

A Review, Bibliography, and Tabulation of K , L , and Higher Atomic Shell X-Ray Fluorescence Yields

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The measured K , L , and higher atomic shell x-ray fluorescence yield data, covering the period 1978 to 1993, following the major previous compilations by Bambynek *et al.* (1972) and Krause (1979), are reviewed. An annotated bibliography of x-ray fluorescence yield measurements, analyses, fits and tables 1978–1993 is presented. Comparisons of the fluorescence yields ω_K , ω_L , and ω_M , based on measurements, and on theoretical models, are presented. Values of ω_K , ω_L , and ω_M , fitted to standard empirical parametric formulations, are presented. In addition, selected well-characterized measured ω_K , ω_L , and ω_M results restricted to the period 1978–1993 are listed. These selected measured values are fitted by least squares to polynomials in Z of the form $\sum_n a_n Z^n$ and compared with theoretical and with earlier fitted values. A section on application of fluorescence yield data to computations of x-ray energy-absorption coefficients is included.

Key words: Auger effect; energy-absorption coefficient; fluorescence yield; ionization; photon; radiationless transitions; vacancies; x-ray.

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1. Introduction: Definitions, History

Values of the fluorescence yield $\omega_i(Z)$, where i represents a given atomic electron shell or subshell, of an atom with atomic number Z , are required in a variety of applications including, for example, atomic physics studies, x-ray fluorescence (XRF) surface chemical analysis, and dosimetric computations for health physics, cancer therapy, and industrial irradiation processing. The focus of this work is on fluorescence yield information for the latter application, in particular the computation of the photon energy-absorption coefficient μ_{en}/ρ .

An excellent description of the fluorescence yield has been given by Compton and Allison (35Co01), briefly paraphrased as follows. When an inner shell atomic electron is ejected from an atom as a result of a collision process involving a photon or other incoming projectile, a vacancy is thus created in that electron's pre-collision subshell.

For incoming photons, such processes could be the photoelectric effect, Compton collisions, and triplet production (electron-positron pair production in the field of the atomic electrons, by photons with energy in excess of 2.04 MeV). Associated with this inner-shell vacancy is an excess of energy above the atom's ground state.

Deexcitation of the (many-electron) atom then occurs through a cascade of radiative and (mostly) radiationless events and, unless electrons are captured from the outside, the atom ends up in a multiply ionized state as discussed, e.g., by Krause *et al.* (71Kr01), rather than returning to the ground state of the neutral atom. A statistical (probabilistic) analysis of this complex series of events has been attempted by Jacobs and Rozsnyai (86Ja01). However, in many applications, the atom is assumed to return to its ground state, including filling the vacancy, by one of two modes.

In one mode a fluorescence (or characteristic) x-ray is emitted^b from the atom, with photon energy equal to the difference between the vacancy-site inner-shell energy level and the energy level of the particular outer shell which happens to supply the electron to fill the vacancy. In the other mode, no fluorescence x-ray is emitted, but instead the excess energy is used to eject an outer-shell electron from the atom, in addition to the ejected electron causing the original vacancy. This second ejected electron is known as an Auger electron, after the cloud chamber confirmation and explanation by Auger (25Au01, 25Au02) of earlier conjectures by Sadler (09Sa01) and Barkla (17Ba01) and observations by Wilson (23Wi01, 23Wi02).

In simplest terms, the fluorescence yield $\omega_i(Z)$ is

$$\omega_i(Z) = \frac{f_i(Z)}{v_i(Z)} \quad (1)$$

^bHere and in what follows, we make the implicit assumption that excitation occurs well above threshold, so that excitation and deexcitation can be treated as distinct processes between which complete relaxation takes place. Also, the x-ray will not be characteristic, but a satellite, if the excitation process produces multiple vacancies through accompanying shake, which can happen with considerable probability [Åberg (67Åb01); Schaphorst (93Sc01) and references therein].

where $f_i(Z)$ is the average number of fluorescence (characteristic) x-rays emitted as a result of $v_i(Z)$ vacancies created in the i^{th} shell or subshell. In reality the situation is quite complex due to the multiplicity of transitions of varying likelihood involving the participating inner and outer subshells, the details of which are described in the extensive reviews by, e.g., Fink *et al.* (66Fi01, 74Fi01, 91Fi01), Bambynek *et al.* (72Ba01), Krause (79Kr01), Mitchell and Barfoot (81Mi01), and Cohen (87Co01). This work is an annex to, and not a replacement of, these earlier major reviews.

Compton and Allison (35Co01) and some later authors, e.g., Tertian and Claisse (82Te01) consider the Auger effect a two stage process, in which the characteristic x-ray is first emitted but then reabsorbed by an outer electron shell in a photoeffect type process. However, this two-stage picture is not in keeping with present-day understanding [Åberg and Howat (82Åb01)]. Already, Wentzel (27We01) showed that the Auger effect is explained by *Fermi's Golden Rule #2* (50Fe01) [actually Wentzel's, although ascribed by Fermi to Schiff (49Sc01)] of perturbation theory.

The energy that is released when the primary vacancy is filled by an electron from a higher orbit is transferred to the Auger electron by a virtual photon—there is no real intermediate state. This is consistent with the picture given by other authors such as Burhop (52Bu01, 72Bu01), and by N. A. Dyson (73Dy01) who contends that the Auger effect must be a one stage process because of strong lines (such as $K \rightarrow L_1$ transition) in the Auger spectra which are forbidden in radiative transitions. A unified theory of the Auger effect and autoionization as resonance scattering has been developed by Åberg and Howat (82Åb01).

The annotated bibliography in this work combines an automated search of Physics Abstracts with a manual search, and focuses on the time period 1978–1993 following the Krause (79Kr01) 1979 review article. Additional text references are listed separately following the annotated bibliography.

