334 |
| 900 | 137.498 | 457.891 | 381.222 | 69.002 | -406.239 | -213.457 | 12.388 |
| 1000 | 142.157 | 472.625 | 389.634 | 82.991 | -406.354 | -192.028 | 10.030 |
| 1100 | 146.150 | 486.366 | 397.810 | 97.411 | -406.281 | -170.598 | 8.101 |
| 1200 | 149.592 | 499.234 | 405.732 | 112.203 | -406.053 | -149.181 | 6.494 |
| 1300 | 152.573 | 511.328 | 413.394 | 127.315 | -405.702 | -127.785 | 5.134 |
| 1400 | 155.164 | 522.732 | 420.800 | 142.704 | -405.257 | -106.425 | 3.971 |
| 1500 | 157.423 | 533.516 | 427.958 | 158.336 | -404.734 | -85.097 | 2.963 |
| 1600 | 159.399 | 543.741 | 434.879 | 174.180 | -404.160 | -63.786 | 2.082 |
| 1700 | 161.134 | 553.457 | 441.570 | 190.208 | -403.549 | -42.535 | 1.307 |
| 1800 | 162.662 | 562.712 | 448.045 | 206.400 | -402.914 | -21.313 | 0.618 |
| 1900 | 164.013 | 571.543 | 454.315 | 222.735 | -402.270 | -0.130 | 0.004 |
| 2000 | 165.210 | 579.987 | 460.389 | 239.197 | -401.625 | 21.019 | -0.549 |

3. Chloroacetic Acid, $CH_2Cl-COOH$ Chloroacetic acid ($C_2H_3ClO_2$)

Ideal gas

 $M_r = 94.4975$ $S^\circ(298.15 \text{ K}) = 325.9 \pm 5.0 \text{ J K}^{-1} \text{ mol}^{-1}$
 $\Delta_f H^\circ(0 \text{ K}) = -416.0 \pm 1.0 \text{ kJ mol}^{-1}$
 $\Delta_f H^\circ(298.15 \text{ K}) = -427.6 \pm 1.0 \text{ kJ mol}^{-1}$

Molecular constants

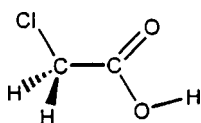
Point group: C_s Symmetry number: $\sigma = 1$ Ground electronic state: \tilde{X}^1A Energy: $\epsilon_X = 0 \text{ cm}^{-1}$ Quantum weight: $g_X = 1$

| | | Vibrational frequencies, ν_i , and degeneracies g_i | | | | | |
|----------|---------|---|-------|----------|-------------------------|------------------|---|
| Symmetry | | ν_i, cm^{-1} | g_i | Symmetry | ν_i, cm^{-1} | g_i | |
| A' | ν_1 | 3566 | 1 | A'' | ν_{10} | 596 | |
| | ν_2 | 3019 | 1 | | ν_{11} | 397 | 1 |
| | ν_3 | 1806 | 1 | | ν_{12} | 216 | 1 |
| | ν_4 | 1428 | 1 | | ν_{13} | 3076 | 1 |
| | ν_5 | 1354 | 1 | | ν_{14} | 1193 | 1 |
| | ν_6 | 1274 | 1 | | ν_{15} | 929 | 1 |
| | ν_7 | 1111 | 1 | | ν_{16} | 611 | 1 |
| | ν_8 | 891 | 1 | | ν_{17} | 492 | 1 |
| | ν_9 | 792 | 1 | | ν_{18} | ... ^a | 1 |

^aInstead of torsional mode $\nu_{18}=62 \text{ cm}^{-1}$, the contributions due to the internal rotation about C—C bond were calculated from the potential $V(\varphi)=\frac{1}{2}V_3(1-\cos 3\varphi)$, where φ is the torsional angle and $V_3=450 \text{ cm}^{-1}$.

CH₂Cl top: Reduced moment of inertia, $I_r=2.4514\times 10^{-39} \text{ g cm}^2$, Symmetry number, $\sigma_m=1$.

Geometry



$$\begin{aligned} r(\text{C—C}) &= 1.508 \pm 0.006 \text{ \AA} \\ r(\text{C=O}) &= 1.223 \pm 0.004 \text{ \AA} \\ r(\text{C—O}) &= 1.352 \pm 0.005 \text{ \AA} \\ r(\text{C—Cl}) &= 1.778 \pm 0.005 \text{ \AA} \end{aligned}$$

$$\begin{aligned} r(\text{C—H}) &= 1.09 \pm 0.02 \text{ \AA} \\ r(\text{O—H}) &= 0.97 \pm 0.015 \text{ \AA} \\ \angle \text{C—C=O} &= 126.1 \pm 0.5^\circ \\ \angle \text{C—C—O} &= 110.6 \pm 0.4^\circ \\ \angle \text{C—C—Cl} &= 112.5 \pm 0.4^\circ \\ \angle \text{C—C—H} &= 109.5 \text{ (assumed)} \\ \angle \text{H—C—H} &= 109.5 \text{ (assumed)} \\ \angle \text{C—O—H} &= 105.8 \pm 1.1^\circ \\ \varphi(\text{O=C—C—Cl}) &= 0.0^\circ \\ \varphi(\text{O=C—O—H}) &= 0.0^\circ \end{aligned}$$

Rotational constants in cm^{-1} :

$$A_0 = 0.350738 \quad B_0 = 0.078433 \quad C_0 = 0.064913$$

Product of moments of inertia: $I_A I_B I_C = 12284 \times 10^{-117} \text{ g}^3 \text{ cm}^6$.

3.1. Enthalpy of Formation

The recommended value of the enthalpy of formation of gaseous chloroacetic acid at 298.15 K is the sum of the enthalpy of formation of the α form of the solid and the enthalpy of sublimation both at 298.15 K. The value of $\Delta_f H^\circ(298.15 \text{ K}, \alpha\text{-cr}) = -(509.74 \pm 0.49) \text{ kJ mol}^{-1}$ is from the work of Lagoa *et al.*¹ who measured the enthalpy of combustion of the α form of the solid with a rotating bomb calorimeter. Their result is in agreement with the result $\Delta_f H^\circ(298.15 \text{ K}, \text{cr}) = -(510.5 \pm 8.3) \text{ kJ mol}^{-1}$ of Smith *et al.*² from their re-evaluation of earlier static bomb calorimetry measurements.³

The enthalpy of sublimation, $\Delta_{\text{sub}} H^\circ(298.15 \text{ K}) = (82.19 \pm 0.92) \text{ kJ mol}^{-1}$, used here is also from Lagoa *et al.*,¹ and is derived from vapor pressures of the solid obtained from Knudsen effusion experiments. This value is supported by application of Hess's law where the enthalpy of sublimation is the sum of five processes at the standard pressure of 1 bar. Thus,

$$\Delta_{\text{sub}} H^\circ(298.15 \text{ K})$$

$$\begin{aligned} &= [H^\circ(\text{cr}, 334.8 \text{ K}) - H^\circ(\text{cr}, 298.15 \text{ K})] \\ &\quad + \Delta_{\text{fus}} H^\circ(334.8 \text{ K}) + [H^\circ(1,462 \text{ K}) - H^\circ(1,334.8 \text{ K})] \\ &\quad + \Delta_{\text{vap}} H^\circ(462 \text{ K}) - [H^\circ(\text{g}, 462 \text{ K}) \\ &\quad - H^\circ(\text{cr}, 298.15 \text{ K})] = 80.9 \pm 2.3 \text{ kJ mol}^{-1}. \end{aligned}$$

The first and third terms $[H^\circ(T_2) - H^\circ(T_1)] \approx C_p^\circ(\text{cr}) \times \Delta T$ are (3.9 ± 0.1) and $(21.6 \pm 0.5) \text{ kJ mol}^{-1}$, respectively, based on $C_p^\circ(\text{cr}) = (106.7 \pm 2.0) \text{ J K}^{-1} \text{ mol}^{-1}$ from differential scanning calorimetry (DSC) measurements¹ and $C_p^\circ(1) = (168.9 \pm 4.0) \text{ J K}^{-1} \text{ mol}^{-1}$ from Pickering.⁴ [This value of $C_p^\circ(\text{cr})$ differs significantly from the average value of $144 \text{ J K}^{-1} \text{ mol}^{-1}$ over the range of 288–318 K reported in Pickering⁴ and adopted in NIST Chemistry WebBook⁵ and Donalski and Hearing.⁶ The value of $C_p^\circ(1) = 179.9 \text{ J K}^{-1} \text{ mol}^{-1}$ from Urazovskii and Sidorov⁷ was not used.] $\Delta_{\text{fus}} H^\circ(334.8 \text{ K}) = 16.3 \pm 0.7 \text{ kJ mol}^{-1}$ from DSC measurements¹ while $\Delta_{\text{vap}} H^\circ(462 \text{ K}) = (54.5 \pm 2.0) \text{ kJ mol}^{-1}$

TABLE 2. Ideal gas thermodynamic properties of chloroacetic acid $C_2H_3ClO_2(g)$ at the standard state pressure, $p^\circ = 0.1 \text{ MPa}$ ($T_r = 298.15 \text{ K}$)

| T (K) | C_p° ($J K^{-1} mol^{-1}$) | S° ($J K^{-1} mol^{-1}$) | $-(G^\circ - H^\circ(T_r))/T$ ($J K^{-1} mol^{-1}$) | $H^\circ - H^\circ(T_r)$ ($kJ mol^{-1}$) | $\Delta_f H^\circ$ ($kJ mol^{-1}$) | $\Delta_f G^\circ$ ($kJ mol^{-1}$) | $\log K_f^\circ$ |
|------------|--|--------------------------------------|--|---|---|---|------------------|
| 0 | 0.000 | 0.000 | ∞ | -16.514 | -416.032 | -416.032 | ∞ |
| 25 | 36.170 | 203.287 | 829.838 | -15.664 | -417.069 | -414.897 | 866.858 |
| 50 | 40.498 | 229.818 | 523.832 | -14.701 | -418.460 | -412.226 | 430.639 |
| 75 | 44.011 | 246.902 | 428.827 | -13.644 | -419.862 | -408.793 | 284.702 |
| 100 | 47.641 | 260.056 | 385.047 | -12.499 | -420.999 | -404.926 | 211.506 |
| 150 | 55.261 | 280.802 | 346.988 | -9.928 | -422.890 | -396.468 | 138.059 |
| 180 | 59.960 | 291.292 | 336.846 | -8.200 | -423.925 | -391.086 | 113.487 |
| 190 | 61.543 | 294.576 | 334.536 | -7.592 | -424.261 | -389.252 | 107.010 |
| 200 | 63.133 | 297.773 | 332.618 | -6.969 | -424.594 | -387.401 | 101.176 |
| 210 | 64.730 | 300.892 | 331.033 | -6.330 | -424.922 | -385.533 | 95.894 |
| 220 | 66.332 | 303.940 | 329.733 | -5.674 | -425.246 | -383.649 | 91.088 |
| 230 | 67.938 | 306.924 | 328.676 | -5.003 | -425.565 | -381.752 | 86.696 |
| 240 | 69.546 | 309.849 | 327.831 | -4.316 | -425.880 | -379.841 | 82.668 |
| 250 | 71.155 | 312.721 | 327.169 | -3.612 | -426.190 | -377.915 | 78.959 |
| 260 | 72.762 | 315.543 | 326.668 | -2.892 | -426.493 | -375.978 | 75.533 |
| 270 | 74.365 | 318.319 | 326.308 | -2.157 | -426.792 | -374.030 | 72.359 |
| 280 | 75.962 | 321.053 | 326.072 | -1.405 | -427.085 | -372.071 | 69.409 |
| 290 | 77.551 | 323.746 | 325.945 | -0.638 | -427.372 | -370.101 | 66.661 |
| 298.15 | 78.838 | 325.913 | 325.913 | 0.000 | -427.600 | -368.487 | 64.556 |
| 300 | 79.129 | 326.402 | 325.916 | 0.146 | -427.651 | -368.121 | 64.094 |
| 350 | 86.790 | 339.180 | 326.908 | 4.296 | -428.953 | -358.094 | 53.441 |
| 400 | 93.933 | 351.242 | 329.202 | 8.816 | -430.094 | -347.892 | 45.429 |
| 450 | 100.456 | 362.689 | 332.293 | 13.678 | -431.085 | -337.556 | 39.181 |
| 500 | 106.342 | 373.583 | 335.881 | 18.851 | -431.941 | -327.116 | 34.173 |
| 600 | 116.373 | 393.891 | 343.883 | 30.005 | -433.310 | -306.019 | 26.641 |
| 700 | 124.509 | 412.462 | 352.372 | 42.063 | -434.295 | -284.721 | 21.246 |
| 800 | 131.218 | 429.540 | 360.965 | 54.859 | -434.957 | -263.304 | 17.192 |
| 900 | 136.840 | 445.329 | 369.473 | 68.270 | -435.342 | -241.822 | 14.035 |
| 1000 | 141.610 | 460.000 | 377.801 | 82.199 | -435.499 | -220.311 | 11.508 |
| 1100 | 145.691 | 473.693 | 385.903 | 96.569 | -435.459 | -198.792 | 9.440 |
| 1200 | 149.204 | 486.524 | 393.760 | 111.318 | -435.259 | -177.285 | 7.717 |
| 1300 | 152.241 | 498.590 | 401.364 | 126.394 | -434.930 | -155.797 | 6.260 |
| 1400 | 154.877 | 509.971 | 408.719 | 141.753 | -434.500 | -134.342 | 5.012 |
| 1500 | 157.173 | 520.736 | 415.831 | 157.358 | -433.991 | -112.918 | 3.932 |
| 1600 | 159.180 | 530.946 | 422.709 | 173.178 | -433.426 | -91.512 | 2.987 |
| 1700 | 160.940 | 540.650 | 429.364 | 189.186 | -432.819 | -70.164 | 2.156 |
| 1800 | 162.490 | 549.894 | 435.805 | 205.359 | -432.185 | -48.846 | 1.417 |
| 1900 | 163.858 | 558.716 | 442.044 | 221.678 | -431.536 | -27.566 | 0.758 |
| 2000 | 165.071 | 567.153 | 448.090 | 238.125 | -430.883 | -6.322 | 0.165 |

is derived from vapor pressure data over a temperature range from 385.45 K to the normal boiling point.^{9,10} [Other values of $\Delta_{\text{fus}}H^\circ(334.8 \text{ K})$ are 16.3 kJ mol^{-1} from Pickering⁴ and 12.3 kJ mol^{-1} from Acree.⁸] The value [$H^\circ(g, 462 \text{ K}) - H^\circ(\text{cr}, 298.15 \text{ K})$] = (15.4 \pm 0.7) kJ mol^{-1} is interpolated from Table 2 of the present work. This corrects the estimate $\Delta_{\text{sub}}H^\circ(298.15) = (75.3 \pm 4.2)$ kJ mol^{-1} by Cox and Pilcher^{5,6,11} based on $\Delta_{\text{fus}}H^\circ(334 \text{ K}) = 19.4 \text{ kJ mol}^{-1}$ from Steiner and Johnson¹² and $\Delta_{\text{vap}}H^\circ(462 \text{ K}) = 54.5 \text{ kJ mol}^{-1}$ mentioned above. Cox and Pilcher also use Hess's law, but assume that the values of C_p° for the gas, liquid, and solid phases are equal.

3.2. Heat Capacity and Entropy

According to the experimental¹³⁻¹⁸ and theoretical^{15,17-19} studies, the lowest-energy conformer of chloroacetic acid

($\text{CH}_2\text{Cl}-\text{CO}-\text{OH}$) has the *cis-syn* structure (C_s symmetry) with the *cis* configuration for the carboxylic group and with the chlorine atom lying in the carboxylic plane eclipsed with the carbonyl group. The second more stable form corresponds to the *cis-gauche* structure with a $\text{Cl}-\text{C}-\text{C}=\text{O}$ angle of $\sim 130^\circ$ (C_1 symmetry). This form differs from *cis-syn* by internal rotation of the CH_2Cl group about the $\text{C}-\text{C}$ bond. Structural parameters of the most stable conformer of chloroacetic acid were determined by gas phase electron diffraction.¹⁴ These parameters are in good agreement with results from *ab initio*^{18,19} and molecular mechanics¹⁵ calculations. In this work, the product of the principal moments of inertia for the most stable *cis-syn* conformer of chloroacetic acid was calculated using rotational constants determined from the microwave study.¹³ Structural parameters given above are those obtained from the electron diffraction

study.¹⁴ These parameters reproduce the product of the principal moments of inertia calculated above within 3%.

Three conformations with respect to internal rotation around the C—C bond were found from an electron diffraction investigation,¹⁴ namely, 56% of a *cis-syn* conformation, 30% of a *cis-gauche* conformation with the CH₂Cl group rotated 131° from the former position, and the remaining 14% of a *cis-gauche* conformation with 79° rotation of the CH₂Cl group. Three conformers of chloroacetic acid were identified by vibrational spectroscopy.^{15–18} Two of them were *cis-syn* and *cis-gauche* in agreement with the electron diffraction data. The third stable form was found to have the *trans* structure, in which the carboxylic hydrogen atom is in the *trans* position with respect to the C=O bond. *Trans* conformers arise from rotation of the OH group about the C—O bond. *Ab initio* calculations^{17,18} and molecular mechanics studies¹⁵ strongly suggest that the third stable form should be the *trans* form. The barrier height for the rotation of the OH group around the C—O bond is predicted to be 1700–3500 cm⁻¹.^{16–18} Due to the height of this barrier and the temperature range of the present tabulation, this internal rotation was ignored in this work.

Two *cis* conformers with respect to internal rotation about the C—C bond were considered in this work: the *cis-syn* conformer of C_s symmetry and two enantiomeric forms of *cis-gauche* conformer of C₁ symmetry. The observed data and the calculations^{14,16–18} consistently predict a slight energy preference for the *cis-syn* form and a small barrier of (400–450) cm⁻¹ for its interconversion. The simple potential

$$V(\varphi) = \frac{1}{2}V_3(1 - \cos 3\varphi),$$

where φ is the Cl—C—C=O torsional angle, is used here for a very approximate calculation of the internal rotational contributions to the thermodynamic functions of chloroacetic acid. The value of the reduced moment of inertia for the CH₂Cl top was derived from the electron diffraction structural parameters.¹⁴

Vibrational spectra of chloroacetic acid were studied in the liquid and solid phases,^{15,20,21} in the vapor²² and matrix.^{16,18} The fundamental frequencies adopted in this work are those derived by Nieminen *et al.*¹⁸ from matrix isolation infrared spectra ($\nu_1, \nu_3 - \nu_5, \nu_7, \nu_8, \nu_9, \nu_{15} - \nu_{17}$) and *ab initio* calculation ($\nu_2, \nu_6, \nu_{10} - \nu_{14}$). These frequencies are in good agreement with results of normal coordinate analysis,²² molecular mechanics,¹⁵ and *ab initio*¹⁷ calculations. The value for the torsional frequency, ν_{18} , was estimated from the microwave spectrum¹³ and it coincides with the value calculated by the *ab initio* method.¹⁸

The uncertainties in the calculated thermodynamic functions (Table 2) may amount to as much as (3–5) J K⁻¹ mol⁻¹ for C_p^o(T) and 5–10 J K⁻¹ mol⁻¹ for S^o(T). They are caused by the uncertainties in the adopted vibrational frequencies and the approximate treatment of the internal rotation.