2. Application of Fluorescence Yield Data in Computing Mass Energy-absorption Coefficients

The mass energy-absorption coefficient μ_{en}/ρ (cm^2/g or m^2/kg , where $1 \text{ cm}^2/\text{g} = 0.1 \text{ m}^2/\text{kg}$), a key parameter in computations of energy deposited in media subjected to photon irradiation, is defined see, e.g., R. T. Berger (61Be01), Hubbell (77Hu01), Hubbell (82Hu01), Higgins *et al.* (92Hi01), and Seltzer (93Se01) as

$$\mu_{en}/\rho = (\mu/\rho)f \quad (2)$$

In this expression, f is the average fraction of the incident photon energy E_0 which does not leave the site of the primary collision in the form of secondary photon radiation (fluorescence, Compton scattered photons, annihilation radiation, bremsstrahlung, etc.) but goes into kinetic energy of particles, particularly electrons, for dissipation locally in the medium via collision losses as ionization and excitation. Here μ/ρ is the total mass attenuation coefficient see, e.g., Hubbell (69Hu01) for a given incident photon energy E_0 in a given substance:

$$\mu/\rho = \sigma_{\text{incoh}}/\rho + \sigma_{\text{coh}}/\rho + \tau/\rho + \kappa_n/\rho + \kappa_e/\rho + \sigma_{\text{ph.n.}}/\rho \quad (3)$$

in which the significant interaction components are incoherent (Compton) scattering $\sigma_{\text{incoh}}/\rho$, coherent (Rayleigh) scattering σ_{coh}/ρ , atomic photoeffect τ/ρ , electron-positron pair production in the field of the atomic nucleus κ_n/ρ , pair production in the field of the atomic electrons (triplet production) κ_e/ρ , and photonuclear interactions (e.g., (γ,n) , $(\gamma,2n)$, (γ,p) , (γ,fiss) , etc.) $\sigma_{\text{ph.n.}}/\rho$.

In Eqs. (2), (3), and (4) the units of μ/ρ and μ_{en}/ρ have been customarily cm^2/g , as well as all of the terms on the right-hand side of Eq. (3) (e.g., $\sigma_{\text{incoh}}/\rho$ in cm^2/g). However, the preferred units for these quantities (μ/ρ , μ_{en}/ρ etc.) are in the SI base units m^2/kg (where $1 \text{ m}^2/\text{kg} = 10 \text{ cm}^2/\text{g}$) as in the tables of Hubbell (82Hu01) and in ICRU Report 44 (89Ic01).

For computational purposes Eq. (2), using the information on the individual processes indicated in Eq. (3), becomes:

$$\mu_{\text{en}}/\rho = (\mu/\rho)f = (\sigma_{\text{incoh}}/\rho)f_{\text{incoh}} + (\tau/\rho)f_{\tau} + (\kappa_n/\rho)f_{\kappa_n} + (\kappa_e/\rho)f_{\kappa_e} + (\sigma_{\text{ph.n.}}/\rho)f_{\sigma_{\text{ph.n.}}} \quad (4)$$

in which the σ_{coh}/ρ term in Eq. (3) is omitted because the coherent scattering process deposits negligible energy at the collision site. Also, in existing μ_{en}/ρ tables, the photonuclear term $(\sigma_{\text{ph.n.}}/\rho)f_{\sigma_{\text{ph.n.}}}$ has been ignored. In Eq. (4) the energy-absorption weighting fractions f_{incoh} , f_{τ} , f_{κ_n} , and f_{κ_e} follow the f definition in Eq. (2), each for the individual interaction process indicated.

Although the incoherent scattering $\sigma_{\text{incoh}}/\rho$ and triplet production κ_e/ρ processes produce vacancies in the atomic electron subshells, resulting in the emission of either fluorescence photons or Auger electrons, existing tables ignore fluorescence emission (in effect, set $\omega_i = 0$) for these processes in which the subshell distribution of vacancy creation is not well known. An example of one of the few available vacancy cascade studies mentioning this effect (Compton ionization) is the work on Kr by Krause and Carlson (67Kr01). Computational procedures for evaluating f_{incoh} and f_{κ} (treating κ_n and κ_e alike), ignoring fluorescence, are given, e.g., by Hubbell (77Hu01) and by Higgins *et al.* (92Hi01). However, a treatment incorporating the full cascade is now available in newer and more detailed calculations by Seltzer (93Se01).

For the atomic photoeffect term $(\tau/\rho)f_{\tau}$, Hubbell (77Hu01) calculated f_{τ} using the formulation by R. T. Berger (61Be01), modified as suggested by Carlsson (71Ca01) to include additional fluorescence cascade effects. For each atomic electron shell group i , where i signifies K , L , M , ..., a photoeffect energy-absorption fraction f_{τ_i} was calculated for each incident photon energy E_0 and element Z as

$$f_{\tau_i}(E_0, Z) = \left(1 - \frac{\eta_i}{E_0}\right) \cdot [1 - G_{\text{br}}(E_0 - \eta_i)] + \frac{\eta_i}{E_0} \left(1 - \frac{\omega_i E_i}{\eta_i} - N_{i,i+1} \frac{\omega_{i+1} E_{i+1}}{\eta_i}\right) \quad (5)$$

in which η_i is the mean absorption-edge energy of the i -shell, $G_{\text{br}}(E_0 - \eta_i)$ is the bremsstrahlung yield see, e.g., Berger and Seltzer (83Be01) and ICRU (84Ic01) for the photoelectron of

energy $E_0 - \eta_i$, ω_i is the i -shell fluorescence yield, E_i is the i -shell fluorescence x-ray mean energy, and $N_{i,i+1}$ is the average number of vacancies created in the $(i+1)$ -shell per primary i -shell vacancy.