Thermodynamic properties of chloroacetic acid were calculated earlier by Banerjee²³ using molecular constants known at that time. A value for the barrier height of ~1750

cm⁻¹ was adopted for calculating the internal rotation contributions of the CH₂Cl and OH groups. The difference between the values of C_p^o(T) and S^o(T) given here and those by Banerjee²³ amounts to 57 and 35 J K⁻¹ mol⁻¹, respectively. Such a difference could not be due to the discrepancy in molecular constants used. The calculation of Banerjee²³ seems to be in error. The rough estimate by the method of group equations,

$$\begin{aligned} C_p^o(\text{CH}_2\text{Cl}-\text{COOH}) &= C_p^o(\text{CH}_3-\text{COOH}) + C_p^o(\text{CH}_2\text{Cl}-\text{CH}_3) \\ &\quad - C_p^o(\text{CH}_3-\text{CH}_3) \\ &= 63.4 + 62.6 - 52.5 = 73.5 \text{ J K}^{-1} \text{ mol}^{-1}, \end{aligned}$$

is close to the value obtained in this work (78.8 J K⁻¹ mol⁻¹) and is very different from the value of Banerjee²³ (136.0 J K⁻¹ mol⁻¹). The calculation of Banerjee²³ is reproduced in the reference book of Frenkel *et al.*²⁴

3.3. References

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4. Oxopropanedinitrile, NC—CO—CN

Oxopropanedinitrile (C₃N₂O)

Ideal gas

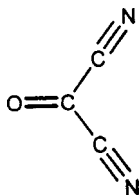
 $M_r = 80.0458$ $S^\circ(298.15\text{ K}) = 310.0 \pm 1.0\text{ J K}^{-1}\text{ mol}^{-1}$ $\Delta_f H^\circ(0\text{ K}) = 246.5 \pm 6.4\text{ kJ mol}^{-1}$ $\Delta_f H^\circ(298.15\text{ K}) = 247.5 \pm 6.4\text{ kJ mol}^{-1}$

Molecular constants

Point group: C_{2v} Symmetry number: $\sigma = 2$ Ground electronic state: \tilde{X}^1A_1 Energy: $\epsilon_X = 0\text{ cm}^{-1}$ Quantum weight: $g_X = 1$ Vibrational frequencies, ν_i , and degeneracies g_i

| Symmetry | ν_i , cm ⁻¹ | g_i | Symmetry | ν_i , cm ⁻¹ | g_i |
|----------------|----------------------------|-------|----------------|----------------------------|-------|
| A ₁ | ν_1 2230 | 1 | B ₁ | ν_7 712 | 1 |
| | ν_2 1711 | 1 | | ν_8 208.2 | 1 |
| | ν_3 712 | 1 | B ₂ | ν_9 2230 | 1 |
| | ν_4 553 | 1 | | ν_{10} 1124 | 1 |
| | ν_5 127.5 | 1 | | ν_{11} 550 | 1 |
| A ₂ | ν_6 307 | 1 | | ν_{12} 245.2 | 1 |

Geometry

 $r(\text{C}=\text{O}) = 1.204 \pm 0.005\text{ \AA}$ $r(\text{C}-\text{C}) = 1.461 \pm 0.005\text{ \AA}$ $r(\text{C}\equiv\text{N}) = 1.159 \pm 0.015\text{ \AA}$ $\angle \text{C}-\text{C}-\text{C} = 114.7 \pm 0.5^\circ$ $\angle \text{C}-\text{C}\equiv\text{N} = 179.2 \pm 0.5^\circ$ Rotational constants in cm⁻¹: $A_0 = 0.225\,529$ $B_0 = 0.097\,556$ $C_0 = 0.067\,980$.Product of moments of inertia: $I_A I_B I_C = 14\,666 \times 10^{-117}\text{ g}^3\text{ cm}^6$.

4.1. Enthalpy of Formation

The recommended value of enthalpy of formation of oxopropanedinitrile is based on calorimetric measurements by von Glemser and Häusser¹ as evaluated by Cox and Pilcher.²

4.2. Heat Capacity and Entropy

Spectroscopic,³⁻¹⁰ electron diffraction,¹¹ and theoretical¹²⁻²⁰ investigations have shown that oxopropanedinitrile, CO(CN)₂, is planar in its ground electronic state \tilde{X}^1A_1 and belongs to the C_{2v} symmetry group. In this work, the product of the principal moments of inertia of CO(CN)₂ was calculated using the rotational constants determined by microwave spectroscopy.⁴ In the absence of isotopic data, a unique set of geometrical parameters cannot be obtained from the microwave spectrum. Structural parameters given above are r_e parameters determined by combining the results of electron diffraction, microwave spectroscopy, and *ab initio* calculations.²⁰ These parameters give values for rotational constants which are only 0.3%–0.9% different from the observed values. The C—C≡N chain appears to be nearly lin-

ear, the deviation from linearity being 0.8°. A small inward bend (0.5°–2°) was also found by *ab initio* calculations^{12,13,19} but is contradicted by *ab initio* calculations of Tyrrell¹⁵ where the C—C≡N is bent outwards by 1.2°.

Vibrational spectra of oxopropanedinitrile were studied in the gas, liquid, and solid phase⁵⁻¹⁰ but there are still several uncertainties in the assignment of the fundamentals. The vibrational frequencies accepted in this work are those assigned by Miller *et al.*⁹ from infrared and Raman spectra of gaseous and liquid oxopropanedinitrile. Their assignment was supported by *ab initio* calculation.¹⁵

According to the semiempirical calculation,²¹ the excited electronic states of CO(CN)₂ lie above 24 000 cm⁻¹. They are not taken into account in the calculation of thermodynamic functions.

The uncertainties in the calculated thermodynamic functions (Table 3) are estimated to be (1–2) J K⁻¹ mol⁻¹ for $C_p^\circ(T)$ and (1–2.5) J K⁻¹ mol⁻¹ for $S^\circ(T)$.

Thermodynamic properties of CO(CN)₂ were calculated earlier by Natarajan and Rajendran²² using electron diffrac-

tion structural parameters¹¹ and the vibrational assignment of Bates and Smith.⁸ The numerical data in the four columns of the table of thermodynamic functions in Natarajan and Rajendran²² are transposed. Moreover, these functions do not correspond to the molecular constants used by the authors. The difference between the $C_p^\circ(T)$ and $S^\circ(T)$ values given here and those by Natarajan and Rajendran²² amounts to 20 $\text{J K}^{-1} \text{mol}^{-1}$ and could not be due to the discrepancy in molecular constants used.

4.3. References

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TABLE 3. Ideal gas thermodynamic properties of oxopropanedinitrile C₃N₂O(g) at the standard state pressure, $p^\circ=0.1$ MPa ($T_r=298.15$ K)

| T (K) | C_p° (J K ⁻¹ mol ⁻¹) | S° (J K ⁻¹ mol ⁻¹) | $-(G^\circ - H^\circ(T_r))/T$ (J K ⁻¹ mol ⁻¹) | $H^\circ - H^\circ(T_r)$ (kJ mol ⁻¹) | $\Delta_f H^\circ$ (kJ mol ⁻¹) | $\Delta_f G^\circ$ (kJ mol ⁻¹) | $\log K_{fs}^\circ$ |
|------------|---|---|---|---|---|---|---------------------|
| 0 | 0.000 | 0.000 | ∞ | -17.147 | 246.525 | 246.525 | ∞ |
| 25 | 33.558 | 186.154 | 838.755 | -16.315 | 246.274 | 246.273 | -514.553 |
| 50 | 37.466 | 210.389 | 519.096 | -15.435 | 246.040 | 246.369 | -257.376 |
| 75 | 43.512 | 226.703 | 419.028 | -14.424 | 245.902 | 246.567 | -171.722 |
| 100 | 49.389 | 240.041 | 372.658 | -13.262 | 245.872 | 246.796 | -128.911 |
| 150 | 59.439 | 262.030 | 332.251 | -10.533 | 246.052 | 247.230 | -86.092 |
| 180 | 64.718 | 273.341 | 321.506 | -8.670 | 246.268 | 247.447 | -71.806 |
| 190 | 66.369 | 276.884 | 319.064 | -8.014 | 246.354 | 247.510 | -68.044 |
| 200 | 67.965 | 280.329 | 317.042 | -7.342 | 246.444 | 247.569 | -64.657 |
| 210 | 69.506 | 283.683 | 315.374 | -6.655 | 246.539 | 247.623 | -61.592 |
| 220 | 70.992 | 286.951 | 314.008 | -5.953 | 246.639 | 247.672 | -58.804 |
| 230 | 72.425 | 290.138 | 312.901 | -5.235 | 246.742 | 247.716 | -56.257 |
| 240 | 73.804 | 293.250 | 312.018 | -4.504 | 246.848 | 247.756 | -53.922 |
| 250 | 75.131 | 296.290 | 311.328 | -3.760 | 246.956 | 247.792 | -51.773 |
| 260 | 76.408 | 299.262 | 310.807 | -3.002 | 247.067 | 247.823 | -49.788 |
| 270 | 77.636 | 302.169 | 310.434 | -2.232 | 247.179 | 247.851 | -47.949 |
| 280 | 78.818 | 305.013 | 310.189 | -1.449 | 247.292 | 247.873 | -46.241 |
| 290 | 79.956 | 307.799 | 310.059 | -0.655 | 247.407 | 247.892 | -44.650 |
| 298.15 | 80.853 | 310.028 | 310.028 | 0.000 | 247.500 | 247.905 | -43.431 |
| 300 | 81.052 | 310.529 | 310.029 | 0.150 | 247.521 | 247.907 | -43.164 |
| 350 | 85.979 | 323.404 | 311.035 | 4.329 | 248.092 | 247.925 | -37.000 |
| 400 | 90.165 | 335.165 | 313.327 | 8.735 | 248.634 | 247.864 | -32.367 |
| 450 | 93.807 | 346.000 | 316.363 | 13.337 | 249.125 | 247.738 | -28.756 |
| 500 | 97.037 | 356.054 | 319.835 | 18.109 | 249.559 | 247.560 | -25.862 |
| 600 | 102.566 | 374.251 | 327.420 | 28.099 | 250.253 | 247.090 | -21.511 |
| 700 | 107.123 | 390.415 | 335.286 | 38.590 | 250.757 | 246.521 | -18.395 |
| 800 | 110.898 | 404.974 | 343.102 | 49.497 | 251.123 | 245.892 | -16.055 |
| 900 | 114.022 | 418.222 | 350.724 | 60.748 | 251.390 | 245.219 | -14.232 |
| 1000 | 116.608 | 430.373 | 358.090 | 72.284 | 251.851 | 244.520 | -12.772 |
| 1100 | 118.754 | 441.591 | 365.177 | 84.055 | 251.719 | 243.809 | -11.577 |
| 1200 | 120.542 | 452.003 | 371.984 | 96.023 | 251.810 | 243.084 | -10.581 |
| 1300 | 122.038 | 461.712 | 378.517 | 108.154 | 251.862 | 242.358 | -9.738 |
| 1400 | 123.297 | 470.804 | 384.788 | 120.422 | 251.871 | 241.623 | -9.015 |
| 1500 | 124.364 | 479.348 | 390.810 | 132.807 | 251.846 | 240.895 | -8.389 |
| 1600 | 125.273 | 487.404 | 396.598 | 145.290 | 251.779 | 240.197 | -7.842 |
| 1700 | 126.052 | 495.022 | 402.165 | 157.857 | 251.679 | 239.474 | -7.358 |
| 1800 | 126.724 | 502.247 | 407.526 | 170.497 | 251.541 | 238.765 | -6.929 |
| 1900 | 127.307 | 509.114 | 412.694 | 183.199 | 251.366 | 238.058 | -6.545 |
| 2000 | 127.815 | 515.658 | 417.680 | 195.956 | 251.156 | 237.363 | -6.199 |

5. Glycolic Acid, HO—CH₂—COOHGlycolic acid (C₂H₄O₃)

Ideal gas

 $M_r=76.0518$ $S^\circ(298.15\text{ K})=318.6\pm 5.0\text{ J K}^{-1}\text{ mol}^{-1}$ $\Delta_f H^\circ(0\text{ K})=-567.9\pm 10.0\text{ kJ mol}^{-1}$ $\Delta_f H^\circ(298.15\text{ K})=-583.0\pm 10.0\text{ kJ mol}^{-1}$

Molecular constants

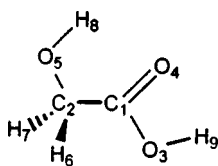
Point group: C_1 Symmetry number: $\sigma=1$ Number of optical isomers: $n=2$ Ground electronic state: \tilde{X}^1A Energy: $\epsilon_X=0\text{ cm}^{-1}$ Quantum weight: $g_X=1$

| Symmetry | | Vibrational frequencies, ν_i , and degeneracies g_i | | | | | |
|----------|------------|---|-------|----------|------------|-------------------------|-------|
| A' | | ν_i, cm^{-1} | g_i | Symmetry | | ν_i, cm^{-1} | g_i |
| | ν_1 | 3561 | 1 | | ν_{12} | 642 | 1 |
| | ν_2 | 3561 | 1 | | ν_{13} | 468 | 1 |
| | ν_3 | 2928 | 1 | | ν_{14} | 270 | 1 |
| | ν_4 | 1774 | 1 | A'' | ν_{15} | 2919 | 1 |
| | ν_5 | 1452 | 1 | | ν_{16} | 1231 | 1 |
| | ν_6 | 1439 | 1 | | ν_{17} | 1019 | 1 |
| | ν_7 | 1332 | 1 | | ν_{18} | 621 | 1 |
| | ν_8 | 1265 | 1 | | ν_{19} | 495 | 1 |
| | ν_9 | 1143 | 1 | | ν_{20} | 281 | 1 |
| | ν_{10} | 1090 | 1 | | ν_{21} | ... ^a | 1 |
| | ν_{11} | 854 | 1 | | | | |

^aInstead of torsional mode $\nu_{21}=152 \text{ cm}^{-1}$, the contributions due to the internal rotation round the C—C bond were calculated from the potential $V(\varphi) = \frac{1}{2} \sum_{n=1}^8 V_n (1 - \cos n\varphi)$, where φ is the torsional angle, $V_1 = 971.1$, $V_2 = 1192.6$, $V_3 = -196.0$, $V_4 = -244.1$, $V_5 = 137.1$, $V_6 = -173.5$, $V_7 = 118.8$, and $V_8 = -48.0$ (in cm^{-1}).

COOH top: Reduced moment of inertia, $I_r = 1.9292 \times 10^{-39} \text{ g cm}^2$, Symmetry number, $\sigma_m = 1$.

Geometry



| | |
|--|---|
| $r(\text{C}_1\text{—C}_2) = 1.495 \pm 0.006 \text{ \AA}$ | $\angle \text{C}_2\text{—C}_1\text{—O}_3 = 112.6 \pm 0.5^\circ$ |
| $r(\text{C}_1\text{—O}_3) = 1.349 \pm 0.006 \text{ \AA}$ | $\angle \text{C}_2\text{—C}_1\text{—O}_4 = 124.2 \pm 0.4^\circ$ |
| $r(\text{C}_1\text{=O}_4) = 1.210 \pm 0.006 \text{ \AA}$ | $\angle \text{C}_1\text{—C}_2\text{—O}_5 = 111.3 \pm 0.4^\circ$ |
| $r(\text{C}_2\text{—O}_5) = 1.406 \pm 0.004 \text{ \AA}$ | $\angle \text{O}_5\text{—C}_2\text{—H}_{6,7} = 110.6 \pm 0.3^\circ$ |
| $r(\text{O}_3\text{—H}_9) = 0.989 \pm 0.019 \text{ \AA}$ | $\angle \text{C}_1\text{—C}_2\text{—H}_{6,7} = 108.8 \pm 0.3^\circ$ |
| $r(\text{O}_5\text{—H}_8) = 0.956 \pm 0.003 \text{ \AA}$ | $\angle \text{C}_1\text{—O}_3\text{—H}_9 = 105.5 \pm 1.1^\circ$ |
| $r(\text{C}_2\text{—H}_{6,7}) = 1.097 \pm 0.003 \text{ \AA}$ | $\angle \text{C}_2\text{—O}_5\text{—H}_8 = 105.2 \pm 0.1^\circ$ |

Rotational constants in cm^{-1} :

$$A_0 = 0.356783 \quad B_0 = 0.135128 \quad C_0 = 0.099891.$$

Product of moments of inertia: $I_A I_B I_C = 4555 \times 10^{-117} \text{ g}^3 \text{ cm}^6$.

| Other stable conformers: | Point group | Symmetry number, σ | Number of optical isomers, n | Energy, cm^{-1} |
|--------------------------|-------------|---------------------------|--------------------------------|--------------------------|
| | C_1 | 1 | 2 | 1200 |
| | C_1 | 1 | 2 | 1300 |

5.1. Enthalpy of Formation

No experimental or theoretical data on enthalpy of formation of gaseous glycolic acid are known from the literature. The value accepted in this work, $\Delta_f H^\circ$ (298.15 K) = $-(583 \pm 10) \text{ kJ mol}^{-1}$, is based on two estimates by additivity methods.

The first value was estimated by group additivity using the equation

$$\begin{aligned} \Delta_f H^\circ(\text{HO—CH}_2\text{—COOH}) &= \Delta_f H^\circ[\text{O—(H)(C)}] + \Delta_f H^\circ[\text{C—(H)}_2(\text{CO})(\text{O})] \\ &\quad + \Delta_f H^\circ[\text{CO—(C)(O)}] + \Delta_f H^\circ[\text{O—(CO)(H)}] \end{aligned}$$

$$= (-158.6) + (-33.5) + (-147.3) + (-241.8)$$

$$= -581.2 \text{ kJ mol}^{-1}$$

with group values generated by Cohen¹ except for the missing value for the $[\text{C—(H)}_2(\text{CO})(\text{O})]$ group. The latter was evaluated from values known for related groups:

$$\begin{aligned} \Delta_f H^\circ[\text{C—(H)}_2(\text{CO})(\text{O})] &= \Delta_f H^\circ[\text{C—(H)}_2(\text{CO})(\text{C})] \\ &\quad + \Delta_f H^\circ[\text{C—(H)(CO)(C)(O)}] \\ &\quad - \Delta_f H^\circ[\text{C—(H)(CO)(C)}_2] \end{aligned}$$

$$\begin{aligned}
 &= (-21.8) + (-18.8) \\
 &\quad - (-7.1) = -33.5 \text{ kJ mol}^{-1}.
 \end{aligned}$$

The other $\Delta_f H^\circ$ value of glycolic acid may be predicted by the method of group equations² using

$$\begin{aligned}
 \Delta_f H^\circ(\text{HO—CH}_2\text{—COOH}) \\
 &= \Delta_f H^\circ(\text{CH}_3\text{—COOH}) + \Delta_f H^\circ(\text{HO—CH}_2\text{—CH}_3) \\
 &\quad - \Delta_f H^\circ(\text{CH}_3\text{—CH}_3) \\
 &= (-432.8) + (-235.2) \\
 &\quad - (-83.8) = -584.2 \text{ kJ mol}^{-1}.
 \end{aligned}$$

Values for $\Delta_f H^\circ$ for CH_3COOH , $\text{CH}_3\text{CH}_2\text{OH}$, and C_2H_6 were taken from a compilation by Pedley.³ A value intermediate between the above two estimates was assigned to the enthalpy of formation of glycolic acid.

5.2. Heat Capacity and Entropy

From microwave spectroscopic studies,^{4–6} the lowest-energy structure of glycolic acid was concluded to be of C_s symmetry with the alcoholic hydroxyl group pointing toward the carbonyl oxygen of carboxyl group. This structure involves the intramolecular bonding between the alcohol hydrogen and the carbonyl oxygen. With the exception of the methylene hydrogen all the atoms in the molecule were found to be coplanar. The pathways and energy barriers involved in possible conformational interconversions of glycolic acid were investigated by *ab initio* calculations.^{7–12} Along with rotamers of *cis*-glycolic acid, where the carbonyl group has the *cis* conformation with its hydroxyl group, the conformers of *trans*-glycolic acid were predicted from theoretical studies. Based on electron diffraction data¹³ the second lowest conformer has been assigned to a *cis*-glycolic acid with hydrogen bonding between two hydroxyl groups. The energy difference between these conformers was found to be 1470 cm^{-1} based on fitting to the diffraction data. However, this result is in conflict with theoretical data^{8,10–12} predicting the energy difference for the two lowest conformers to be $530\text{--}880 \text{ cm}^{-1}$. Moreover, it has been shown by Godfrey *et al.*¹¹ from *ab initio* calculations and a microwave study that the two experimentally observed glycolic acid species need not necessarily be the two of lowest energy. The authors in Ref. 11 have assigned the second conformer detected by microwave spectroscopy as the *trans*-glycolic acid with relative energy of $\sim 1200 \text{ cm}^{-1}$. This conformer was structurally quite similar to the *trans*-glycolic acid conformer detected earlier in an infrared matrix isolation study.⁹

The symmetries and relative stabilities of the conformers of glycolic acid adopted in this work are based on the detailed *ab initio* calculations of Godfrey *et al.*¹¹ who tested some of their predictions experimentally. The lowest energy

conformer identified by *ab initio* calculation¹¹ was found to be the C_1 conformer which is a slightly twisted version of the C_s from detected by microwave spectroscopy.^{4–6} A conformer of C_1 symmetry exists in two enantiomeric forms, and there is a small barrier (1.5 cm^{-1}) between this conformer and its mirror image where the saddle point is of C_s symmetry. The ground vibrational state energy in these symmetric double wells may be greater than the height of the saddle point, in which case the effective structure of the observed conformer would closely match the C_s conformer of the saddle point. It should be noted that the decision between assigning C_s or C_1 symmetry is of great importance for the calculation of the thermodynamic functions. For C_1 symmetry the term $R \ln 2$ must be added to both the entropy and Gibbs energy function because two optically isomeric forms are present. C_1 symmetry was accepted in this work as the point group of the lowest-energy form of glycolic acid.

The product of the principal moments of inertia for the most stable conformer of glycolic acid was calculated using the rotational constants determined from microwave spectrum investigation.⁴ Structural parameters given above were determined from the microwave spectra of normal and isotopically substituted species of glycolic acid.⁵ In general, these geometric molecular parameters are close to those determined from electron diffraction analysis¹³ and *ab initio* calculations.^{7,8,10}

According to the *ab initio* calculation of Godfrey *et al.*¹¹ the C_1 conformer, a twisted version of *cis*-form with hydrogen bonding between the hydroxyl groups, is expected to be the second lowest conformer with relative energy of 693 cm^{-1} . In this work, the energy profile between two lowest-energy C_1 conformers¹¹ was approximated by a potential energy function for internal rotation around the C—C bond,

$$V(\varphi) = \frac{1}{2} \sum_{n=1}^8 V_n (1 - \cos n\varphi),$$

where φ is the $\text{O}_4\text{—C}_1\text{—C}_2\text{—C}_5$ torsional angle. The eight coefficients (V_n) in the expansion for this moderately complex potential energy function were determined using data from the *ab initio* calculation by Godfrey *et al.*¹¹ The flat minimum at $\varphi = 0^\circ$ corresponds to the lowest-energy C_1 conformer of glycolic acid. Because of the small barrier height between its enantiomeric forms (1.5 cm^{-1}), they are not represented by the above potential and their contribution was taken into account by adding the $R \ln 2$ to the entropy and Gibbs energy function. The barrier of 1658 cm^{-1} at $\varphi = 110^\circ$ separates the lowest-energy conformer from the second C_1 stable conformer with an energy minimum of 693 cm^{-1} at $\varphi = 155^\circ$. There is a barrier of 338 cm^{-1} at $\varphi = 180^\circ$ between this conformer and its mirror image. The value of the reduced moment of inertia I_r was calculated using the molecular structural parameters of Blom and Bauder.⁵ The next most abundant C_1 conformers with relative energies of ~ 1200 and 1300 cm^{-1} (Godfrey *et al.*¹¹) were taken into account in this work ignoring their internal

rotation and adopting their molecular constants to be the same as those of the basic conformer. The conformers with energies of about 2000 cm^{-1} and higher were ignored because of their negligible contribution to the thermodynamic functions.

Hollenstein *et al.*¹⁴ have measured the infrared spectra of 11 isotopic modifications of glycolic acid isolated in an argon matrix and have evaluated the transferable valence force field that reproduced the observed frequencies and isotopic shifts very satisfactorily. The vibrational assignment by Hollenstein *et al.*¹⁴ for C_s symmetry is adopted in this work. Fundamentals of glycolic acid calculated by an *ab initio* method¹⁰ are close to experimental values except for the low-frequency torsional mode.

The uncertainties in the calculated thermodynamic functions (Table 4) may reach $(3\text{--}5)\text{ J K}^{-1}\text{ mol}^{-1}$ for $C_p^\circ(T)$ and $(5\text{--}10)\text{ J K}^{-1}\text{ mol}^{-1}$ for $S^\circ(T)$. They are essentially due to the approximate treatment of internal rotation in glycolic acid.

Ideal gas thermodynamic properties of glycolic acid have not been reported previously.

5.3. References

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TABLE 4. Ideal gas thermodynamic properties of glycolic acid $C_2H_4O_3(g)$ at the standard state pressure, $p^\circ = 0.1 \text{ MPa}$ ($T_r = 298.15 \text{ K}$)

| T (K) | C_p° ($\text{J K}^{-1} \text{ mol}^{-1}$) | S° ($\text{J K}^{-1} \text{ mol}^{-1}$) | $-(G^\circ - H^\circ(T_r))/T$ ($\text{J K}^{-1} \text{ mol}^{-1}$) | $H^\circ - H^\circ(T_r)$ (kJ mol^{-1}) | $\Delta_f H^\circ$ (kJ mol^{-1}) | $\Delta_f G^\circ$ (kJ mol^{-1}) | $\log K_f^\circ$ |
|------------|---|---|---|--|--|--|------------------|
| 0 | 0.000 | 0.000 | ∞ | -17.007 | -567.940 | -567.940 | ∞ |
| 25 | 37.880 | 194.730 | 839.900 | -16.129 | -569.214 | -566.334 | 1183.274 |
| 50 | 39.654 | 221.605 | 524.756 | -15.158 | -571.011 | -562.818 | 587.964 |
| 75 | 42.201 | 238.111 | 426.614 | -14.138 | -572.893 | -558.300 | 388.829 |
| 100 | 46.011 | 250.752 | 381.121 | -13.037 | -574.448 | -553.193 | 288.955 |
| 150 | 55.327 | 271.119 | 341.183 | -10.510 | -577.040 | -541.987 | 188.734 |
| 180 | 61.435 | 281.742 | 330.402 | -8.759 | -578.430 | -534.844 | 155.206 |
| 190 | 63.524 | 285.120 | 327.931 | -8.134 | -578.873 | -532.410 | 146.368 |
| 200 | 65.636 | 288.431 | 325.873 | -7.488 | -579.308 | -529.953 | 138.408 |
| 210 | 67.769 | 291.685 | 324.168 | -6.821 | -579.732 | -527.475 | 131.200 |
| 220 | 69.920 | 294.888 | 322.764 | -6.133 | -580.147 | -524.977 | 124.644 |
| 230 | 72.089 | 298.043 | 321.621 | -5.423 | -580.551 | -522.460 | 118.653 |
| 240 | 74.272 | 301.158 | 320.704 | -4.691 | -580.945 | -519.926 | 113.157 |
| 250 | 76.467 | 304.234 | 319.983 | -3.937 | -581.327 | -517.376 | 108.098 |
| 260 | 78.669 | 307.276 | 319.436 | -3.162 | -581.698 | -514.811 | 103.425 |
| 270 | 80.876 | 310.286 | 319.042 | -2.364 | -582.057 | -512.231 | 99.096 |
| 280 | 83.082 | 313.267 | 318.782 | -1.544 | -582.403 | -509.639 | 95.073 |
| 290 | 85.285 | 316.221 | 318.643 | -0.702 | -582.737 | -507.034 | 91.325 |
| 298.15 | 87.074 | 318.610 | 318.610 | 0.000 | -583.000 | -504.902 | 88.456 |
| 300 | 87.480 | 319.150 | 318.611 | 0.161 | -583.059 | -504.418 | 87.826 |
| 350 | 98.188 | 333.446 | 319.716 | 4.806 | -584.470 | -491.195 | 73.306 |
| 400 | 108.146 | 347.217 | 322.298 | 9.968 | -585.568 | -477.792 | 62.392 |
| 450 | 117.092 | 360.482 | 325.809 | 15.603 | -586.386 | -464.267 | 53.890 |
| 500 | 124.943 | 373.235 | 329.918 | 21.658 | -586.965 | -450.666 | 47.080 |
| 600 | 137.622 | 397.192 | 339.162 | 34.818 | -587.557 | -423.343 | 36.855 |
| 700 | 147.128 | 419.154 | 349.044 | 49.077 | -587.601 | -395.964 | 29.547 |
| 800 | 154.453 | 439.298 | 359.085 | 64.171 | -587.259 | -368.606 | 24.067 |
| 900 | 160.307 | 457.840 | 369.041 | 79.919 | -586.638 | -341.310 | 19.809 |
| 1000 | 165.136 | 474.988 | 378.789 | 96.198 | -585.808 | -314.095 | 16.406 |
| 1100 | 169.214 | 490.923 | 388.267 | 112.921 | -584.815 | -286.970 | 13.627 |
| 1200 | 172.712 | 505.800 | 397.449 | 130.022 | -583.696 | -259.942 | 11.315 |
| 1300 | 175.742 | 519.747 | 406.326 | 147.448 | -582.482 | -233.008 | 9.362 |
| 1400 | 178.385 | 532.870 | 414.900 | 165.157 | -581.200 | -206.174 | 7.692 |
| 1500 | 180.703 | 545.258 | 423.182 | 183.114 | -579.868 | -179.431 | 6.248 |
| 1600 | 182.744 | 556.987 | 431.181 | 201.289 | -578.508 | -152.761 | 4.987 |
| 1700 | 184.547 | 568.121 | 438.912 | 219.655 | -577.132 | -126.198 | 3.878 |
| 1800 | 186.144 | 578.716 | 446.387 | 238.191 | -575.750 | -99.707 | 2.893 |
| 1900 | 187.564 | 588.818 | 453.620 | 256.878 | -574.375 | -73.299 | 2.015 |
| 2000 | 188.830 | 598.472 | 460.623 | 275.699 | -573.016 | -46.960 | 1.226 |

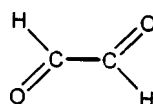
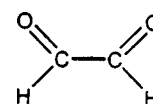
6. Glyoxal, O=CH—CH=O

Glyoxal (C₂H₂O₂)

Ideal gas

 $M_r = 58.0366$ $S^\circ(298.15\text{ K}) = 272.5 \pm 1.0 \text{ J K}^{-1} \text{ mol}^{-1}$ $\Delta_f H^\circ(0\text{ K}) = -206.4 \pm 0.5 \text{ kJ mol}^{-1}$
 $\Delta_f H^\circ(298.15\text{ K}) = -212.0 \pm 0.8 \text{ kJ mol}^{-1}$

Molecular constants

trans-Glyoxal*cis*-Glyoxal

Point group:

 C_{2h} C_{2v} Symmetry number, σ :

2

2

Ground electronic state:

 \tilde{X}^1A_g \tilde{X}^1A_1 Quantum weight, g_x :

1

1

Relative energy, ϵ_X , cm^{-1} :

0

1555

Vibrational frequencies, ν_i , and degeneracies, g_i

| | Symmetry | ν_i , cm^{-1} | g_i | Symmetry | ν_i , cm^{-1} | g_i |
|------------|----------|----------------------------|-------|----------|----------------------------|-------|
| ν_1 | A_g | 2843.27 | 1 | A_1 | 2841 | 1 |
| ν_2 | | 1744.12 | 1 | | 1746 | 1 |
| ν_3 | | 1352.60 | 1 | | 1369 | 1 |
| ν_4 | | 1065.81 | 1 | | 827 | 1 |
| ν_5 | | 550.53 | 1 | | 284.5 | 1 |
| ν_6 | A_u | 801.36 | 1 | A_2 | 1050 | 1 |
| ν_7 | | ... ^a | 1 | | ... ^a | 1 |
| ν_8 | B_g | 1047.81 | 1 | B_1 | 750 | 1 |
| ν_9 | B_u | 2835.07 | 1 | B_2 | 2810 | 1 |
| ν_{10} | | 1732.10 | 1 | | 1761 | 1 |
| ν_{11} | | 1312.38 | 1 | | 1360 | 1 |
| ν_{12} | | 338.55 | 1 | | 825 | 1 |

Geometry:

| | | |
|-----------------------|-------------------------------|---------------------|
| $r(\text{C—C})$ | $1.525 \pm 0.003 \text{ \AA}$ | 1.522 \AA |
| $r(\text{C=O})$ | $1.207 \pm 0.002 \text{ \AA}$ | 1.208 \AA |
| $r(\text{C—H})$ | $1.116 \pm 0.008 \text{ \AA}$ | 1.109 \AA |
| $\angle \text{C—C=O}$ | $121.2 \pm 0.2^\circ$ | 123.1° |
| $\angle \text{O=C—H}$ | $126.6 \pm 1.7^\circ$ | 122.8° |

Rotational constants, cm^{-1} :

| | | |
|-------|-----------|-----------|
| A_0 | 1.844 319 | 0.891 062 |
| B_0 | 0.160 021 | 0.206 503 |
| C_0 | 0.147 351 | 0.167 864 |

Product of moments of inertia,

 $I_A I_B I_C \times 10^{117} \text{ g}^3 \text{ cm}^6$:

504.42

710.17

Reduced moment of inertia for CHO

top, $I_r \times 10^{39} \text{ g cm}^2$:

0.8199

0.4807

Symmetry number for CHO top, σ_m :

1

1

^aInstead of torsional mode ν_7 (126.7 cm^{-1} for *trans*- and 89.6 cm^{-1} for *cis*-glyoxal), the contributions due to the internal rotation were calculated from the potential $V(\varphi) = \frac{1}{2} \sum_{n=1}^6 V_n (1 - \cos n\varphi)$, where φ is the torsional angle, $V_1 = 1587.6$, $V_2 = 1139.5$, $V_3 = -59.0$, $V_4 = -110.9$, $V_5 = 40.0$, and $V_6 = 0.0$ (in cm^{-1}). For the internal rotational constant B , the Fourier expansion coefficients were used $B = B_0 + \sum_{n=1}^6 B_n \cos n\varphi$ where $B_0 = 4.213$, $B_1 = -1.117$, $B_2 = 0.421$, $B_3 = -0.126$, $B_4 = 0.040$, $B_5 = -0.015$, and $B_6 = 0$ (in cm^{-1}).

6.1. Enthalpy of Formation

The value of $\Delta_f H^\circ$ (298.15 K) for glyoxal accepted in this work was determined by Fletcher and Pilcher¹ from measurements of heat of combustion by flame calorimetry. A lower value of $-224.3 \text{ kJ mol}^{-1}$ was obtained by Curtiss *et al.*² from *ab initio* calculation.

6.2. Heat Capacity and Entropy

It has been shown both spectroscopically^{3–18} and theoretically^{19–35} that glyoxal ($\text{O}=\text{CH}-\text{CH}=\text{O}$) undergoes rotational isomerization and exists in two planar, *trans* and *cis*, forms. The structure of the more stable *trans* conformer has been determined by a combination of electron diffraction and rotational data.³⁶ As *trans*-glyoxal has no dipole moment, there is no microwave spectrum, and accurate rotational constants for the ground state can only be obtained by rotational analysis of infrared or electronic absorption bands. The values obtained for the rotational constants of *trans*-glyoxal^{4,5,13,14,37} are in good agreement with each other. Using the rotational constants for the five isotopic species, Birss *et al.*¹¹ have evaluated the structural parameters for *trans*-glyoxal. These parameters are in excellent agreement with electron diffraction results.³⁶ In general, these experimental geometries are in good agreement with theoretical results.^{24,26,28,33–35,38–43} The product of the principal moments of inertia for the planar structure of *trans*-glyoxal of C_{2h} symmetry was calculated here using the rotational constants determined from a high-resolution Fourier-transform study of *trans*-glyoxal.¹⁵ Structural parameters of *trans*-glyoxal given above are those obtained from an electron diffraction investigation.³⁶ These parameters reproduce the spectroscopic moments of inertia with in an accuracy of 0.2%–0.3%.