For μ_{en}/ρ in mixtures, where the bremsstrahlung yield $G_{\text{br}}(E_0 - \eta_i)$ is a function of the matrix rather than the atom suffering the primary collision, further refinements in the computations have been proposed by Attix (84At01), and have been implemented by Seltzer (93Se01).

Since the K fluorescence x-rays predominate over the L x-rays, and L x-rays over M x-rays, etc., both in photon energy and yield, particularly for low- Z elements, the approximation f_{τ} for the photoeffect energy-absorption fraction f_{τ} has sometimes been used, e.g., by R. T. Berger (61Be01), in computing μ_{en}/ρ :

For E_0 above the K edge:

$$(\tau/\rho)f_{\tau} \approx (\tau_K/\rho) \left(1 - \frac{\omega_K \bar{E}_K}{E_0}\right) + (\tau_L/\rho) + (\tau_M/\rho) + \dots \quad (6)$$

which reduces to

$$f_{\tau} \approx 1 - \frac{(\tau_K/\rho)}{(\tau/\rho)} \left(\frac{\omega_K \bar{E}_K}{E_0}\right). \quad (7)$$

For E_0 between the K and L edge:

$$(\tau/\rho)f_{\tau} \approx (\tau_L/\rho) \left(1 - \frac{\omega_L \bar{E}_L}{E_0}\right) + (\tau_M/\rho) + \dots \quad (8)$$

or

$$f_{\tau} \approx 1 - \frac{(\tau_L/\rho)}{(\tau/\rho)} \left(\frac{\omega_L \bar{E}_L}{E_0}\right). \quad (9)$$

and similarly between the L and M edges:

$$f_{\tau} \approx 1 - \frac{(\tau_M/\rho)}{(\tau/\rho)} \left(\frac{\omega_M \bar{E}_M}{E_0}\right) \quad (10)$$

with further ad hoc approximations within the L and M multiple-edge regions.

3. K-Shell Fluorescence Yield ω_K

The bulk of the fluorescence yield measurements reported in the literature have been for the K shell, and this trend has continued through the 1978–1993 period, as indicated in the annotated bibliography in Sec. 11 in this report. A semi-empirical fitting formula for ω_K , introduced by Burhop (55Bu01), has become established in the literature, of the form

$$\left(\frac{\omega_K}{1-\omega_K}\right)^{1/4} = C_0 + C_1 Z + C_2 Z^2 + C_3 Z^3 \quad (11)$$

$$= \sum_{i=0}^3 C_i Z^i$$

which can be rewritten

$$\omega_K = \frac{\left[\sum_{i=0}^3 C_i Z^i\right]^4}{1 + \left[\sum_{i=0}^3 C_i Z^i\right]^4}. \quad (12)$$

This same fitting formula, with different sets of C_i coefficients, has been applied to average L -shell fluorescence yields ω_L , as will be seen in Sec. 4 following.

Bambynek *et al.* (72Ba01), in their 1972 review article (the most comprehensive and widely quoted fluorescence yield reference to date), have fitted their collection of "selected 'most reliable' experimental values," listed in their Table III.IV with parameters for Eq. (12), above, of values:

$$\begin{aligned} C_0 &= 0.015 \pm 0.010 \\ C_1 &= 0.0327 \pm 0.0005 \\ C_2 &= 0 \\ C_3 &= -(0.64 \pm 0.07) \times 10^{-6}. \end{aligned} \quad (13)$$

In Table 1 of this work, ω_K values for $3 \leq Z \leq 110$, matching the range of the Krause (79Kr01) table, generated from the above Eqs. (12) and (13), are given in the first ω_K column, and are compared with measured values (averaged when more than one measurement per element is given) in the next column.

In a subsequent review, Krause (79Kr01) incorporated additional new data in a revised evaluation, and presented a table of ω_K adopted values for all elements $5 \leq Z \leq 110$, but did not provide a corresponding parametric fit or fits. These 1979 Krause adopted values are included in Table 1 of this work, also percent differences of these adopted values from the 1972 Bambynek *et al.* (72Ba01) fit.

In 1984 Bambynek (84Ba01) presented a further reevaluation of ω_K , incorporating about 100 new measurements subsequent to the 1972 Bambynek *et al.* (72Ba01) evaluation. Using a stepwise regression analysis with 119 selected ω_K measurements, Bambynek fitted his new evaluation to the form in Eq. (12) above, with parameters C_i now:

$$\begin{aligned} C_0 &= 0.0370 \pm 0.0052 \\ C_1 &= 0.03112 \pm 0.00044 \\ C_2 &= (5.44 \pm 0.11) \times 10^{-5} \\ C_3 &= -(1.250 \pm 0.070) \times 10^{-6}. \end{aligned} \quad (14)$$

This revised fit was used to generate the ω_K values in the next to the last column in Table 1 for all elements $3 \leq Z \leq 110$.

For $35 \leq Z \leq 107$ the 1984 values differ from the 1972 values by less than 1%. For $3 \leq Z \leq 34$ the new values exceed the old by more than 1%, and for $3 \leq Z \leq 11$ by more than 10%. However, in the range $3 \leq Z \leq 11$ ω_K is small, ranging from 2.928×10^{-4} for $Z = 3$ to 0.02133 for $Z = 11$.

Since ω_K enters the computations of photon energy-absorption coefficients μ_{en}/ρ approximately as $[1 - \omega_K(E_K/E_0)]$, see Eq. (6), where $E_K/E_0 < 1$, these differences, even the 78.5% at $Z = 3$, result in less than 1% differences in μ_{en}/ρ for all elements $3 \leq Z \leq 100$.