Rotational constants for the ground state of *cis*-glyoxal were obtained from rotational analysis of the 0–0 band^{6,8} and by microwave spectroscopy.^{7,10,12,17,18} The product of the principal moments of inertia for the planar structure of *cis*-glyoxal of C_{2v} symmetry was calculated in this work using the rotational constants determined from microwave studies.^{17,18} Sets of possible structural parameters of *cis*-glyoxal were evaluated using the rotational constants for its isotopomers combined with assumptions of the values of some parameters.^{7,8,10,28,44} The geometry of *cis*-glyoxal was also calculated by *ab initio*^{22,24,26–29,31–33,35,38,39,41,42} and molecular mechanics⁴³ methods. Structural parameters of *cis*-glyoxal shown above were proposed by Tyulin *et al.*⁴⁴ from analysis of microwave data.¹⁷

Vibrational spectra of glyoxal were studied in the gas phase^{4–8,10,13,15,16,45–55} and with matrix isolation techniques.^{56–58} The experimental assignments have been confirmed by normal coordinate analyses^{44,59–64} and theoretical calculations.^{26,28,29,31,34,35,40,43,65–70} The adopted values for vibrational frequencies of *trans*-glyoxal were taken from the investigation of dispersed fluorescence spectra⁵³ ($\nu_1 - \nu_5, \nu_8$), high-resolution infrared spectra⁴⁸ ($\nu_6, \nu_7, \nu_{11}, \nu_{12}$), a rotational fine structure of the C–H

stretching band⁵ (ν_9) and a high-resolution infrared Fourier-transform study¹⁵ (ν_{10}). The uncertainties in these values are within 0.5 cm^{-1} . The adopted values of $\nu_1 - \nu_7$ and ν_{10} fundamental frequencies of *cis*-glyoxal are those observed from gas and Ar-matrix spectra.^{13,18,55,58} For unobserved frequencies ($\nu_8, \nu_9, \nu_{11}, \nu_{12}$), the values were selected on the basis of *ab initio* calculations,^{28,31,35,68} and their uncertainties are estimated to be $25\text{--}50 \text{ cm}^{-1}$.

The torsional potential function for glyoxal has been investigated experimentally^{3,6,7,9,13,16,18,71} and theoretically^{19–34,70,72} by many authors. An early infrared study³ suggested a very high *trans-cis* rotation barrier (V_{rot}) of 4810 cm^{-1} based on a torsional frequency of 128 cm^{-1} . Currie and Ramsay⁶ provided the first estimate for *cis-trans* enthalpy difference, $\Delta H = (1125 \pm 100) \text{ cm}^{-1}$, from the temperature dependence of *cis* and *trans* absorption bands in the visible spectrum. Durig *et al.*⁹ obtained the torsional potential function for glyoxal that fits the infrared data of the *trans* conformer and the microwave intensity data of the *cis* conformer and yields the $\Delta H = (1180 \pm 150) \text{ cm}^{-1}$ with a barrier height $V_{\text{rot}} = 1770 \text{ cm}^{-1}$. The energy difference between the *cis* and *trans* conformers has been revised upwards twice by Butz *et al.*, first to $(1350 \pm 200) \text{ cm}^{-1}$ derived from a spectroscopic temperature study¹³ and then to $(1688 \pm 100) \text{ cm}^{-1}$ by fitting spectroscopic data to a torsional potential.¹⁶ The potential function obtained by Butz *et al.*¹⁶ has a barrier to *trans-cis* rotation of 2077 cm^{-1} . Recently Hübner *et al.*¹⁸ have determined the *cis-trans* enthalpy difference, $\Delta H = (1555 \pm 48) \text{ cm}^{-1}$, from absorption intensities of glyoxal by microwave spectroscopy. This value is intermediate between the values reported by Butz *et al.*^{13,16} Based on this value for ΔH , Hübner *et al.*¹⁸ have recalculated the potential curve for internal rotation in glyoxal ($V_{\text{rot}} = 2003 \text{ cm}^{-1}$). Numerous *ab initio* calculations yield different results for ΔH depending on the basis sets and levels of calculation (1050 ,²⁰ $1320\text{--}1970$,³⁴ 1490 ,²⁹ $1500\text{--}1990$,³¹ 1504 ,⁷² ~ 1600 ,^{26,27,32} $1700\text{--}2200$,²³ ~ 2000 ,^{22,24,25,28,30} and 2422 cm^{-1}).⁷⁰ Due to this wide range they are not of help in deciding among the available experimental data.

The torsional potential function determined by Hübner *et al.*¹⁸ is accepted in this work in order to account for the internal rotation in glyoxal. The Fourier expansion coefficients for the internal rotation constant B were adopted to be the same as those used by Durig *et al.*⁹ and later by other authors.^{16,18} The values of the reduced moment of inertia, I_r , were calculated from accepted structural parameters (see above).

Experimentally observed absorption spectra of glyoxal have been identified with transitions to the excited electronic states \tilde{a}^3A_u ($T_0 = 19\,199 \text{ cm}^{-1}$) and \tilde{A}^1A_u ($T_0 = 21\,973 \text{ cm}^{-1}$).^{45,73,74} These assignments agree with other experimental^{75–79} and theoretical^{35,80,81} studies. Dykstra and Schaefer⁸² have predicted two low-lying ($\sim 15\,000 \text{ cm}^{-1}$) unobserved triplet states from *ab initio* calculation. Because of high energies of excited electronic states of glyoxal, they are not considered in this work. These electronic states would only make an appreciable contribution to the thermo-

TABLE 5. Ideal gas thermodynamic properties of glyoxal $C_2H_2O_2(g)$ at the standard state pressure, $p^\circ = 0.1$ MPa ($T_r = 298.15$ K)

| T (K) | C_p° ($J K^{-1} mol^{-1}$) | S° ($J K^{-1} mol^{-1}$) | $-(G^\circ - H^\circ(T_r))/T$ ($J K^{-1} mol^{-1}$) | $H^\circ - H^\circ(T_r)$ ($kJ mol^{-1}$) | $\Delta_f H^\circ$ ($kJ mol^{-1}$) | $\Delta_f G^\circ$ ($kJ mol^{-1}$) | $\log K_f^\circ$ |
|------------|--|--------------------------------------|--|---|---|---|------------------|
| 0 | 0.000 | 0.000 | ∞ | -13.690 | -206.432 | -206.432 | ∞ |
| 25 | 33.558 | 168.135 | 682.424 | -12.857 | -206.857 | -206.353 | 431.146 |
| 50 | 36.353 | 192.179 | 431.899 | -11.986 | -207.559 | -205.605 | 214.791 |
| 75 | 39.053 | 207.448 | 354.677 | -11.042 | -208.268 | -204.466 | 142.401 |
| 100 | 41.428 | 219.011 | 319.370 | -10.306 | -208.805 | -203.114 | 106.095 |
| 150 | 45.964 | 236.671 | 289.008 | -7.851 | -209.665 | -200.077 | 69.672 |
| 180 | 48.675 | 245.290 | 281.018 | -6.431 | -210.146 | -198.114 | 57.490 |
| 190 | 49.594 | 247.947 | 279.209 | -5.940 | -210.306 | -197.441 | 54.280 |
| 200 | 50.525 | 250.514 | 277.710 | -5.439 | -210.467 | -196.761 | 51.388 |
| 210 | 51.471 | 253.002 | 276.474 | -4.929 | -210.627 | -196.071 | 48.769 |
| 220 | 52.433 | 255.419 | 275.463 | -4.410 | -210.788 | -195.374 | 46.387 |
| 230 | 53.414 | 257.771 | 274.643 | -3.880 | -210.948 | -194.670 | 44.210 |
| 240 | 54.414 | 260.066 | 273.987 | -3.341 | -211.107 | -193.959 | 42.213 |
| 250 | 55.434 | 262.307 | 273.476 | -2.792 | -211.265 | -193.241 | 40.375 |
| 260 | 56.473 | 264.502 | 273.089 | -2.233 | -211.422 | -192.517 | 38.677 |
| 270 | 57.533 | 266.653 | 272.810 | -1.663 | -211.576 | -191.787 | 37.103 |
| 280 | 58.611 | 268.765 | 272.628 | -1.082 | -211.729 | -191.051 | 35.641 |
| 290 | 59.707 | 270.841 | 272.531 | -0.490 | -211.880 | -190.310 | 34.278 |
| 298.15 | 60.612 | 272.508 | 272.508 | 0.000 | -212.000 | -189.702 | 33.235 |
| 300 | 60.819 | 272.883 | 272.509 | 0.112 | -212.027 | -189.563 | 33.005 |
| 350 | 66.557 | 282.686 | 273.270 | 3.296 | -212.711 | -185.765 | 27.723 |
| 400 | 72.398 | 291.956 | 275.031 | 6.770 | -213.294 | -181.874 | 23.750 |
| 450 | 78.086 | 300.814 | 277.407 | 10.533 | -213.765 | -177.917 | 20.652 |
| 500 | 83.427 | 309.321 | 280.176 | 14.572 | -214.126 | -173.913 | 18.168 |
| 600 | 92.680 | 325.380 | 286.389 | 23.395 | -214.547 | -165.827 | 14.436 |
| 700 | 99.929 | 340.236 | 293.034 | 33.041 | -214.638 | -157.696 | 11.767 |
| 800 | 105.446 | 353.957 | 299.803 | 43.323 | -214.488 | -149.568 | 9.766 |
| 900 | 109.622 | 366.629 | 306.534 | 54.086 | -214.174 | -141.471 | 8.211 |
| 1000 | 112.807 | 378.351 | 313.137 | 65.214 | -213.761 | -133.416 | 6.969 |
| 1100 | 115.270 | 389.223 | 319.566 | 76.623 | -213.288 | -125.404 | 5.955 |
| 1200 | 117.204 | 399.339 | 325.797 | 88.251 | -212.789 | -117.436 | 5.112 |
| 1300 | 118.749 | 408.784 | 331.821 | 100.051 | -212.289 | -109.508 | 4.400 |
| 1400 | 120.001 | 417.631 | 337.638 | 111.991 | -211.806 | -101.621 | 3.791 |
| 1500 | 121.030 | 425.947 | 343.251 | 124.044 | -211.349 | -93.765 | 3.265 |
| 1600 | 121.887 | 433.786 | 348.667 | 136.191 | -210.932 | -85.921 | 2.805 |
| 1700 | 122.609 | 441.198 | 353.894 | 148.417 | -210.556 | -78.122 | 2.400 |
| 1800 | 123.223 | 448.224 | 358.941 | 160.709 | -210.227 | -70.337 | 2.041 |
| 1900 | 123.749 | 454.900 | 363.817 | 173.058 | -209.948 | -62.575 | 1.720 |
| 2000 | 124.204 | 461.260 | 368.531 | 185.457 | -209.719 | -54.823 | 1.432 |

dynamic functions at temperatures above 3000 K.

The uncertainties in the calculated thermodynamic functions (Table 5) are estimated to be (0.5–3.0) $J K^{-1} mol^{-1}$ for $C_p^\circ(T)$ and (1.0–3.0) $J K^{-1} mol^{-1}$ for $S^\circ(T)$.

The thermodynamic functions of glyoxal as an equilibrium mixture of *trans* and *cis* conformers were calculated by Compton.⁸³ The discrepancies between $C_p^\circ(T)$ and $S^\circ(T)$ values calculated in this work and in Compton⁸³ increase with temperature and reach 3.1 and 2.2 $J K^{-1} mol^{-1}$, respectively, at 1000 K. These discrepancies are due to different

molecular constants used in the calculations and the noninclusion of internal rotation in glyoxal by Compton.⁸³ Values of $H^\circ(T) - H^\circ(0)$ given by Compton appear to be in error. Hollenstein *et al.*⁸⁴ derived the thermodynamic functions of glyoxal using the semiclassical approximation for the sum-over-states calculation for nonrigid molecules. The difference between their $C_p^\circ(T)$ and $S^\circ(T)$ values and those calculated in this work range from 0.1 to 8.8 $J K^{-1} mol^{-1}$ depending on the level of approximation used. The discrepancies with the results of statistical calculations by Natarajan

*et al.*⁸⁵ amount to (50–60) J K⁻¹ mol⁻¹ and cannot be due to differences in the molecular constants used.

6.3. References

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7. Cyanooxomethyl Radical, $\dot{\text{O}}\text{CCN}$ Cyanooxomethyl radical (C_2NO)

Ideal gas

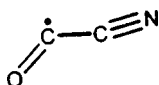
 $M_r = 54.0281$ $S^\circ(298.15\text{ K}) = 278.2 \pm 1.5\text{ J K}^{-1}\text{ mol}^{-1}$ $\Delta_f H^\circ(0\text{ K}) = 207.2 \pm 10.0\text{ kJ mol}^{-1}$ $\Delta_f H^\circ(298.15\text{ K}) = 210.0 \pm 10.0\text{ kJ mol}^{-1}$

Molecular Constants

Point group: C_s Ground electronic state: \tilde{X}^2A' Excited electronic state: \tilde{A}^2A'' Energy: $\epsilon_X = 0\text{ cm}^{-1}$ Energy: $\epsilon_A = 15\,500\text{ cm}^{-1}$ Symmetry number: $\sigma = 1$ Quantum weight: $g_X = 2$ Quantum weight: $g_A = 2$ Vibrational frequencies, ν_i , and degeneracies, g_i

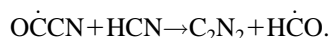
| Symmetry | ν_i, cm^{-1} | g_i | Symmetry | ν_i, cm^{-1} | g_i |
|----------|-------------------------|-------|----------|-------------------------|-------|
| A' | ν_1 2249 | 1 | | ν_4 488 | |
| | ν_2 1703 | 1 | | ν_5 174 | 1 |
| | ν_3 909 | 1 | A'' | ν_6 233 | 1 |

Geometry

 $r(\text{C}=\text{O}) = 1.19 \pm 0.02\text{ \AA}$ $r(\text{C}-\text{C}) = 1.47 \pm 0.02\text{ \AA}$ $r(\text{C}\equiv\text{N}) = 1.16 \pm 0.02\text{ \AA}$ $\angle \text{O}=\text{C}-\text{C} = 128 \pm 3^\circ$ $\angle \text{C}-\text{C}\equiv\text{N} = 170 \pm 5^\circ$ Product of moments of inertia: $I_A I_B I_C = 183 \times 10^{-117}\text{ g}^3\text{ cm}^6$.

7.1. Enthalpy of Formation

No experimental data are available for enthalpy of formation of the cyanooxomethyl radical. The value accepted in this work was estimated by Francisco and Liu¹ by using the isodesmic reaction approach with the following reaction:



The relationship between the enthalpy of formation of $\dot{\text{O}}\text{CCN}$ and the enthalpy change of the above reaction is

$$\begin{aligned} \Delta_f H^\circ(\dot{\text{O}}\text{CCN}) &= \Delta_f H^\circ(\text{C}_2\text{N}_2) + \Delta_f H^\circ(\dot{\text{H}}\text{C}\ddot{\text{O}}) \\ &\quad - \Delta_f H^\circ(\text{HCN}) - \Delta_r H^\circ. \end{aligned}$$

Based on known experimental enthalpies of formation of HCN, C_2N_2 , $\dot{\text{H}}\text{C}\ddot{\text{O}}$, and the $\Delta_r H^\circ$ value predicted by *ab initio* calculations, Francisco and Liu¹ have estimated the enthalpy of formation for $\dot{\text{O}}\text{CCN}$ to be 207.5–210.0 kJ mol^{-1} .

7.2. Heat Capacity and Entropy

There are no experimental data on structure and vibrational spectra of $\dot{\text{O}}\text{CCN}$. According to *ab initio* calculations,^{1,2} a bent structure of C_s symmetry is adopted in this work for $\dot{\text{O}}\text{CCN}$ in the ground electronic state \tilde{X}^2A' . The product of the principal moments of inertia was calculated using the structural parameters shown above. These

values are based on the results of *ab initio* calculations^{1,2} and comparison with structural parameters of $\text{CO}(\text{CN})_2$, COX_2 , and XCO ($X=\text{F}, \text{Cl}$) molecules.

Vibrational frequencies of $\dot{\text{O}}\text{CCN}$ were calculated in this work using the following force constants:

$$f_{\text{C}=\text{O}} = 11.413\text{ m dyn/\AA}, \quad f_{\text{C}-\text{C}\equiv\text{N}} = 0.110\text{ m dyn/\AA},$$

$$f_{\text{C}-\text{C}} = 7.225\text{ m dyn/\AA}, \quad f_{\text{CC},\text{CO}} = 0.513\text{ m dyn/\AA},$$

$$f_{\text{C}\equiv\text{N}} = 13.674\text{ m dyn/\AA}, \quad f_{\text{CC},\text{CN}} = -1.543\text{ m dyn/\AA},$$

$$f_{\text{C}-\text{C}=\text{O}} = 0.750\text{ m dyn/\AA}, \quad f_{\text{CC},\text{CCO}} = 0.585\text{ m dyn/\AA}.$$

These constants were transferred from the $\text{CO}(\text{CN})_2$ molecule except for $f_{\text{C}-\text{C}=\text{O}}$ and $f_{\text{C}-\text{C}\equiv\text{N}}$ whose values were reduced by comparison with the bending force constants in COX_2 and XCO ($X=\text{F}, \text{Cl}, \text{Br}$). Normal coordinate analysis for $\text{CO}(\text{CN})_2$ was carried using the vibrational assignment of Miller *et al.*³ The uncertainties in the calculated frequencies of $\dot{\text{O}}\text{CCN}$ can reach 50 cm^{-1} . These frequencies agree overall with those resulting from the *ab initio* calculation.¹

The first excited electronic state \tilde{A}^2A'' of $\dot{\text{O}}\text{CCN}$ was predicted at an energy of $\sim 15\,500\text{ cm}^{-1}$ by an *ab initio* calculation¹ and was included in the calculations in this work. Structural parameters and vibrational frequencies for the \tilde{A}^2A'' state were accepted as identical to those of the ground state.

The uncertainties in the calculated thermodynamic functions (Table 6) are estimated to be (1.5–2.0) $\text{J K}^{-1} \text{mol}^{-1}$ for $C_p^\circ(T)$ and (1.5–3.0) $\text{J K}^{-1} \text{mol}^{-1}$ for $S^\circ(T)$.

Ideal gas thermodynamic properties of the cyanooxomethyl radical have not been reported previously.

7.3. References

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TABLE 6. Ideal gas thermodynamic properties of the cyanooxomethyl radical $\text{C}_2\text{NO}(\text{g})$ at the standard state pressure, $p^\circ = 0.1 \text{ MPa}$ ($T_r = 298.15 \text{ K}$)

| T (K) | C_p° ($\text{J K}^{-1} \text{mol}^{-1}$) | S° ($\text{J K}^{-1} \text{mol}^{-1}$) | $-(G^\circ - H^\circ(T_r))/T$ ($\text{J K}^{-1} \text{mol}^{-1}$) | $H^\circ - H^\circ(T_r)$ (kJ mol^{-1}) | $\Delta_f H^\circ$ (kJ mol^{-1}) | $\Delta_f G^\circ$ (kJ mol^{-1}) | $\log K_f^\circ$ |
|------------|--|--|--|--|--|--|------------------|
| 0 | 0.000 | 0.000 | ∞ | -13.594 | 207.190 | 207.190 | ∞ |
| 25 | 33.297 | 174.512 | 685.016 | -12.763 | 207.298 | 206.094 | -430.604 |
| 50 | 35.132 | 197.975 | 436.251 | -11.914 | 207.405 | 204.851 | -214.003 |
| 75 | 38.803 | 212.901 | 359.439 | -10.990 | 207.563 | 203.542 | -141.757 |
| 100 | 42.195 | 224.543 | 324.309 | -9.977 | 207.781 | 202.171 | -105.602 |
| 150 | 47.280 | 242.677 | 294.223 | -7.732 | 208.327 | 199.251 | -69.384 |
| 180 | 49.604 | 251.509 | 286.384 | -6.278 | 208.683 | 197.403 | -57.284 |
| 190 | 50.296 | 254.210 | 284.620 | -5.778 | 208.802 | 196.773 | -54.096 |
| 200 | 50.955 | 256.807 | 283.165 | -5.272 | 208.921 | 196.137 | -51.225 |
| 210 | 51.582 | 259.308 | 281.970 | -4.759 | 209.039 | 195.495 | -48.626 |
| 220 | 52.182 | 261.721 | 280.995 | -4.240 | 209.156 | 194.847 | -46.262 |
| 230 | 52.757 | 264.054 | 280.208 | -3.715 | 209.271 | 194.194 | -44.102 |
| 240 | 53.309 | 266.311 | 279.582 | -3.185 | 209.384 | 193.536 | -42.121 |
| 250 | 53.839 | 268.498 | 279.095 | -2.649 | 209.496 | 192.873 | -40.298 |
| 260 | 54.350 | 270.620 | 278.729 | -2.108 | 209.605 | 192.206 | -38.614 |
| 270 | 54.842 | 272.680 | 278.467 | -1.562 | 209.712 | 191.535 | -37.054 |
| 280 | 55.318 | 274.683 | 278.296 | -1.012 | 209.817 | 190.860 | -35.605 |
| 290 | 55.779 | 276.632 | 278.205 | -0.456 | 209.919 | 190.181 | -34.255 |
| 298.15 | 56.145 | 278.183 | 278.183 | 0.000 | 210.000 | 189.625 | -33.221 |
| 300 | 56.226 | 278.531 | 278.184 | 0.104 | 210.018 | 189.499 | -32.994 |
| 350 | 58.286 | 287.356 | 278.877 | 2.968 | 210.473 | 186.041 | -27.765 |
| 400 | 60.126 | 295.262 | 280.439 | 5.929 | 210.852 | 182.524 | -23.835 |
| 450 | 61.811 | 302.442 | 282.491 | 8.978 | 211.153 | 178.965 | -20.773 |
| 500 | 63.373 | 309.036 | 284.820 | 12.108 | 211.379 | 175.376 | -18.321 |
| 600 | 66.186 | 320.845 | 289.863 | 18.589 | 211.634 | 168.146 | -14.638 |
| 700 | 68.609 | 331.235 | 295.046 | 25.332 | 211.683 | 160.892 | -12.006 |
| 800 | 70.666 | 340.534 | 300.161 | 32.299 | 211.585 | 153.643 | -10.032 |
| 900 | 72.395 | 348.960 | 305.122 | 39.454 | 211.379 | 146.410 | -8.497 |
| 1000 | 73.839 | 356.665 | 309.897 | 46.768 | 211.093 | 139.204 | -7.271 |
| 1100 | 75.044 | 363.761 | 314.475 | 54.214 | 210.748 | 132.033 | -6.270 |
| 1200 | 76.051 | 370.335 | 318.860 | 61.770 | 210.353 | 124.893 | -5.436 |
| 1300 | 76.896 | 376.457 | 323.057 | 69.419 | 209.918 | 117.792 | -4.733 |
| 1400 | 77.609 | 382.182 | 327.078 | 77.145 | 209.441 | 110.720 | -4.131 |
| 1500 | 78.214 | 387.558 | 330.933 | 84.937 | 208.931 | 103.689 | -3.611 |
| 1600 | 78.731 | 392.622 | 334.632 | 92.785 | 208.384 | 96.710 | -3.157 |
| 1700 | 79.175 | 397.409 | 338.185 | 100.681 | 207.808 | 89.746 | -2.758 |
| 1800 | 79.559 | 401.946 | 341.602 | 108.618 | 207.199 | 82.823 | -2.403 |
| 1900 | 79.894 | 406.257 | 344.893 | 116.591 | 206.559 | 75.929 | -2.087 |
| 2000 | 80.189 | 410.362 | 348.064 | 124.596 | 205.890 | 69.071 | -1.804 |

8. Oxalic Acid, HO—CO—CO—OH

Oxalic acid (C₂H₂O₄)

Ideal gas

 $M_r = 90.0354$ $S^\circ(298.15\text{ K}) = 320.6 \pm 5.0 \text{ J K}^{-1} \text{ mol}^{-1}$ $\Delta_f H^\circ(0\text{ K}) = -721.2 \pm 2.0 \text{ kJ mol}^{-1}$ $\Delta_f H^\circ(298.15\text{ K}) = -731.8 \pm 2.0 \text{ kJ mol}^{-1}$

Molecular constants

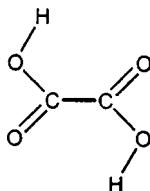
Point group: C_{2h} Symmetry number: $\sigma = 2$ Ground electronic state: \tilde{X}^1A_g Energy: $\epsilon_X = 0 \text{ cm}^{-1}$ Quantum weight: $g_X = 1$ Vibrational frequencies, ν_i , and degeneracies, g_i

| Symmetry | ν_i, cm^{-1} | g_i | Symmetry | ν_i, cm^{-1} | g_i | | | |
|----------|-------------------------|-------|----------|-------------------------|------------|----------------|------|---|
| A_g | ν_1 | 3484 | 1 | B_g | ν_{10} | - ^a | 1 | |
| | ν_2 | 1800 | 1 | | ν_{11} | 851 | 1 | |
| | ν_3 | 1423 | 1 | | ν_{12} | 563 | 1 | |
| | ν_4 | 1195 | 1 | | B_u | ν_{13} | 3484 | 1 |
| | ν_5 | 815 | 1 | | | ν_{14} | 1826 | 1 |
| | ν_6 | 608 | 1 | | | ν_{15} | 1278 | 1 |
| | ν_7 | 405 | 1 | | | ν_{16} | 1127 | 1 |
| A_u | ν_8 | 666 | 1 | ν_{17} | | 651 | 1 | |
| | ν_9 | 460 | 1 | ν_{18} | 264 | 1 | | |

^aInstead of torsional mode $\nu_{10} \sim 90 \text{ cm}^{-1}$, the contributions due to the internal rotation about C—C bond were calculated from the potential $V(\varphi) = \frac{1}{2}V_1(1 - \cos \varphi)$, where φ is the torsional angle and $V_1 = 700 \text{ cm}^{-1}$.

COOH top: Reduced moment of inertia, $I_r = 3.6454 \times 10^{-39} \text{ g cm}^2$, Symmetry number, $\sigma_m = 1$

Geometry



| | |
|-----------------------|-------------------------------|
| $r(\text{C—C})$ | $1.548 \pm 0.004 \text{ \AA}$ |
| $r(\text{C=O})$ | $1.208 \pm 0.001 \text{ \AA}$ |
| $r(\text{C—O})$ | $1.339 \pm 0.002 \text{ \AA}$ |
| $r(\text{O—H})$ | $1.056 \pm 0.014 \text{ \AA}$ |
| $\angle \text{C—C—O}$ | $123.1 \pm 0.9^\circ$ |
| $\angle \text{O=C—O}$ | $125.0 \pm 0.2^\circ$ |
| $\angle \text{C—O—H}$ | $104.4 \pm 2.3^\circ$ |

Product of moments of inertia: $I_A I_B I_C = 11\,950 \times 10^{-117} \text{ g}^3 \text{ cm}^6$.

8.1. Enthalpy of Formation

Wilhoit and Shiao¹ have measured the enthalpy of combustion of solid oxalic acid in a rotating platinum bomb calorimeter and have calculated $\Delta_f H^\circ(\text{C}_2\text{H}_2\text{O}_4, \text{cr}, 298.15\text{ K}) = -(829.9 \pm 1.0) \text{ kJ mol}^{-1}$. The enthalpy of formation of gaseous oxalic acid is calculated from this value by adding the enthalpy of sublimation, $\Delta_{\text{sub}} H^\circ = 98.1 \text{ kJ mol}^{-1}$, obtained from vapor pressure measurements.²

8.2. Heat Capacity and Entropy

Nahlovska *et al.*³ carried out an electron diffraction and an IR study of oxalic acid indicating that the structure was in a planar *trans* conformation (C_{2h} symmetry) in which the

hydrogen atom of one COOH group participated in intramolecular hydrogen bonding with the carbonyl oxygen of the other COOH group (“hydrogen bonded” *trans* conformer). Infrared matrix-isolation spectra of oxalic acid⁴ were interpreted in terms of the same model and the tentative conclusion was made that a second conformer of oxalic acid exists in the vapor phase. This conformer was suggested to be a *trans* form (C_{2h} symmetry) with the hydrogen atom of each COOH group oriented towards the carbonyl oxygen of the same COOH group (“free” *trans* form). Oxalic acid was determined to be in its *trans* conformation from x-ray crystallographic studies⁵ and no other conformers were observed. Studies of oxalic acid in solution⁶ favored almost free rota-

tion, which would include a range of possible conformations. *Ab initio* studies of oxalic acid⁷⁻⁹ have obtained a most stable hydrogen bonded *trans* planar structure in agreement with the results of an electron diffraction investigation.³ The hydrogen bonded or free *trans* form was found to be lowest energy depending on the basis set used in the *ab initio* calculations.¹⁰ The free *trans* conformer was reported to be most stable by an earlier *ab initio* study.¹¹

The hydrogen bonded *trans* form (C_{2h} symmetry) is accepted for oxalic acid in this work. Its product of principal moments of inertia was calculated using the structural parameters determined from the electron diffraction study.³ Following the results of a theoretical calculation by Tyrrell,⁹ it was assumed that there is a slight energy preference for this conformer and only a small barrier of about 2 kcal mol⁻¹ (700 cm⁻¹) separating it from other stable conformers. The simple approximate potential,

$$V(\varphi) = \frac{1}{2}V_1(1 - \cos \varphi),$$

where φ is the OC—CO torsional angle, was used here for the calculation of internal rotational contributions to the thermodynamic functions of oxalic acid. The value of the reduced moment of inertia for the COOH top was derived from the molecular constants shown above.

There are several studies of the vibrational spectra of the free C₂H₂O₄ molecule^{3,4,12-14} but some fundamentals are still unobserved. Stace and Oralratmanee¹⁴ have reported the infrared vapor phase measurements and the first vapor phase Raman spectra. They have proposed a new vibrational frequencies assignment based on their experimental results as well as calculated in-plane frequencies using an Urey–Bradley force field derived from formic and acetic acid. Redington and Redington⁴ have investigated the infrared spectra of oxalic acid vapor and the infrared spectra of matrix-isolated C₂H₂O₄, C₂HDO₄, and C₂D₂O₄ and have carried out a normal coordinate analysis using a general valence force field for the in-plane infrared-active modes. None of the above studies, however, gives a complete vibrational assignment for oxalic acid. De Villepin and Novak¹⁵ have developed a general valence force field for both infrared and Raman frequencies using the relatively small number of experimentally observed fundamentals. Their force field reproduces the experimental vibrational spectra of oxalic and is comparable with force fields already known for other carboxylic acids. The fundamental frequencies ν_2 – ν_5 and ν_7 adopted here were observed in the Raman spectra of gaseous oxalic acid.¹⁴ The 405 cm⁻¹ mode was subsequently re-assigned to ν_7 .⁴ The values of ν_8 , ν_9 , ν_{13} – ν_{16} , and ν_{18} are those determined from gas-phase infrared spectra.⁴ The value for ν_{17} was obtained from spectra of neon matrix-isolated oxalic acid.⁴ The uncertainties in these experimental frequencies are in the range of 5–10 cm⁻¹.

Redington and Redington⁴ have suggested tentative values for $\nu_6(A_g) = 538$ cm⁻¹, $\nu_{11}(B_g) = 590$ cm⁻¹, and $\nu_{12}(B_g) = 512$ cm⁻¹ fundamentals using possible combination bands.

The adopted values of these frequencies were taken from normal coordinate calculations.¹⁵ Calculations for the other frequencies are in good agreement with experimental results. The value of the symmetric O—H stretching frequency ν_1 was accepted to be the same as the value of antisymmetric O—H stretching (ν_{13}) and it coincides with the calculated value.¹⁵ The uncertainties in these frequencies are estimated to be 25–50 cm⁻¹.

Evidently due to its weak intensity, the C—C torsion mode $\nu_{10}(A_u)$ was not observed in a vapor phase spectra search that extended down to 35 cm⁻¹.⁴ *Ab initio* calculations^{8,10} have suggested the value of 160 cm⁻¹ for the torsional mode, however, it should be noted that all other calculated values are much greater than those from experiment. Cyvin and Alfheim¹⁶ have calculated $\nu_{10} = 90$ cm⁻¹ using force constants transferred from formic and acetic acids. The value of 110 cm⁻¹ was obtained from other normal coordinate calculation.¹⁵ The corresponding torsional mode for oxalyl fluoride (F—CO—CO—F) was observed in the gas phase at 54 cm⁻¹ and in the solid phase at 94 cm⁻¹.¹⁷ From the ratio of $\nu_{\text{tors}}(\text{C}_2\text{F}_2\text{O}_2, \text{solid})/\nu_{\text{tors}}(\text{C}_2\text{F}_2\text{O}_2, \text{gas}) = 1.74$ and the value of 138 cm⁻¹ for the torsional mode of solid oxalic acid,^{15,18} the value of ~ 80 cm⁻¹ for the torsional mode of gaseous oxalic acid is expected.

The uncertainties in the calculated thermodynamic functions (Table 7) may be as much as (5–7) JK⁻¹mol⁻¹ for $C_p^\circ(T)$ and (5–10) JK⁻¹mol⁻¹ for $S^\circ(T)$. These uncertainties arise from the uncertainties in the adopted values of the vibrational frequencies and the simple and approximate model for internal rotation.

Ideal gas thermodynamic properties of oxalic acid have not been reported previously.

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TABLE 7. Ideal gas thermodynamic properties of oxalic acid C₂H₂O₄(g) at the standard state pressure, $p^\circ = 0.1$ MPa ($T_r = 298.15$ K)

| T (K) | C_p° (J K ⁻¹ mol ⁻¹) | S° (J K ⁻¹ mol ⁻¹) | $-(G^\circ - H^\circ(T_r))/T$ (J K ⁻¹ mol ⁻¹) | $H^\circ - H^\circ(T_r)$ (kJ mol ⁻¹) | $\Delta_f H^\circ$ (kJ mol ⁻¹) | $\Delta_f G^\circ$ (kJ mol ⁻¹) | $\log K_f^\circ$ |
|------------|---|---|---|---|---|---|------------------|
| 0 | 0.000 | 0.000 | ∞ | -17.321 | -721.180 | -721.180 | ∞ |
| 25 | 40.556 | 193.337 | 848.905 | -16.389 | -722.227 | -719.030 | 1502.311 |
| 50 | 41.771 | 221.869 | 529.039 | -15.359 | -723.498 | -715.371 | 747.333 |
| 75 | 43.565 | 239.102 | 429.695 | -14.295 | -724.813 | -711.016 | 495.189 |
| 100 | 46.685 | 252.031 | 383.721 | -13.169 | -725.959 | -706.240 | 368.897 |
| 150 | 56.008 | 272.618 | 343.382 | -10.615 | -727.905 | -695.943 | 242.346 |
| 180 | 62.509 | 283.401 | 332.497 | -8.837 | -728.902 | -689.455 | 200.072 |
| 190 | 64.693 | 286.839 | 330.004 | -8.201 | -729.209 | -687.255 | 188.937 |
| 200 | 66.861 | 290.213 | 327.931 | -7.544 | -729.504 | -685.040 | 178.912 |
| 210 | 69.005 | 293.527 | 326.214 | -6.864 | -729.786 | -682.809 | 169.837 |
| 220 | 71.116 | 296.786 | 324.802 | -6.164 | -730.057 | -680.566 | 161.585 |
| 230 | 73.190 | 299.993 | 323.654 | -5.442 | -730.316 | -678.311 | 154.047 |
| 240 | 75.223 | 303.151 | 322.734 | -4.700 | -730.564 | -676.044 | 147.135 |
| 250 | 77.213 | 306.262 | 322.013 | -3.938 | -730.801 | -673.768 | 140.774 |
| 260 | 79.160 | 309.329 | 321.466 | -3.156 | -731.027 | -671.482 | 134.901 |
| 270 | 81.061 | 312.352 | 321.073 | -2.355 | -731.243 | -669.187 | 129.460 |
| 280 | 82.918 | 315.334 | 320.815 | -1.535 | -731.450 | -666.886 | 124.407 |
| 290 | 84.731 | 318.275 | 320.677 | -0.696 | -731.646 | -664.576 | 119.701 |
| 298.15 | 86.176 | 320.644 | 320.644 | 0.000 | -731.800 | -662.689 | 116.099 |
| 300 | 86.501 | 321.178 | 320.645 | 0.160 | -731.834 | -662.260 | 115.308 |
| 350 | 94.724 | 335.142 | 321.729 | 4.695 | -732.644 | -650.598 | 97.095 |
| 400 | 101.995 | 348.276 | 324.235 | 9.616 | -733.273 | -638.832 | 83.422 |
| 450 | 108.424 | 360.668 | 327.601 | 14.880 | -733.761 | -626.996 | 72.779 |
| 500 | 114.112 | 372.392 | 331.499 | 20.446 | -734.137 | -615.113 | 64.260 |
| 600 | 123.628 | 394.073 | 340.152 | 32.352 | -734.633 | -591.258 | 51.473 |
| 700 | 131.195 | 413.720 | 349.281 | 45.107 | -734.871 | -567.341 | 42.335 |
| 800 | 137.323 | 431.652 | 358.473 | 58.543 | -734.902 | -543.401 | 35.480 |
| 900 | 142.375 | 448.127 | 367.531 | 72.536 | -734.765 | -519.472 | 30.149 |
| 1000 | 146.605 | 463.353 | 376.362 | 86.991 | -734.487 | -495.566 | 25.885 |
| 1100 | 150.187 | 477.499 | 384.921 | 101.836 | -734.087 | -471.692 | 22.398 |
| 1200 | 153.247 | 490.701 | 393.192 | 117.011 | -733.590 | -447.859 | 19.495 |
| 1300 | 155.880 | 503.074 | 401.173 | 132.471 | -733.013 | -424.068 | 17.039 |
| 1400 | 158.158 | 514.711 | 408.872 | 148.175 | -732.379 | -400.327 | 14.936 |
| 1500 | 160.137 | 525.692 | 416.297 | 164.092 | -731.700 | -376.632 | 13.115 |
| 1600 | 161.866 | 536.084 | 423.462 | 180.195 | -730.994 | -352.965 | 11.523 |
| 1700 | 163.381 | 545.943 | 430.379 | 196.458 | -730.273 | -329.363 | 10.120 |
| 1800 | 164.714 | 555.320 | 437.062 | 212.865 | -729.544 | -305.796 | 8.874 |
| 1900 | 165.891 | 564.258 | 443.523 | 229.396 | -728.823 | -282.276 | 7.760 |
| 2000 | 166.934 | 572.794 | 449.775 | 246.038 | -728.113 | -258.789 | 6.759 |