The fitting formula Eq. (12), rewritten from Eq. (11), with the 1984 Bambynek (84Ba01) parameters in Eq. (14) is the ω_K evaluation which was used in the Higgins *et al.* (92Hi01) and Seltzer (93Se01) μ_{en}/ρ computations, and the numerical values are repeated for convenience in the summary Table 8 for $3 \leq Z \leq 100$. However a new polynomial fit to weighted averages of more-recent well-characterized measured data for $11 \leq Z \leq 99$ is described in Sec. 7 and presented in Table 4

of this work, with coefficients given in Table 7, and is here recommended for future applications as discussed in Sec. 8.

4. L-Shell Average Fluorescence Yield ω_L

The status of fluorescence yield data for the L subshells L_1 , L_2 , and L_3 , and particularly for the average yield ω_L , is well summarized in the recent reviews by Cohen (87Co01) and by Jitschin (90Ji01). Jitschin summarizes and discusses measurements of ω_{L1} , ω_{L2} and ω_{L3} , also Coster-Kronig and Auger yields, over the time period 1980–1990. Jitschin classifies the measurement methods as (a) crystal spectrometer, (b) $K_\alpha - L_\alpha$ coincidence, and (c) synchrotron photoionization. In his summary table Jitschin (90Ji01) also includes two earlier papers by Campbell *et al.* (74Ca01, 75Mc02) which report $e^- - L$ -x-ray coincidence measurements in Pb of ω_{L1} and the Coster-Kronig yields f_{12} and f_{13} .

For the atomic photoeffect process, the average L -shell fluorescence yield ω_L is defined (see, for example Krause in 78Kr01, 79Kr01) as

$$\omega_L = (\sigma_{L1}\nu_1 + \sigma_{L2}\nu_2 + \sigma_{L3}\nu_3)/\sigma_L \quad (15)$$

in which σ_{L1} , σ_{L2} , σ_{L3} and σ_L are the L_1 , L_2 and L_3 individual L -subshell and total L -shell photoeffect cross sections (e.g., from Scofield 73Sc01). ν_1 , ν_2 and ν_3 are the effective subshell fluorescence yields which are discussed and defined in more detail by Krause in (80Kr01), who has also pointed out the generally weak dependence of the average L fluorescence yield on the initial vacancy distribution (93Kr01).

On the other hand, there are circumstances in which the L fluorescence yield can depend critically on the vacancy distribution among the three L subshells, which can have quite different effective fluorescence yields. Where Coster-Kronig channels are open, the vacancies in the more tightly bound subshells will predominantly bubble up to the $L_3(2p_{3/2})$ subshell before the next step in the cascade takes place. However, some Coster-Kronig transitions are energetically cut off in certain regions of the periodic table [Chen *et al.* (71Ch01), Crasemann *et al.* (71Cr01), Chen *et al.* (77Ch01)]. Then the vacancy distribution among the subshells can make a big difference, and it in turn depends critically on the mode of excitation. Thus, what is the average yield in one experiment can be very different from the average yield in another experiment if the excitation mechanisms (and energies) are different in the two. In this work, photoionization substantially above threshold, rather than, for example, collisions with charged particles, or decays of radioactive atoms (see, e.g., 75Ra01), is assumed to be the excitation mechanism.

Cohen (87Co01) presents, besides tables and standard-form fits (Eq. (12)) of the ω_L values given by Bambynek *et al.* (72Ba01) and by Mitchell and Barfoot (81Mi01), new theoretical values of ω_L calculated using the ECPSSR (*E*nergy loss, Coulomb deflection, Perturbed Stationary State, Relativistic effects) theory of Brandt and Lapicki (79Br01) and the resulting L -shell ionization cross sections computed by Cohen and Harrigan (85Co02).

Cohen's (87Co01) fitting parameters for ϖ_L in Eq. (12) are:
Bambynek *et al.* (72Ba01), ϖ_L , $23 \leq Z \leq 96$:

$$\begin{aligned} C_0 &= 0.238209 \\ C_1 &= -0.00216099 \\ C_2 &= 1.85156 \times 10^{-4} \\ C_3 &= -7.47647 \times 10^{-7} \end{aligned} \quad (16)$$

Mitchell and Barfoot (81Mi01), ϖ_L , $23 \leq Z \leq 92$:

$$\begin{aligned} C_0 &= 0.326968 \\ C_1 &= -0.00242879 \\ C_2 &= 1.71660 \times 10^{-4} \\ C_3 &= -6.96583 \times 10^{-7} \end{aligned} \quad (17)$$

Cohen (87Co01), ϖ_L , ECPSSR, $30 \leq Z \leq 96$:

$$\begin{aligned} C_0 &= 0.177650 \\ C_1 &= 0.00298937 \\ C_2 &= 8.91297 \times 10^{-5} \\ C_3 &= -2.67184 \times 10^{-7} \end{aligned} \quad (18)$$

Numerical values generated from these fits are compared with the tables presented by Cohen (87Co01), in Table 2 of this work.

The Cohen ECPSSR fit, using the above Eq. (18) parameters, was extended to lower Z 's, down to $Z = 3$, the first element with an L -shell electron, under the assumption that electrons external to the atom could participate in the fluorescence process. The ϖ_L values computed from the fit for the lowest Z 's appeared too high, exceeding the ω_K values. However, the Mitchell and Barfoot (81Mi01) table actually extends down to $Z = 12$, based on results in a thesis by Hoffmann (78Ho01).