9. Methyl Hydroperoxide, CH₃—O—O—H

Methyl hydroperoxide (CH₄O₂)

Ideal gas

 $M_r = 48.0414$ $S^\circ(298.15 \text{ K}) = 276.5 \pm 3.0 \text{ J K}^{-1} \text{ mol}^{-1}$ $\Delta_f H^\circ(0 \text{ K}) = -126.2 \pm 5.0 \text{ kJ mol}^{-1}$ $\Delta_f H^\circ(298.15 \text{ K}) = -139.0 \pm 5.0 \text{ kJ mol}^{-1}$

Molecular constants

Point group: C_1 Symmetry number: $\sigma = 1$ Number of optical isomers: $n = 2$ Ground electronic state: \tilde{X}^1A Energy: $\epsilon_X = 0 \text{ cm}^{-1}$ Quantum weight: $g_X = 1$ Vibrational frequencies, ν_i , and degeneracies, g_i

| Symmetry | ν_i, cm^{-1} | g_i | Symmetry | ν_i, cm^{-1} | g_i |
|----------|-------------------------|-------|----------|-------------------------|-------|
| A | ν_1 | 3604 | 1 | ν_9 | 1145 |
| | ν_2 | 2957 | 1 | ν_{10} | 1115 |
| | ν_3 | 2955 | 1 | ν_{11} | 1003 |

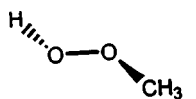
| | | | | | |
|---------|------|---|------------|------------------|---|
| ν_4 | 2861 | 1 | ν_{12} | 800 | 1 |
| ν_5 | 1509 | 1 | ν_{13} | 415 | 1 |
| ν_6 | 1453 | 1 | ν_{14} | ... ^a | 1 |
| ν_7 | 1450 | 1 | ν_{15} | ... ^b | 1 |
| ν_8 | 1348 | 1 | | | |

^aInstead of torsional mode $\nu_{14}=240\text{ cm}^{-1}$, the contributions due to the internal rotation of CH_3 group around the C—O bond were calculated from the potential $V(\varphi)=\frac{1}{2}V_3(1-\cos 3\varphi)$, where φ is the H—C—O—O torsional angle and $V_3=1120\text{ cm}^{-1}$.

^bInstead of torsional mode $\nu_{15}=149\text{ cm}^{-1}$, the contributions due to the internal rotation of OH group around the O—O bond were calculated from the potential: $V(\varphi)=V_0+V_1\cos\varphi+V_2\cos 2\varphi+V_3\cos 3\varphi$, where φ is the C—O—O—H torsional angle, $V_0=780.7$, $V_1=1111.1$, $V_2=555.6$, and $V_3=52.6$ (in cm^{-1}).

CH_3 top: Reduced moment of inertia, $I_r=0.4282\times 10^{-39}\text{ g cm}^2$, Symmetry number, $\sigma_m=3$.

OH top: Reduced moment of inertia, $I_r=0.138\times 10^{-39}\text{ g cm}^2$, Symmetry number, $\sigma_m=1$.



Geometry

| | |
|---------------------------|-----------------------------|
| $r(\text{C—H})$ | $1.090\pm 0.010\text{ \AA}$ |
| $r(\text{O—H})$ | $0.970\pm 0.005\text{ \AA}$ |
| $\angle\text{C—O—O}$ | $105.3\pm 0.5^\circ$ |
| $\angle\text{O—O—H}$ | $100.0\pm 2.0^\circ$ |
| $\angle\text{H—C—O}$ | $109.5\pm 4.0^\circ$ |
| $\varphi(\text{C—O—O—H})$ | $120\pm 5^\circ$ |
| $r(\text{C—O})$ | $1.420\pm 0.005\text{ \AA}$ |
| $r(\text{O—O})$ | $1.460\pm 0.010\text{ \AA}$ |

Rotational constants in cm^{-1} :

$$A_0=1.434\ 544\quad B_0=0.350\ 826\quad C_0=0.301\ 985.$$

Product of moments of inertia: $I_A I_B I_C=144.3\times 10^{-117}\text{ g}^3\text{ cm}^6$.

9.1. Enthalpy of Formation

Experimental data on enthalpy of formation are not available for methyl hydroperoxide. Values of $\Delta_f H^\circ(\text{CH}_3\text{OOH}, \text{g}, 298.15\text{ K})$ between -122 and -138 kJ mol^{-1} were estimated from semiempirical,¹⁻³ molecular mechanics,⁴⁻⁶ *ab initio*,⁷ and group additivity⁸⁻¹⁰ calculations. Lay *et al.*¹¹ have estimated $\Delta_f H^\circ$ (298.15 K) to be $(139.7\pm 5.0)\text{ kJ mol}^{-1}$ using the experimentally determined value of $\Delta_f H^\circ(\text{CH}_3\text{OO})$ with an average bond energy for the O—H bond in ROOH compounds. Based on the experimental values of $\Delta_f H^\circ(\text{CH}_3\text{O})$, $\Delta_f H^\circ(\text{OH})$,¹² and the bond energy of the O—O bond in CH_3OOH ,¹³ the value of $-134.3\text{ kJ mol}^{-1}$ may be obtained for the enthalpy of formation of methyl hydroperoxide. Benassi *et al.*⁵ proposed $\Delta_f H^\circ$ (298.15 K) = $-138.1\text{ kJ mol}^{-1}$ by employing theoretical and empirical approaches. Similar values of $\Delta_f H^\circ$ are predicted by the method of group equations from

$$\Delta_f H^\circ(\text{CH}_3\text{OOH})$$

$$\begin{aligned} &= \Delta_f H^\circ(\text{CH}_3\text{OOCH}_3) + \Delta_f H^\circ(\text{CH}_3\text{OH}) \\ &\quad - \Delta_f H^\circ(\text{CH}_3\text{OCH}_3) \\ &= (-125.5) + (-201.5) - (-184.1) \\ &= -142.9\text{ kJ mol}^{-1}, \end{aligned}$$

$$\Delta_f H^\circ(\text{CH}_3\text{OOH})$$

$$\begin{aligned} &= \Delta_f H^\circ(\text{CH}_3\text{CH}_2\text{OOH}) + \Delta_f H^\circ(\text{CH}_3\text{OH}) \\ &\quad - \Delta_f H^\circ(\text{CH}_3\text{CH}_2\text{OH}) \\ &= (-173.6) + (-201.5) - (-235.2) \\ &= -139.9\text{ kJ mol}^{-1}, \end{aligned}$$

(the $\Delta_f H^\circ$ values were taken from Pedley¹⁴ (CH_3OH , CH_3OCH_3 , and $\text{CH}_3\text{CH}_2\text{OH}$), Baker *et al.*¹⁵ (CH_3OOCH_3), and Lay *et al.*¹¹ ($\text{CH}_3\text{CH}_2\text{OOH}$). The average of these esti-

mates, -139 ± 5.0 kJ mol⁻¹, is accepted in this work for the enthalpy of formation of methyl hydroperoxide at 298.15 K.

9.2. Heat Capacity and Entropy

The microwave spectrum of methyl hydroperoxide, CH₃OOH, has been investigated by Tyblewski *et al.*¹⁶ However, the assignment of the spectrum was complicated because of the widespread effects of the internal rotation around the O—O bond. The experimental results provide definitive evidence that the minimum of the potential energy relative to internal rotation around the O—O bond corresponds to a *skew* conformation. The adjustment to the observed data resulted in a small *trans* barrier of 172.5 cm⁻¹. This barrier separates two equivalent potential minima associated with two enantiomeric *skew* forms. From a preliminary interpretation of microwave data, Tyblewski *et al.*¹⁷ reported almost all the structural parameters of the CH₃OOH molecule. In a more detailed study, Tyblewski *et al.*¹⁶ gave only one parameter adjusted to the observed data, $\angle C—C—O = 105.3^\circ$. All other structural parameters were adopted by those authors from *ab initio* calculations. Conformational properties of methyl hydroperoxide have been investigated theoretically.^{1,2,4-6,11,16,18-22} *Ab initio*^{5,16,19-22} and molecular mechanics^{4,6} calculations provide an equilibrium molecular structure and *trans* barrier in close agreement with the conclusions of experimental study of Tyblewski *et al.*¹⁶

The product of the principal moments of inertia of methyl hydroperoxide was calculated in this work using the rotational constants determined from a microwave study.¹⁶ Structural parameters for the *skew* C₁ symmetry conformation given above are based on the experimental¹⁶ and theoretical^{4-6,16,19-22} data for CH₃OOH and from comparison with the structural parameters of CH₃OOCH₃.²³ These parameters give values for the rotational constants which differ by only 0.4%–0.6% from the observed values.

Methyl hydroperoxide contains OH and CH₃ groups, which rotate around O—O and C—O bonds. Contributions to the thermodynamic functions from these hindered rotors were evaluated in this work based on available data on rotational barriers in CH₃OOH. The value of $V_3 = 1120$ cm⁻¹ was accepted for the barrier to internal rotation of the methyl group around the C—O bond. This value was found from the microwave study¹⁶ and agrees closely with values obtained from *ab initio* calculations.^{4,16} The double-minimum potential energy function,

$$V(\varphi) = V_0 + V_1 \cos \varphi + V_2 \cos 2\varphi + V_3 \cos 3\varphi,$$

was used in this work for the O—O internal rotation. This function was chosen earlier for the hindered rotation potential function in hydrogen peroxide.²⁴ The expansion coefficients V_0 , V_1 , V_2 , and V_3 can be expressed in terms of the *trans* barrier height V_{trans} , the *cis* barrier height V_{cis} , and the COOH dihedral angle φ_e corresponding to a minimum of the potential function.²⁴ The values of V_0 , V_1 , V_2 , and V_3 were calculated in this work assuming $V_{trans} = 172.5$ cm⁻¹, $V_{cis} = 2500$ cm⁻¹, and $\varphi_e = 120^\circ$. The value of V_{trans}

$= 172.5$ cm⁻¹ ($\angle C—O—O—H = 180^\circ$) was taken from the microwave study.¹⁶ It should be noted that substantially lower values of V_{trans} from 80 to 126 cm⁻¹ were predicted by *ab initio*^{5,16,19,20} and molecular mechanics⁶ calculations, whereas high V_{trans} values of 240–590 cm⁻¹ were found from other molecular mechanics studies.^{4,5} No experimental data for the V_{cis} rotational barrier have been published. Its value accepted in this work is based on theoretical results^{4-6,16,25} which agree closely with each other. The equilibrium COOH dihedral angle, $\varphi_e = 120^\circ$, was estimated in this work (see above). The values of the reduced moments of inertia for CH₃ and OH tops were calculated from structural parameters given above.

Experimental data on the vibrational spectra of methyl hydroperoxide are unknown. Vibrational frequencies of CH₃OOH were predicted from *ab initio* calculations.^{16,22} In this work, fundamental frequencies of methyl hydroperoxide were estimated by normal coordinate calculations using force constants transferred from the CH₃OOCH₃ and H₂O₂ molecules:

| | | | | | |
|-------------|-------|-----------------|--------|-----------------|--------|
| $f_{O—O}$ | 5.000 | $f_{O—O—H}$ | 0.940 | $f_{C—H,H—C—O}$ | -0.056 |
| $f_{C—O}$ | 5.529 | $f_{tors(C—O)}$ | 0.094 | $f_{C—H,H—C—H}$ | 0.090 |
| $f_{C—H}$ | 4.692 | $f_{tors(O—O)}$ | 0.011 | $f_{C—O,H—C—O}$ | 0.240 |
| $f_{O—H}$ | 7.173 | $f_{C—H,C—H}$ | 0.041 | $f_{C—O,C—O—O}$ | 1.735 |
| $f_{C—O—O}$ | 1.512 | $f_{C—H,C—O}$ | 0.250 | $f_{O—O,C—O—O}$ | 0.617 |
| $f_{H—C—O}$ | 0.820 | $f_{C—O,O—O}$ | 0.267 | $f_{O—O,O—O—H}$ | 0.919 |
| $f_{H—C—H}$ | 0.516 | $f_{O—O,O—H}$ | -0.317 | $f_{O—H,O—O—H}$ | 0.925 |

(stretching and stretch-stretch interaction constants are in units of mdyn/Å; bend and torsion constants are in units of mdyn Å; stretch-bend interaction constants are in units of mdyn). The valence force fields for CH₃OOCH₃ and H₂O₂ were obtained from normal coordinate calculations using known vibrational assignments.^{6,26} Comparison of our frequencies with those obtained in *ab initio* calculations^{16,22} shows satisfactory agreement taking into account the fact that the *ab initio* frequencies are unscaled and the anharmonicity contributions are neglected.

The uncertainties in the calculated thermodynamic functions (Table 8) may reach (2–4) J K⁻¹ mol⁻¹ for $C_p^\circ(T)$ and (3–5) J K⁻¹ mol⁻¹ for $S^\circ(T)$. They are caused by the uncertainties in the adopted vibrational frequencies and the approximate model used for internal rotation.

Ideal gas thermodynamic properties of methyl hydroperoxide have not been reported previously.

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TABLE 8. Ideal gas thermodynamic properties of methyl hydroperoxide $\text{CH}_3\text{O}_2(\text{g})$ at the standard state pressure, $p^\circ = 0.1 \text{ MPa}$ ($T_r = 298.15 \text{ K}$)

| T (K) | C_p° ($\text{J K}^{-1} \text{ mol}^{-1}$) | S° ($\text{J K}^{-1} \text{ mol}^{-1}$) | $-(G^\circ - H^\circ(T_r))/T$ ($\text{J K}^{-1} \text{ mol}^{-1}$) | $H^\circ - H^\circ(T_r)$ (kJ mol^{-1}) | $\Delta_f H^\circ$ (kJ mol^{-1}) | $\Delta_f G^\circ$ (kJ mol^{-1}) | $\log K_f^\circ$ |
|------------|---|---|---|--|--|--|------------------|
| 0 | 0.000 | 0.000 | ∞ | -13.919 | -126.247 | -126.247 | ∞ |
| 25 | 34.489 | 171.468 | 692.887 | -13.036 | -127.153 | -125.355 | 261.912 |
| 50 | 35.540 | 195.488 | 438.841 | -12.168 | -128.683 | -123.025 | 128.522 |
| 75 | 38.512 | 210.437 | 360.349 | -11.243 | -130.278 | -119.836 | 83.460 |
| 100 | 41.627 | 221.948 | 324.360 | -10.241 | -131.534 | -116.158 | 60.674 |
| 150 | 46.975 | 239.876 | 293.352 | -8.021 | -133.584 | -108.013 | 37.613 |
| 180 | 49.820 | 248.694 | 285.189 | -6.569 | -134.723 | -102.791 | 29.829 |
| 190 | 50.758 | 251.413 | 283.340 | -6.066 | -135.098 | -101.007 | 27.768 |
| 200 | 51.703 | 254.040 | 281.810 | -5.554 | -135.472 | -99.203 | 25.909 |
| 210 | 52.659 | 256.586 | 280.549 | -5.032 | -135.845 | -97.380 | 24.222 |
| 220 | 53.627 | 259.058 | 279.516 | -4.501 | -136.215 | -95.540 | 22.684 |
| 230 | 54.610 | 261.463 | 278.679 | -3.960 | -136.583 | -93.683 | 21.276 |
| 240 | 55.609 | 263.808 | 278.011 | -3.408 | -136.949 | -91.810 | 19.982 |
| 250 | 56.624 | 266.099 | 277.489 | -2.847 | -137.312 | -89.922 | 18.788 |
| 260 | 57.653 | 268.340 | 277.094 | -2.276 | -137.671 | -88.019 | 17.683 |
| 270 | 58.696 | 270.535 | 276.810 | -1.694 | -138.026 | -86.103 | 16.657 |
| 280 | 59.751 | 272.689 | 276.625 | -1.102 | -138.376 | -84.173 | 15.703 |
| 290 | 60.816 | 274.804 | 276.526 | -0.499 | -138.722 | -82.231 | 14.811 |
| 298.15 | 61.690 | 276.502 | 276.502 | 0.000 | -139.000 | -80.640 | 14.128 |
| 300 | 61.889 | 276.884 | 276.503 | 0.114 | -139.063 | -80.277 | 13.977 |
| 350 | 67.290 | 286.830 | 277.276 | 3.344 | -140.679 | -70.350 | 10.499 |
| 400 | 72.583 | 296.163 | 279.059 | 6.841 | -142.143 | -60.202 | 7.861 |
| 450 | 77.604 | 305.005 | 281.455 | 10.597 | -143.453 | -49.879 | 5.790 |
| 500 | 82.279 | 313.427 | 284.234 | 14.596 | -144.620 | -39.418 | 4.118 |
| 600 | 90.573 | 329.181 | 290.432 | 23.250 | -146.559 | -18.190 | 1.584 |
| 700 | 97.633 | 343.688 | 297.017 | 32.670 | -148.043 | 3.328 | -0.248 |
| 800 | 103.693 | 357.130 | 303.701 | 42.743 | -149.133 | 25.032 | -1.634 |
| 900 | 108.935 | 369.654 | 310.342 | 53.381 | -149.884 | 46.850 | -2.719 |
| 1000 | 113.485 | 381.372 | 316.865 | 64.507 | -150.352 | 68.737 | -3.590 |
| 1100 | 117.439 | 392.378 | 323.235 | 76.058 | -150.582 | 90.659 | -4.305 |
| 1200 | 120.878 | 402.748 | 329.433 | 87.978 | -150.618 | 112.593 | -4.901 |
| 1300 | 123.871 | 412.544 | 335.453 | 100.219 | -150.500 | 134.525 | -5.405 |
| 1400 | 126.478 | 421.822 | 341.294 | 112.739 | -150.261 | 156.442 | -5.837 |
| 1500 | 128.752 | 430.627 | 346.958 | 125.503 | -149.928 | 178.338 | -6.210 |
| 1600 | 130.741 | 439.001 | 352.451 | 138.480 | -149.526 | 200.219 | -6.536 |
| 1700 | 132.485 | 446.981 | 357.779 | 151.643 | -149.075 | 222.061 | -6.823 |
| 1800 | 134.018 | 454.598 | 362.948 | 164.970 | -148.588 | 243.882 | -7.077 |
| 1900 | 135.369 | 461.881 | 367.965 | 178.441 | -148.080 | 265.672 | -7.304 |
| 2000 | 136.564 | 468.855 | 372.836 | 192.039 | -147.563 | 287.436 | -7.507 |