A log-log plot of the Mitchell-Barfoot-Hoffmann values was found to be linear over the range $12 \leq Z \leq 50$ except for two anomalously low points for $Z = 17$ and 18 , perhaps due to round-off in the two-digit table. Using the slope of this log-log linear plot, and normalizing from 0.024 to the 3-digit 0.0242 ECPSSR Cohen-fit value at $Z = 37$ a fit is here presented:

For ϖ_L , $11 \leq Z \leq 36$:

$$\varpi_L = 1.9390 \times 10^{-8} \times Z^{3.8874} \quad (19)$$

In the summary Table 8 from (89Hu01), the values of ϖ_L were generated using Eq. (19) for $11 \leq Z \leq 36$, and Eq. (12) with the Cohen (87Co01) ECPSSR fit parameters in the above Eq. (18) for $37 \leq Z \leq 100$, and these ϖ_L values were used in the Higgins *et al.* (92Hi01) and Seltzer (93Se01) μ_{en}/ρ computations. However, a new polynomial fit to weighted averages of 1983–1993 evaluated measured data for $26 \leq Z \leq 92$ is described in Sec. 7 with numerical values presented in Table 5 of this work and coefficients given in Table 7, and is here recommended for future ϖ_L applications for $Z \geq 26$, as discussed in Sec. 8.

5. M-Shell Average Fluorescence Yield ϖ_M

An extended-range table of M -shell average fluorescence yield values ϖ_M , suitable for systematic computations of mass

energy-absorption coefficients μ_{en}/ρ , was not found in the literature, so this work has undertaken to provide recommended ϖ_M fits and tables.

Burhop (55Bu01) fitted the ϖ_M Lay (34La01) and Jaffe (54Ja01) data available in 1955 by a formula

$$\varpi_M = 1.7 \times 10^{-9} (Z-13)^4 \quad (20)$$

Subsequent measurements by Jopson *et al.* (65Jo01), corrected by Bambynek *et al.* (72Ba01) for a 20% correction from double M -shell vacancies to convert ω_{LM} into ϖ_M data, and by Konstantinov and Sazonova (68Ko01), by Hribar *et al.* (82Hr01), and by Shatendra *et al.* (84Sh01) are shown in Table 3. The average difference of all three measurements 1934–1984 from the Burhop (55Bu01) Eq. (20) values is a factor 0.758, with no significant trend as a function of Z over this limited high- Z range $76 \leq Z \leq 92$.

Hence a recommended fit has been proposed (89Hu01):

$$\begin{aligned} \varpi_M &= 0.758 \times 1.7 \times 10^{-9} (Z-13)^4 \\ &= 1.29 \times 10^{-9} (Z-13)^4 \end{aligned} \quad (21)$$

from which numerical values are given in the last column in Table 3, and for the range $19 \leq Z \leq 100$ in the summary Table 8.

Also shown in Table 3, for comparison, are some theoretical ϖ_M values derived from Chen *et al.* (80Ch02) and given by Sarkar *et al.* (81Sa01). The Chen *et al.* (80Ch02) values listed were obtained from their theoretical ω_{M_4} and ω_{M_5} values using a recipe quoted by McGuire (72Mc01) from Jopson *et al.* (65Jo01):

$$\varpi_M \approx \varpi_{LM} \approx 0.4 \omega_{M_4} + 0.6 \omega_{M_5} \quad (22)$$

The Sarkar *et al.* (81Sa01) values listed in Table 3 were also derived from the Chen *et al.* (80Ch02) theoretical ω_{M_i} subshell values, but by Sarkar *et al.* using the relation

$$\varpi_M = \sum_{i=1}^5 \frac{1}{18} N_{M_i} \omega_{M_i} \quad (23)$$

in which N_{M_i} are the numbers of electrons in each M_i subshell.

Equation (21), used in generating the ϖ_M values listed in the summary Table 8 for $19 \leq Z \leq 100$, could be further refined by including additional ω_{M_i} subshell measurements by Karttunen *et al.* (71Ka01) and by Baker *et al.* (74Ba01), perhaps combined using Eq. (23) above.

In the present work, a further ϖ_M analysis has been undertaken, incorporating more-recent measurements $64 \leq Z \leq 92$ up through 1993 as described in Sec. 7, with the results given in Table 6 including values generated from a new polynomial fit.

6. N- and Higher-Shell Average Fluorescence Yields

No N -shell or higher shell measurements were found in the literature. For the N shell, the best source of ϖ_N data, if

required, is probably the theoretical work of McGuire (74Mc01), which provides ω_{N_1} , ω_{N_2} , and ω_{N_3} values for 25 elements over the range $38 \leq Z \leq 103$, and ω_{N_4} , ω_{N_5} , and $\omega_{N_6,7}$ values for 20 elements over the range $50 \leq Z \leq 103$. For average $\bar{\omega}_N$ values, this information would need to be combined using an expression analogous to either Eq. (22) or (23) above.

As a cautionary note on the McGuire (74Mc01) results, Chantler (93Ch01) has pointed out to the authors a comment in a paper by Ohno and Wendum (85Oh01) that, although the McGuire results (75Mc01) have become a standard reference, an extensive comparison by Fugle and Alvarado (80Fu01) of experimental core-hole linewidths in the N shell with McGuire's calculations (74Mc01) revealed large differences between theory and experiment for Coster-Kronig type of transitions. However, good agreement is found over extensive ranges of the periodic table, hence the McGuire (74Mc01) results are here recommended for the N shell until replaced by improved systematic calculations and tabulations.

7. A Further Analysis and Fits Using 1978–1993 Measured Data

In this section we present a new analysis, using selected x-ray production (XRP) cross section and fluorescence yield measurements for the K -, L - and M -shells by both photons and charged particles from the period 1978–1993 (including also one 1977 paper). For those works in which the XRP cross sections were given, the fluorescence yields were obtained from the cross section values using the relation

$$\omega_{K/L/M} = \frac{\sigma_{K/L/M}}{\sigma_{K/L/M}} \quad (24)$$

where $\sigma_{K/L/M}$ denotes the total measured x-ray production cross section and $\sigma_{K/L/M}$ is the K -, L - or M -shell ionization cross section, respectively. We used the values of the ionization cross sections given by Scofield (73Sc01) and by Cohen (85Co02) for photons and protons, respectively.

$\sigma_{K/L/M}$ measured x-ray production cross sections were evaluated by adding the (K_α and K_β) XRP cross sections for K x-rays, (L_α , L_α , L_β and L_γ) XRP cross sections for the L x-rays, and in the case of M x-rays, since different components were not resolved, therefore the total M XRP cross sections were directly obtained. The fluorescence yield data for the K , L and M shells thus derived from XRP measurements according to Eq. (24) are presented in column 2 of the Tables 4, 5, and 6, respectively. The uncertainties shown are quoted from the originating authors.

In the cases where the XRP cross sections were reported at more than one incident energy of the exciting particle (photons, charged particles), the values of the i^{th} shell fluorescence yields, ω_i ($i = K, L, M$), for an element were obtained by taking the weighted average of the ω_j values available at different incident energies, using the expression

$$\omega_i = W \sum_j \left[\frac{\omega_j}{(\Delta\omega_j)^2} \right], \quad W = \frac{1}{\sum_j (\Delta\omega_j)^{-2}} \quad (25)$$

where ω_j denotes the j^{th} experimentally deduced fluorescence yield and $(\Delta\omega_j)$ represents the quoted uncertainty in the j^{th} experimental value. These weighted-average ω_i values are listed in column 3 of the Tables 4, 5, and 6.

The averaged (for each Z) experimental values of fluorescence yields were least squares fitted to polynomials in Z of the form

$$\omega_i = \sum_{n=0} \alpha_n Z^n . \quad (26)$$

Precaution was taken to avoid any undesirable oscillations in the regions between the calculated points. The values of the fitting coefficients a_n and the atomic-number ranges of validity of the polynomials are given in Table 7. The fitted values are presented in column 4 of the Tables 4, 5, and 6.

The theoretical average L shell fluorescence yields based on the RDHS (relativistic Dirac-Hartree-Slater) model, $\bar{\omega}_L(\text{RDHS})$, were taken from Puri *et al.* (93Pu04). In this reference the L_i subshell fluorescence (ω_i , $i = 1, 2, 3$) and Coster-Kronig (f_{ij}) yields have been evaluated using RDHS model based radiative and non-radiative transition rates from Scofield (74Sc01) and Chen *et al.* (79Ch02), respectively. The cutoffs and onsets of different Coster-Kronig transitions have been properly considered in these calculations. The $\bar{\omega}_L$ values were further evaluated using these calculated ω_i and f_{ij} values.

In the case of the M shell, theoretical data (72Mc01, 80Ch02, 83Ch01) on ω_i ($i = 1, 2, 3, 4, 5$), f_{ij} , and the super Coster-Kronig transition probabilities S_{ij} are available for a limited number of elements in the atomic-number region of interest. The values of these parameters for intermediate elements were obtained using spline interpolation. The interpolated values were used in the present calculations of theoretical average M shell fluorescence yields. As a word of caution, Krause (93Kr01) points out that there are inherent limitations in the single particle calculations for M and N shell parameters such as the early calculations by McGuire (72Mc01) and by Chen *et al.* (80Ch02, 83Ch01). These limitations are discussed by Karim and Crasemann (85Ka01) in their calculations of L -shell Coster-Kronig transition rates for Ar, including the effects of final-state channel mixing.

8. General Discussion

The values of K shell fluorescence yields ω_K generated from the polynomial fitting of the experimental data, using Eq. (26) and the K shell coefficients a_n in Table 7, are compared in Table 4 with the semi-empirical values of Krause (79Kr01) and with the values listed in the earlier report by Hubbell (89Hu01) generated from the 1984 Bambynek fit (84Ba01). The latter three data sets are seen to be in general agreement within the uncertainties indicated in the experimental-value column, which range from the order of 10% for $Z = 11$ to the order of 1% for $Z = 80$ and above.

The present fitted values of the average L shell fluorescence yields, $\bar{\omega}_L(\text{fit})$, are compared with the $\bar{\omega}_L(\text{RDHS})$ (93Pu04) and $\bar{\omega}_L(\text{ECPSSR})$ (87Co01) values, and also with the 1989 Hubbell listing (89Hu01), in Table 5. The $\bar{\omega}_L(\text{RDHS})$ (93Pu04) values are higher than the $\bar{\omega}_L(\text{ECPSSR})$ (87Co01)

by up to 12% in the atomic-number region $32 \leq Z \leq 77$ and up to 21% for elements with $Z < 32$. These two sets are in good agreement with each other for elements with $Z > 77$. It is clear from Table 5 that the $\omega_L(\text{fit})$ values are in good agreement with the $\omega_L(\text{RDHS})$ (93Pu04) values for the elements in the region $28 \leq Z \leq 92$. For the elements with $48 \leq Z \leq 77$, the fitted values are higher than the $\omega_L(\text{ECPSSR})$ (87Co01) by up to 20% whereas outside this atomic-number region the two sets are in agreement with each other. In our polynomial fitting in this work, we omitted the here-listed experimental ω_L values in the region $40 \leq Z \leq 53$ from Singh *et al.* (83Si02) because of the large deviations from the other available but limited data. As further measured data become available, this omission may be reconsidered in future fitting attempts.

The fitted average M shell fluorescence yields $\omega_M(\text{fit})$ and the two sets of theoretical values ω_M (Chen) (80Ch02, 83Ch01) and ω_M (McGuire) (72Mc01), also the 1989 Hubbell fit (89Hu01), are compared in Table 6. The two sets of theoretical values of ω_M agree with each other for the elements with $70 \leq Z \leq 92$ except for $Z = 90$ where the ω_M (Chen) is lower than the ω_M (McGuire) value by 20%. Above $Z = 90$, the theoretical data based on AHS ([non-relativistic] approximate Herman Skillman) calculations of McGuire (72Mc01) are not available.