10. Dimethyl Peroxide, CH₃—O—O—CH₃Dimethyl peroxide (C₂H₆O₂)

Ideal gas

 $M_r = 62.0682$ $S^\circ(298.15\text{ K}) = 308.4 \pm 3.0\text{ J K}^{-1}\text{ mol}^{-1}$ $\Delta_f H^\circ(0\text{ K}) = -106.5 \pm 5.0\text{ kJ mol}^{-1}$
 $\Delta_f H^\circ(298.15\text{ K}) = -125.5 \pm 5.0\text{ kJ mol}^{-1}$

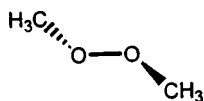
Molecular constants

Point group: C_2 Symmetry number: $\sigma = 2$ Number of optical isomers: $n = 2$ Ground electronic state: \tilde{X}^1A Energy: $\epsilon_X = 0\text{ cm}^{-1}$ Quantum weight: $g_X = 1$ Vibrational frequencies, ν_i , and degeneracies, g_i

| Symmetry | ν_i, cm^{-1} | g_i | Symmetry | ν_i, cm^{-1} | g_i |
|----------|-------------------------|------------------|----------|-------------------------|------------|
| A | ν_1 | 2945 | 1 | ν_{13} | ... |
| | ν_2 | 2917 | 1 | B | ν_{14} |
| | ν_3 | 2900 | 1 | | ν_{15} |
| | ν_4 | 1487 | 1 | | ν_{16} |
| | ν_5 | 1474 | 1 | | ν_{17} |
| | ν_6 | 1433 | 1 | | ν_{18} |
| | ν_7 | 1198 | 1 | | ν_{19} |
| | ν_8 | 1165 | 1 | | ν_{20} |
| | ν_9 | 1020 | 1 | | ν_{21} |
| | ν_{10} | 786 | 1 | | ν_{22} |
| | ν_{11} | 448 | 1 | | ν_{23} |
| | ν_{12} | ... ^a | 1 | | ν_{24} |
| | | | | | ... |

^aInstead of torsional modes $\nu_{12} = 218\text{ cm}^{-1}$ and $\nu_{24} = 231\text{ cm}^{-1}$, the contributions due to the internal rotation of CH₃ groups around C—O bonds were calculated from the potential $V(\varphi) = \frac{1}{2}V_3(1 - \cos 3\varphi)$, where φ is the H—C—O—O torsional angle and $V_3 = 900\text{ cm}^{-1}$.

^bInstead of torsional mode $\nu_{13} = 73\text{ cm}^{-1}$, the contributions due to the internal rotation of OCH₃ group around O—O bond were calculated from the potential $V(\varphi) = V_0 + V_1 \cos \varphi + V_2 \cos 2\varphi + V_3 \cos 3\varphi$, where φ is the C—O—O—C torsional angle, $V_0 = 1341.3$, $V_1 = 2081.0$, $V_2 = 1052.2$, and $V_3 = 225.5$ (in cm^{-1}).

CH₃ top: Reduced moment of inertia, $I_r = 0.4910 \times 10^{-39}\text{ g cm}^2$, Symmetry number, $\sigma_m = 3$.OCH₃ top: Reduced moment of inertia, $I_r = 1.5928 \times 10^{-39}\text{ g cm}^2$, Symmetry number, $\sigma_m = 1$.

Geometry

 $r(\text{C—O}) = 1.420 \pm 0.007\text{ \AA}$ $r(\text{O—O}) = 1.457 \pm 0.012\text{ \AA}$ $r(\text{C—H}) = 1.099 \pm 0.004\text{ \AA}$ $\angle \text{C—O—O} = 105.2 \pm 0.5^\circ$ $\angle \text{H—C—H} = 110.1 \pm 0.7^\circ$ $\phi(\text{C—O—O—C}) = 119 \pm 4^\circ$ Product of moments of inertia: $I_A I_B I_C = 1123 \times 10^{-117}\text{ g}^3\text{ cm}^6$.

10.1. Enthalpy of Formation

The enthalpy of formation of dimethyl peroxide recommended in this work ($-125.5\text{ kJ mol}^{-1}$) was determined by Baker *et al.*¹ from calorimetric measurements of the enthalpy of combustion. Slightly lower values of

$\Delta_f H^\circ(298.15\text{ K})$ from $-129.3\text{ kJ mol}^{-1}$ to $-133.9\text{ kJ mol}^{-1}$ were estimated by group additivity^{2,3} and molecular mechanics calculations.^{4,5} A lower value is estimated from the experimental value of $\Delta_f H^\circ(\text{CH}_3\text{O})$ Gurvich *et al.*⁶ and the bond energy for the O—O bond in CH_3OOCH_3 .⁷

$$\begin{aligned}\Delta_f H^\circ(\text{CH}_3\text{OOCH}_3) &= 2\Delta_f H^\circ(\text{CH}_3\text{O}) \\ &\quad - D^\circ(\text{CH}_3\text{O}-\text{OCH}_3) \\ &= 2 \times 13 - 160 = -134 \text{ kJ mol}^{-1}.\end{aligned}$$

A still more negative value, $-138.1 \text{ kJ mol}^{-1}$, was proposed by Benassi *et al.*⁸ employing theoretical and empirical approaches. Semiempirical MINDO^{9,10} and molecular mechanics¹¹ calculations provide $\Delta_f H^\circ$ (298.15 K) values, -105.1 to $-121.6 \text{ kJ mol}^{-1}$, which are greater than the experimental value. The available theoretical and empirical evaluations are grounds to believe that the experimental value of $\Delta_f H^\circ$ (298.15 K) may be overestimated. Thus, its uncertainty was increased in this work.

10.2. Heat Capacity and Entropy

The molecular structure of dimethyl peroxide, CH_3OOCH_3 , was studied by electron diffraction.¹² The molecular intensities were analyzed using a model which considered the internal rotation about the O—O bond. The equilibrium geometry was determined to be the *skew* conformation with a C—O—O—C dihedral angle of 119° . The barrier in the *trans* configuration ($\angle \text{C—O—O—C} = 180^\circ$) was found to be $87^{+87}_{-52} \text{ cm}^{-1}$. This result agrees with an analysis of infrared and Raman spectra¹³ and a normal coordinate analysis based on these data¹⁴ as well as with results of semiempirical,^{9,10,15,16} *ab initio*¹⁷⁻¹⁹ and molecular mechanics^{4,5,8} calculations. Photoelectron spectra investigations^{20,21} and *ab initio* calculations^{8,22-24} support an exactly planar or nearly planar *trans* configuration. However, it should be noted that the value of a dihedral angle depends strongly on the basis set and method used.¹⁹ Hamada and Morishita²⁵ have described the Raman and infrared spectra of CH_3OOCH_3 in terms of planar structure of D_{3h} symmetry. In this work, the product of the principal moments of inertia of dimethyl peroxide for *skew* conformation of C_2 symmetry was calculated using the structural parameters determined from electron diffraction study.¹²

The dimethyl peroxide molecule undergoes three large-amplitude motions: an internal rotation of the OCH_3 groups about the O—O bond and internal rotation of the CH_3 groups about two C—O bonds. Contributions to the thermodynamic properties from these hindered rotors were calculated in this work based on available data on the rotational barriers in CH_3OOCH_3 . The value of $V_3 = 900 \text{ cm}^{-1}$ was accepted for the internal rotation barrier of methyl groups around the C—O bonds. This value was used by Koput²⁶ in his theoretical model describing the internal rotation in dimethyl peroxide and it is the average of the results of molecular mechanics calculations.⁴ The double-minimum potential energy function,

$$V(\varphi) = V_0 + V_1 \cos \varphi + V_2 \cos 2\varphi + V_3 \cos 3\varphi,$$

was used in this work for the internal rotation around the O—O bond. This function was chosen earlier for the hindered rotation potential function in hydrogen peroxide.²⁷ The expansion coefficients V_0 , V_1 , V_2 , and V_3 can be expressed

in terms of the *trans* barrier height V_{trans} , the *cis* barrier height V_{cis} , and the COOC dihedral angle φ_e corresponding to a minimum of the potential function.²⁷ The values of V_0 , V_1 , V_2 , and V_3 were calculated in this work assuming $V_{trans} = 87 \text{ cm}^{-1}$, $V_{cis} = 4700 \text{ cm}^{-1}$, and $\varphi_e = 119^\circ$. The value of V_{trans} was taken from the electron diffraction study.¹² It agrees with results of theoretical calculations^{4,5,8,19} within the limits of experimental accuracy. The adopted value of V_{cis} is based on results of *ab initio*^{8,18} and molecular mechanics^{4,5} calculations. The values of the reduced moments of inertia for CH_3 and OCH_3 tops were calculated from structural parameters given above.

Christe¹³ has investigated the infrared spectra of CH_3OOCH_3 and CD_3OCD_3 both in the gas phase and in an Ar matrix and the Raman spectra of these molecules in the liquid phase. The vibrational assignment proposed by Christe¹³ has considerable uncertainty for a number of the vibrations, and the low-frequency modes were not assigned at all. Bell and Laane¹⁴ have carried out a normal coordinate analysis of dimethyl peroxide using the experimental data of Christe.¹³ Hamada and Morishita²⁵ have described the vibrational spectra of CH_3OOCH_3 in terms of D_{3h} symmetry. However, this assignment is in conflict with available information on dimethyl peroxide. There are two *ab initio* calculations of vibrational frequencies of CH_3OOCH_3 .^{23,28} Results of Benassi and Taddei²³ show overall agreement with experimental data taking into account the fact that the *ab initio* frequencies are unscaled and the anharmonicity contributions are neglected. Chen and Allinger⁵ have carried out the molecular mechanics calculation of vibrational frequencies of dimethyl peroxide. Their results are in agreement with the available, but somewhat tentative and incomplete, experimental data.¹³ The assignment proposed by Chen and Allinger⁵ is accepted in this work.

The uncertainties in the calculated thermodynamic functions (Table 9) may reach $(2-4) \text{ J K}^{-1} \text{ mol}^{-1}$ for $C_p^\circ(T)$ and $(3-5) \text{ J K}^{-1} \text{ mol}^{-1}$ for $S^\circ(T)$. They are caused by uncertainties in the adopted vibrational frequencies and the approximate treatment of internal rotation.

Ideal gas thermodynamic properties of dimethyl peroxide have not been reported previously.

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TABLE 9. Ideal gas thermodynamic properties of dimethyl peroxide $C_2H_6O_2(g)$ at the standard state pressure, $p^\circ = 0.1 \text{ MPa}$ ($T_r = 298.15 \text{ K}$)

| T (K) | C_p° ($\text{J K}^{-1} \text{ mol}^{-1}$) | S° ($\text{J K}^{-1} \text{ mol}^{-1}$) | $-(G^\circ - H^\circ(T_r))/T$ ($\text{J K}^{-1} \text{ mol}^{-1}$) | $H^\circ - H^\circ(T_r)$ (kJ mol^{-1}) | $\Delta_f H^\circ$ (kJ mol^{-1}) | $\Delta_f G^\circ$ (kJ mol^{-1}) | $\log K_f^\circ$ |
|------------|---|---|---|--|--|--|------------------|
| 0 | 0.000 | 0.000 | ∞ | -17.153 | -106.460 | -106.460 | ∞ |
| 25 | 38.665 | 180.406 | 831.484 | -16.277 | -107.908 | -104.953 | 219.284 |
| 50 | 42.071 | 208.414 | 513.657 | -15.262 | -110.130 | -101.224 | 105.747 |
| 75 | 45.713 | 226.123 | 415.019 | -14.167 | -112.460 | -96.245 | 67.030 |
| 100 | 49.953 | 239.851 | 369.568 | -12.972 | -114.304 | -90.552 | 47.299 |
| 150 | 58.133 | 261.681 | 330.120 | -10.266 | -117.336 | -78.000 | 27.162 |
| 180 | 62.664 | 272.683 | 319.647 | -8.453 | -119.030 | -69.973 | 20.305 |
| 190 | 64.145 | 276.111 | 317.266 | -7.819 | -119.590 | -67.232 | 18.483 |
| 200 | 65.624 | 279.439 | 315.292 | -7.171 | -120.149 | -64.462 | 16.836 |
| 210 | 67.105 | 282.676 | 313.662 | -6.507 | -120.707 | -61.664 | 15.338 |
| 220 | 68.594 | 285.832 | 312.325 | -5.828 | -121.264 | -58.840 | 13.970 |
| 230 | 70.095 | 288.915 | 311.241 | -5.135 | -121.818 | -55.990 | 12.716 |
| 240 | 71.609 | 291.930 | 310.374 | -4.426 | -122.370 | -53.116 | 11.560 |
| 250 | 73.139 | 294.884 | 309.695 | -3.703 | -122.919 | -50.219 | 10.493 |
| 260 | 74.685 | 297.783 | 309.181 | -2.964 | -123.465 | -47.301 | 9.503 |
| 270 | 76.247 | 300.630 | 308.812 | -2.209 | -124.006 | -44.360 | 8.582 |
| 280 | 77.823 | 303.432 | 308.570 | -1.439 | -124.542 | -41.401 | 7.723 |
| 290 | 79.413 | 306.190 | 308.441 | -0.652 | -125.072 | -38.422 | 6.921 |
| 298.15 | 80.717 | 308.409 | 308.409 | 0.000 | -125.500 | -35.981 | 6.304 |
| 300 | 81.014 | 308.910 | 308.411 | 0.150 | -125.597 | -35.425 | 6.168 |
| 350 | 89.107 | 322.005 | 309.427 | 4.402 | -128.110 | -20.196 | 3.014 |
| 400 | 97.131 | 334.429 | 311.782 | 9.059 | -130.424 | -4.619 | 0.603 |
| 450 | 104.858 | 346.320 | 314.963 | 14.110 | -132.528 | 11.235 | -1.304 |
| 500 | 112.164 | 357.750 | 318.674 | 19.538 | -134.426 | 27.313 | -2.853 |
| 600 | 125.390 | 379.396 | 327.012 | 31.431 | -137.633 | 59.970 | -5.221 |
| 700 | 136.887 | 399.609 | 335.955 | 44.558 | -140.119 | 93.110 | -6.948 |
| 800 | 146.892 | 418.557 | 345.108 | 58.759 | -141.956 | 126.561 | -8.263 |
| 900 | 155.605 | 436.373 | 354.269 | 73.893 | -143.219 | 160.207 | -9.298 |
| 1000 | 163.187 | 453.169 | 363.327 | 89.842 | -143.993 | 193.966 | -10.132 |
| 1100 | 169.774 | 469.039 | 372.223 | 106.498 | -144.351 | 227.782 | -10.816 |
| 1200 | 175.490 | 484.062 | 380.922 | 123.768 | -144.366 | 261.615 | -11.388 |
| 1300 | 180.449 | 498.309 | 389.409 | 141.570 | -144.105 | 295.441 | -11.871 |
| 1400 | 184.752 | 511.843 | 397.673 | 159.536 | -143.625 | 329.235 | -12.284 |
| 1500 | 188.491 | 524.721 | 405.719 | 178.502 | -142.971 | 362.990 | -12.640 |
| 1600 | 191.746 | 536.992 | 413.543 | 197.518 | -142.187 | 396.716 | -12.951 |
| 1700 | 194.585 | 548.703 | 421.152 | 216.838 | -141.305 | 430.365 | -13.223 |
| 1800 | 197.069 | 559.898 | 428.551 | 236.423 | -140.351 | 463.971 | -13.464 |
| 1900 | 199.247 | 570.612 | 435.748 | 256.241 | -139.347 | 497.516 | -13.677 |
| 2000 | 201.164 | 580.882 | 442.750 | 276.264 | -138.314 | 531.010 | -13.868 |

11. Diacetyl Peroxide, $\text{CH}_3\text{—CO—O—O—CO—CH}_3$ Diacetyl peroxide ($\text{C}_4\text{H}_6\text{O}_4$)

Ideal gas

 $M_r = 118.089$ $S^\circ(298.15\text{ K}) = (390.7 \pm 6.0) \text{ J K}^{-1} \text{ mol}^{-1}$ $\Delta_f H^\circ(0\text{ K}) = -(477.0 \pm 10.0) \text{ kJ mol}^{-1}$ $\Delta_f H^\circ(298.15\text{ K}) = -(500.0 \pm 10.0) \text{ kJ mol}^{-1}$

Molecular constants

Point group: C_2 Symmetry number: $\sigma = 2$ Number of optical isomers: $n = 2$ Ground electronic state: \tilde{X}^1A Energy: $\epsilon_X = 0 \text{ cm}^{-1}$ Quantum weight: $g_X = 1$ Vibrational frequencies, ν_i , and degeneracies g_i

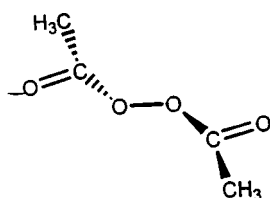
| Symmetry | ν_i, cm^{-1} | g_i | Symmetry | ν_i, cm^{-1} | g_i | |
|----------|-------------------------|------------------|----------|-------------------------|------------|------------------|
| A | ν_1 | 3001 | 1 | ν_{19} | ... | |
| | ν_2 | 2998 | 1 | B | ν_{20} | 3001 |
| | ν_3 | 2934 | 1 | | ν_{21} | 2998 |
| | ν_4 | 1762 | 1 | | ν_{22} | 2934 |
| | ν_5 | 1466 | 1 | | ν_{23} | 1749 |
| | ν_6 | 1428 | 1 | | ν_{24} | 1457 |
| | ν_7 | 1426 | 1 | | ν_{25} | 1427 |
| | ν_8 | 1290 | 1 | | ν_{26} | 1426 |
| | ν_9 | 1087 | 1 | | ν_{27} | 1285 |
| | ν_{10} | 927 | 1 | | ν_{28} | 962 |
| | ν_{11} | 875 | 1 | | ν_{29} | 922 |
| | ν_{12} | 850 | 1 | | ν_{30} | 808 |
| | ν_{13} | 648 | 1 | | ν_{31} | 639 |
| | ν_{14} | 631 | 1 | | ν_{32} | 619 |
| | ν_{15} | 392 | 1 | | ν_{33} | 545 |
| | ν_{16} | 285 | 1 | | ν_{34} | 328 |
| | ν_{17} | 151 | 1 | | ν_{35} | 140 |
| | ν_{18} | ... ^a | 1 | | ν_{36} | ... ^b |

^aInstead of torsional mode $\nu_{18} \sim 70 \text{ cm}^{-1}$, the contributions due to the internal rotation of O—CO—CH_3 group around O—O bond were calculated from the potential $V(\varphi) = V_0 + V_1 \cos \varphi + V_2 \cos 2\varphi + V_3 \cos 3\varphi$, where φ is the C—O—O—C torsional angle, $V_0 = 1433.9$, $V_1 = 2222.2$, $V_2 = 1111.1$, and $V_3 = 232.8$ (in cm^{-1}).

^bInstead of torsional modes ν_{19} and $\nu_{36} \sim 60 \text{ cm}^{-1}$, the contributions due to the internal rotation of CH_3 groups around C—C bonds were calculated from the potential $V(\varphi) = \frac{1}{2}V_3(1 - \cos 3\varphi)$, where φ is the H—C—C—O torsional angle and $V_3 = 80 \text{ cm}^{-1}$.

O—CO—CH_3 top: Reduced moment of inertia, $I_r = 4.5200 \times 10^{-39} \text{ g cm}^2$, Symmetry number, $\sigma_m = 1$.

CH_3 top: Reduced moment of inertia, $I_r = 0.5182 \times 10^{-39} \text{ g cm}^2$, Symmetry number, $\sigma_m = 3$.



Geometry

 $r(\text{C—C}) = 1.495 \pm 0.01 \text{ \AA}$ $\angle \text{O=C—O} = 122 \pm 3^\circ$ $r(\text{C=O}) = 1.21 \pm 0.02 \text{ \AA}$ $\angle \text{C—O—O} = 113 \pm 3^\circ$ $r(\text{C—O}) = 1.34 \pm 0.02 \text{ \AA}$ $\angle \text{H—C—H} = 109.5 \pm 2.0^\circ$ $r(\text{O—O}) = 1.46 \pm 0.02 \text{ \AA}$ $\varphi(\text{C—O—O—C}) = 120 \pm 10^\circ$ $r(\text{C—H}) = 1.09 \pm 0.01 \text{ \AA}$ $\varphi(\text{O=C—O—O}) = 0^\circ$ $\angle \text{C—C=O} = 128 \pm 3^\circ$

Product of moments of inertia: $I_A I_B I_C = 96075 \times 10^{-117} \text{ g}^3 \text{ cm}^6$.

11.1. Enthalpy of Formation

Jaffe *et al.*¹ determined the enthalpy of formation for liquid diacetyl peroxide ($\text{CH}_3\text{CO—OO—COCH}_3$) from calorimetric measurements. Assuming the enthalpy of vaporization of diacetyl peroxide to be equal to the enthalpy of vaporization of acetic anhydride ($\text{CH}_3\text{CO—O—COCH}_3$), the authors obtained the value of -498 kJ mol^{-1} for the enthalpy of formation of gaseous diacetyl peroxide. A similar value of $\Delta_f H^\circ(298.15 \text{ K}) = -502 \text{ kJ mol}^{-1}$ may be derived from the group additivity contributions for peroxyacids and peroxyesters obtained by Benassi and Taddei² using empirical approaches and *ab initio* calculations. The value predicted by the method of group equations,

$$\begin{aligned} \Delta_f H^\circ(\text{CH}_3\text{CO—OO—COCH}_3) \\ = \Delta_f H^\circ(\text{CH}_3\text{CO—O—COCH}_3) + \Delta_f H^\circ(\text{CH}_3\text{OOCH}_3) \\ - \Delta_f H^\circ(\text{CH}_3\text{OCH}_3) = (-572.5) + (-125.5) \\ - (-184.1) = -513.9 \end{aligned}$$

(the $\Delta_f H^\circ$ values for the related compounds were taken from the compilation by Pedley³), agrees with the two above estimates within the limits of the combined errors of those determinations. Considerably lower values of $\Delta_f H^\circ$ for diacetyl peroxide were estimated from a group additivity approximation (-540 kJ mol^{-1})⁴ and a semiempirical calculation ($\sim -585 \text{ kJ mol}^{-1}$).^{5,6} The value recommended in this work, $-(500 \pm 10) \text{ kJ mol}^{-1}$, is based on the estimates.^{1,2}

11.2. Heat Capacity and Entropy

Experimental data on molecular structure and vibrational frequencies of diacetyl peroxide, $\text{CH}_3\text{CO—OO—COCH}_3$, are unknown. Semiempirical MINDO calculations⁵ predict that the *skew* conformation of C_2 symmetry with a C—O—O—C dihedral angle of 108° is the most stable. Although a planar structure was suggested for the peroxy acid group in peroxyacetic acid from a microwave study,⁷ the *skew* conformation of C_2 symmetry is accepted in this work for diacetyl peroxide in accord with a semiempirical calculation⁵ and by analogy with dimethyl peroxide.⁸ The product of the principal moments of inertia was calculated

from structural parameters estimated by comparison with structural parameters of $\text{CH}_3\text{CO—OOH}$,⁷ $\text{CH}_3\text{CO—O—COCH}_3$,⁹ and CH_3OOCH_3 .⁸

The diacetyl peroxide molecule undergoes five large-amplitude motions: an internal rotation about a central O—O bond and about two C—C and two C—O bonds. Contributions to the thermodynamic functions due to internal rotation of CH_3 groups were calculated in this work assuming the V_3 barrier height to be the same as that in peroxyacetic acid.⁷ The double-minimum potential energy function,

$$V(\varphi) = V_0 + V_1 \cos \varphi + V_2 \cos 2\varphi + V_3 \cos 3\varphi,$$

was used for the internal rotation about the O—O bond. This function was chosen earlier for the hindered rotation potential function in hydrogen peroxide.¹⁰ The expansion coefficients V_0 , V_1 , V_2 , and V_3 can be expressed in terms of the *trans* barrier height V_{trans} , the *cis* barrier height V_{cis} , and the COOC dihedral angle φ_e corresponding to a minimum of the potential function.¹⁰ The values of V_0 , V_1 , V_2 , and V_3 were calculated in this work assuming $V_{trans} = 90 \text{ cm}^{-1}$, $V_{cis} = 5000 \text{ cm}^{-1}$, and $\varphi_e = 120^\circ$. The values of V_{trans} and V_{cis} are derived from the results of a MINDO calculation⁵ taking into account that the same calculation overestimated the V_{trans} and underestimated the V_{cis} barriers for the CH_3OOH and CH_3OOCH_3 molecules. The values of the reduced moments of inertia for CH_3 and CH_3COO tops were calculated from structural parameters given above. There are no data on barriers hindering the internal rotation of CH_3CO groups about the C—O bonds. Frequencies $\nu_{17} = 151$ and $\nu_{35} = 140 \text{ cm}^{-1}$ were accepted in this work for the torsional motion about the C—O bonds.

Zyatkov *et al.*¹¹ have calculated vibrational frequencies of diacetyl peroxide using a theoretical simulation of the force field. In this work, vibrational frequencies of $\text{CH}_3\text{CO—OO—COCH}_3$ were calculated using force constants transferred from the $\text{CH}_3\text{CO—OOH}$ molecule. The simplified valence force field for that molecule was obtained by normal coordinate calculations for the vibrational assignment of Cugley *et al.*¹² The calculated force constants reproduce the observed vibrational wave numbers of $\text{CH}_3\text{CO—OOH}$ with a root mean square deviation of 0.6 cm^{-1} . The following 18 force constants were used for calculating the vibrational frequencies of diacetyl peroxide:

| | | | | | |
|--------------------|--------|---------------------------------------|-------|---|--------|
| $f_{\text{O—O}}$ | 3.821 | $f_{\text{C—C—O}}$ | 2.363 | $f_{\text{tors(C—O)}}$ | 0.430 |
| $f_{\text{C—O}}$ | 3.827 | $f_{\text{O=C—O}} = f_{\text{O=C—C}}$ | 1.972 | $f_{\text{tors(O—O)}}$ | 0.104 |
| $f_{\text{C=O}}$ | 11.382 | $f_{\text{C—C—H}}$ | 0.497 | $f_{\text{C=O,C—C}} = f_{\text{C=O,C—O}}$ | 1.679 |
| $f_{\text{C—C}}$ | 3.480 | $f_{\text{H—C—H}}$ | 0.523 | $f_{\text{C—C,C—O}}$ | -0.356 |
| $f_{\text{C—H}}$ | 4.836 | $f_{\text{wag(C=O)}}$ | 0.515 | $f_{\text{C—H,C—C—H}}$ | 0.415 |
| $f_{\text{C—O—O}}$ | 1.659 | $f_{\text{tors(C—C)}}$ | 0.006 | $f_{\text{C—O,O=C—O}}$ | 0.232 |

(stretching and stretch-stretch interaction constants are in units of $\text{mdyn}/\text{\AA}$; bending, wagging, and torsion constants are in units of $\text{mdyn}/\text{\AA}$; stretch-bend interaction constants are in units of mdyn).

The uncertainties in the calculated thermodynamic functions (Table 10) may reach $(8-12) \text{ J K}^{-1} \text{ mol}^{-1}$ for $C_p^\circ(T)$ and $(6-10) \text{ J K}^{-1} \text{ mol}^{-1}$ for $S^\circ(T)$. They are caused by the uncertainties in the adopted vibrational frequencies and the approximate treatment of internal rotation.

Ideal gas thermodynamic properties of diacetyl peroxide have not been reported previously.

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TABLE 10. Ideal gas thermodynamic properties of diacetyl peroxide, $\text{C}_4\text{H}_6\text{O}_4(\text{g})$ at the standard state pressure, $p^\circ=0.1 \text{ MPa}$ ($T_r=298.15 \text{ K}$)

| T (K) | C_p° ($\text{J K}^{-1} \text{ mol}^{-1}$) | S° ($\text{J K}^{-1} \text{ mol}^{-1}$) | $-(G^\circ - H^\circ(T_r))/T$ ($\text{J K}^{-1} \text{ mol}^{-1}$) | $H^\circ - H^\circ(T_r)$ (kJ mol^{-1}) | $\Delta_f H^\circ$ (kJ mol^{-1}) | $\Delta_f G^\circ$ (kJ mol^{-1}) | $\log K_f^\circ$ |
|------------|---|---|---|--|--|--|------------------|
| 0 | 0.000 | 0.000 | ∞ | -23.944 | -476.961 | -476.961 | ∞ |
| 25 | 50.640 | 213.740 | 1130.731 | -22.925 | -478.990 | -473.540 | 989.395 |
| 50 | 58.035 | 251.723 | 682.615 | -21.545 | -481.588 | -467.166 | 488.038 |
| 75 | 62.084 | 276.009 | 543.259 | -20.044 | -484.278 | -459.352 | 319.917 |
| 100 | 66.834 | 294.500 | 478.840 | -18.434 | -486.503 | -450.697 | 235.417 |
| 150 | 78.569 | 323.725 | 422.461 | -14.810 | -490.316 | -431.949 | 150.416 |
| 180 | 86.766 | 338.764 | 407.276 | -12.332 | -492.443 | -420.074 | 121.901 |
| 190 | 89.636 | 343.532 | 403.796 | -11.450 | -493.135 | -416.035 | 114.374 |
| 200 | 92.559 | 348.204 | 400.900 | -10.539 | -493.820 | -411.959 | 107.591 |
| 210 | 95.524 | 352.791 | 398.500 | -9.599 | -494.496 | -407.849 | 101.446 |
| 220 | 98.523 | 357.304 | 396.525 | -8.629 | -495.163 | -403.707 | 95.851 |
| 230 | 101.547 | 361.750 | 394.917 | -7.628 | -495.821 | -399.536 | 90.736 |
| 240 | 104.588 | 366.136 | 393.626 | -6.598 | -496.468 | -395.336 | 86.041 |
| 250 | 107.640 | 370.468 | 392.614 | -5.536 | -497.104 | -391.109 | 81.717 |
| 260 | 110.696 | 374.749 | 391.844 | -4.445 | -497.729 | -386.857 | 77.720 |
| 270 | 113.750 | 378.984 | 391.290 | -3.323 | -498.342 | -382.581 | 74.014 |
| 280 | 116.796 | 383.176 | 390.925 | -2.170 | -498.943 | -378.283 | 70.569 |
| 290 | 119.830 | 387.327 | 390.730 | -0.987 | -499.530 | -373.963 | 67.357 |
| 298.15 | 122.291 | 390.682 | 390.682 | 0.000 | -500.000 | -370.427 | 64.896 |
| 300 | 122.848 | 391.441 | 390.685 | 0.227 | -500.105 | -369.622 | 64.356 |
| 350 | 137.551 | 411.491 | 392.235 | 6.740 | -502.777 | -347.661 | 51.885 |
| 400 | 151.377 | 430.772 | 395.855 | 13.967 | -505.120 | -325.337 | 42.484 |
| 450 | 164.144 | 449.350 | 400.773 | 21.860 | -507.157 | -302.738 | 35.141 |
| 500 | 175.813 | 467.258 | 406.531 | 30.363 | -508.918 | -279.929 | 29.244 |
| 600 | 196.085 | 501.163 | 419.512 | 48.990 | -511.704 | -233.859 | 20.359 |
| 700 | 212.888 | 532.691 | 433.456 | 69.465 | -513.643 | -187.388 | 13.983 |
| 800 | 226.936 | 562.062 | 447.717 | 91.476 | -514.847 | -140.688 | 9.186 |
| 900 | 238.778 | 589.495 | 461.963 | 114.778 | -515.418 | -93.879 | 5.449 |
| 1000 | 248.823 | 615.187 | 476.014 | 139.172 | -515.458 | -47.036 | 2.457 |
| 1100 | 257.379 | 639.314 | 489.774 | 164.493 | -515.048 | -0.210 | 0.010 |
| 1200 | 264.693 | 662.031 | 503.192 | 190.606 | -514.271 | 46.561 | -2.027 |
| 1300 | 270.965 | 683.471 | 516.243 | 217.397 | -513.201 | 93.260 | -3.747 |
| 1400 | 276.360 | 703.754 | 528.919 | 244.770 | -511.906 | 139.862 | -5.218 |
| 1500 | 281.017 | 722.984 | 541.221 | 272.645 | -510.431 | 186.370 | -6.490 |
| 1600 | 285.051 | 741.252 | 553.157 | 300.953 | -508.834 | 232.810 | -7.600 |
| 1700 | 288.556 | 758.641 | 564.737 | 329.637 | -507.144 | 279.105 | -8.576 |
| 1800 | 291.613 | 775.223 | 575.973 | 358.649 | -505.392 | 325.314 | -9.440 |
| 1900 | 294.288 | 791.063 | 586.880 | 387.947 | -503.606 | 371.413 | -10.211 |
| 2000 | 296.638 | 806.219 | 597.471 | 417.496 | -501.807 | 417.423 | -10.902 |

TABLE 11. Summary of the thermodynamic properties at 298.15 K and the standard state pressure, $p^\circ = 0.1$ MPa

| Molecule | S° (298.15 K) ($\text{J K}^{-1} \text{mol}^{-1}$) | C_p° (298.15 K) ($\text{J K}^{-1} \text{mol}^{-1}$) | $\Delta_f H^\circ$ (298.15 K) (kJ mol^{-1}) |
|------------------------|---|---|---|
| Bromoacetic acid | 337.0 ± 5.0 | 80.5 ± 3.0 | -383.5 ± 3.1 |
| Chloroacetic acid | 325.9 ± 5.0 | 78.8 ± 3.0 | -427.6 ± 1.0 |
| Oxopropanedinitrile | 310.0 ± 1.0 | 80.8 ± 1.0 | 247.5 ± 6.4 |
| Glycolic acid | 318.6 ± 5.0 | 87.1 ± 3.0 | -583.0 ± 10.0 |
| Glyoxal | 272.5 ± 1.0 | 60.6 ± 0.5 | -212.0 ± 0.8 |
| Cyanooxomethyl radical | 278.2 ± 1.5 | 56.1 ± 1.5 | 210.0 ± 10.0 |
| Oxalic acid | 320.6 ± 5.0 | 86.2 ± 5.0 | -731.8 ± 2.0 |
| Methyl hydroperoxide | 276.5 ± 3.0 | 61.7 ± 2.0 | -139.0 ± 5.0 |
| Dimethyl peroxide | 308.4 ± 3.0 | 80.7 ± 2.0 | -125.5 ± 5.0 |
| Diacetyl peroxide | 390.7 ± 6.0 | 122.3 ± 8.0 | -500.0 ± 10.0 |

12. Conclusions

The thermodynamic properties of 10 organic molecules have been calculated, based on the critical evaluation of available thermodynamic and spectroscopic information. Where no data were available, estimation techniques were used. Recommended values for the entropy, heat capacity, and the enthalpy of formation at 298.15 K are summarized in Table 11.

For a first round of new experimental studies, substantially improved formation properties could be obtained with new or confirming measurements for the enthalpy of formation for bromoacetic acid, glycolic acid, the cyanooxomethyl radical, and the three peroxides—methyl hydroperoxide, dimethyl peroxide, and diacetyl peroxide. In order to calculate significantly more reliable thermal functions, the structure and the vibrational frequencies are needed for the cyanooxomethyl radical and diacetyl peroxide, since these data were estimated for these two compounds.

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14. Extended Bibliographies

The following bibliography lists articles that were found in the literature pertaining to the molecules discussed above but were not used as sources of information in the evaluations.

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