Below $Z = 70$, RDHS model (80Ch02, 83Ch01) based data for the M shell on ω_i ($i = 1, 2, 3, 4, 5$), f_{ij} and S_{ij} are not available.

From Table 6 we note that the fitted values of ω_M for the elements in the region $64 \leq Z \leq 92$ agree with both the sets of theoretical values except for $Z = 90$, where the $\omega_M(\text{fit})$ value is closer to the ω_M (McGuire) value. The $\omega_M(\text{fit})$ values for the elements in the region $64 \leq Z \leq 92$ are found to be as much as 22% higher than the Burhop-formula (55Bu01) fitted 1989 values of Hubbell (89Hu01) which are included in this comparison since they have been used as input data to some practical applications including the recent μ_{en}/ρ calculations by Higgins *et al.* (92Hi01) and Seltzer (93Se01).

We note from the above discussion that:

(a) The present $\omega_K(\text{fit})$ values in Table 4 are in good agreement with the semi-empirical values of Krause (79Kr01) and are well supported by the RDHS model based values of Chen (80Ch01) over the atomic-number region $17 \leq Z \leq 80$. The 1989 Hubbell (89Hu01) ω_K values generally agree with the present fitted values within the stated experimental uncertainties, in many cases falling between the experimental values and the present fitted values.

(b) The present $\omega_L(\text{fit})$ values in Table 5 are best explained by the $\omega_L(\text{RDHS})$ (93Pu01) values over the atomic number region $28 \leq Z \leq 92$. The earlier Hubbell ω_L listings (89Hu01), generated from the Cohen (87Co01) standard-form parametric fit [Eq. (12), using Eq. (18) parameters, $Z = 37$ to 100, Eq. (19) for $Z = 3$ to 36] to his ECPSSR model results, deviate from the present polynomial fit in the region $46 \leq Z \leq 53$ by amounts ranging from 10% to 18% (for $Z = 52$); however, the listed independent measurements differ by similar amounts, as much as 24% for $Z = 49$, as discussed above.

(c) The present $\omega_M(\text{fit})$ values in Table 6 are seen to be in good agreement with both the Chen (80Ch02, 83Ch01) and the McGuire (72Mc01) theoretical values in the atomic number region $64 \leq Z \leq 92$. The 1989 Hubbell (89Hu01) ω_M values, although systematically lower than the present work for $Z \geq 70$, are seen to be in reasonable agreement for $Z < 70$, and could be used to extend the present fitted values down to $Z = 13$ in applications where this Z range is needed.

In view of these observations, the authors recommend the use of the polynomial fitted values [Eq. (26), using the coefficients given in Table 7] of the K -, L - and M -shell fluorescence yields for various applications. Nevertheless, the 1989 Hubbell (89Hu01) values in Table 8, along with their generating functions, are still included in this paper for applications where these extended ranges are required.

9. Tables 1 to 8: $\omega_K, \omega_L, \omega_M$, Analyses and FitsTABLE 1. ω_K : Comparison of earlier fits, tables, $3 \leq Z \leq 110$

Z	Bambynek <i>et al.</i> , 1972 fit ^a	Bambynek <i>et al.</i> , 1972 most rel. exper.	% Diffs. exper. from 1972 fit	Krause 1979 adopted table	% Diffs. 1979 adopted from 1972 fit	Bambynek 1984 fit ^a	% Diffs. 1984 fit to 1972 fit
3 Li	[1.64(-4)] ^a					[2.928(-4)] ^a	78.5%
4 Be	[4.51(-4)] ^a					[6.929(-4)] ^a	53.6
5 B	[.00101] ^a			.0017	68.3%	[.001409] ^a	39.5
6 C	[.00198] ^a			.0028	41.4	.002575	30.1
7 N	[.00351] ^a			.0052	48.1	.004349	23.9
8 O	[.00579] ^a			.0083	43.4	.006909	19.3
9 F	[.00902] ^a			.013	44.1	.01045	15.9
10 Ne	[.0134] ^a			.018	34.3	.01519	13.4
11 Na	[.0192] ^a			.023	19.8	.02133	11.1
12 Mg	[.0265] ^a			.030	13.3	.02911	9.8
13 Al	.0357	.0380	6.4%	.039	9.2	.03872	8.5
14 Si	.0469	.043	-8.3	.050	6.6	.05037	7.4
15 P	.0603	.060	-.5	.063	4.5	.06422	6.5
16 S	.0760	.082	7.9	.078	2.6	.08038	5.8
17 Cl	.0941	.0955	1.5	.097	3.1	.09892	5.1
18 Ar	.115	.122	6.1	.118	2.6	.1199	4.3
19 K	.138			.140	1.4	.1432	3.8
20 Ca	.163			.163	0	.1687	3.5
21 Sc	.190	.190	0	.188	-1.1	.1962	3.3
22 Ti	.219	.221	.9	.214	-2.3	.2256	3.0
23 V	.249	.253	1.6	.243	-2.4	.2564	3.0
24 Cr	.281	.283	.7	.275	-2.1	.2885	2.7
25 Mn	.314	.313	-.3	.308	-1.9	.3213	2.3
26 Fe	.347	.342	-1.4	.340	-2.0	.3546	2.2
27 Co	.381	.366	-3.9	.373	-2.1	.3880	1.8
28 Ni	.414			.406	-1.9	.4212	1.7
29 Cu	.446	.443	-.7	.440	-1.3	.4538	1.7
30 Zn	.479			.474	-1.0	.4857	1.4
31 Ga	.510	.528	3.5	.507	-.6	.5166	1.3
32 Ge	.540	.554	2.6	.535	-.9	.5464	1.2
33 As	.568	.589	3.6	.562	-1.1	.5748	1.2
34 Se	.596			.589	-1.2	.6019	1.0
35 Br	.622			.618	-.6	.6275	.9
36 Kr	.646	.660	2.2	.643	-.5	.6517	.9
37 Rb	.669	.669	0	.667	-.3	.6744	.8
38 Sr	.691	.702	1.6	.690	-.1	.6956	.7
39 Y	.711			.710	-.1	.7155	.6
40 Zr	.730			.730	0	.7340	.5
41 Nb	.747			.747	0	.7512	.6
42 Mo	.764			.765	.1	.7672	.4
43 Te	.779			.780	.1	.7821	.4
44 Ru	.793			.794	.1	.7958	.4
45 Rh	.806			.808	.2	.8086	.3
46 Pd	.818			.820	.2	.8204	.3
47 Ag	.830	.843	.5	.831	.1	.8313	.2
48 Cd	.840			.843	.4	.8415	.2
49 In	.850			.853	.4	.8508	.1
50 Sn	.859			.862	.3	.8595	.1
51 Sb	.867			.870	.3	.8676	.1
52 Te	.875	.857	-2.1	.877	.2	.8750	0
53 I	.882			.884	.2	.8819	0
54 Xe	.888	.894	.7	.891	.3	.8883	0
55 Cs	.895	.889	-.7	.897	.2	.8942	-.1
56 Ba	.900			.902	.2	.8997	0
57 La	.906			.907	.1	.9047	-.1
58 Ce	.911			.912	.1	.9096	-.2
59 Pr	.915			.917	.2	.9140	-.1
60 Nd	.920			.921	.1	.9181	-.2
61 Pm	.924			.925	.1	.9220	-.2
62 Sm	???			.929	.2	.9255	-.2

TABLE 1. ω_K : Comparison of earlier fits, tables, $3 \leq Z \leq 110$ — Continued

Z	Bambynek <i>et al.</i> , 1972 fit ^a	Bambynek <i>et al.</i> , 1972 most rel. exper.	% Diffs. exper. from 1972 fit	Krause 1979 adopted table	% Diffs. 1979 adopted from 1972 fit	Bambynek 1984 fit ^a	% Diffs. 1984 fit to 1972 fit
63 Eu	.931	.925	-.6	.932	.1	.9289	-.2
64 Gd	.934			.935	.1	.9320	-.2
65 Tb	.937			.938	.1	.9349	-.2
66 Dy	.940	.943	.3	.941	.1	.9376	-.2
67 Ho	.943			.944	.1	.9401	-.2
68 Er	.945			.947	.2	.9425	-.3
69 Tm	.947			.949	.2	.9447	-.2
70 Yb	.950			.951	.1	.9467	-.3
71 Lu	.952			.953	.1	.9487	-.3
72 Hf	.954			.955	.1	.9505	-.4
73 Ta	.956			.957	.1	.9522	-.4
74 W	.957			.958	.1	.9538	-.3
75 Re	.959			.959	0	.9553	-.4
76 Os	.960			.961	.1	.9567	-.3
77 Ir	.962			.962	0	.9580	-.4
78 Pt	.963	.967	.4	.963	0	.9592	-.4
79 Au	.964			.964	0	.9604	-.4
80 Hg	.966	.958	-.8	.965	-.1	.9615	-.5
81 Tl	.967			.966	-.1	.9625	-.5
82 Pb	.968	.972	.4	.967	-.1	.9634	-.5
83 Bi	.968			.968	0	.9643	-.4
84 Po	.970			.968	-.2	.9652	-.5
85 At	.971			.969	-.2	.9659	-.5
86 Rn	.972			.969	-.3	.9667	-.5
87 Fr	.972			.970	-.2	.9674	-.5
88 Ra	.973			.970	-.3	.9680	-.5
89 Ac	.974			.971	-.3	.9686	-.6
90 Th	.975			.971	-.4	.9691	-.6
91 Pa	.975			.972	-.3	.9696	-.6
92 U	.976	.970	-.6	.972	-.4	.9701	-.6
93 Np	[.977] ^a			.973	-.4	.9706	-.7
94 Pu	[.977] ^a			.973	-.4	.9710	-.6
95 Am	[.978] ^a			.974	-.4	.9713	-.7
96 Cm	[.978] ^a			.974	-.4	.9717	-.6
97 Bk	[.979] ^a			.975	-.4	.9720	-.7
98 Cf	[.979] ^a			.975	-.4	.9722	-.7
99 Es	[.980] ^a			.975	-.5	.9725	-.8
100 Fm	[.980] ^a			.976	-.4	[.9727] ^a	-.7
101 Md	[.980] ^a			.976	-.4	[.9729] ^a	-.7
102 No	[.981] ^a			.976	-.5	[.9730] ^a	-.8
103 Lw	[.981] ^a			.977	-.4	[.9732] ^a	-.8
104	[.981] ^a			.977	-.4	[.9732] ^a	-.8
105	[.982] ^a			.977	-.5	[.9733] ^a	-.9
106	[.982] ^a			.978	-.4	[.9733] ^a	-.9
107	[.982] ^a			.978	-.4	[.9734] ^a	-.9
108	[.983] ^a			.978	-.5	[.9733] ^a	-1.0
109	[.983] ^a			.978	-.5	[.9733] ^a	-1.0
110	[.983] ^a			.979	-.4	[.9732] ^a	-1.0

^aValues in square brackets [] are outside the regions of validity of the fits by Bambynek *et al.* (72Ba01) and by Bambynek (84Ba01). These unjustified extrapolations, using the fits, are included only to provide non-zero numerical data where required in some applications.

TABLE 2. Summary and analysis of Cohen (87Co01) ω_L review, fits, and tabulations, including ECPSSR (Energy loss, Coulomb deflection, Perturbed Stationary State, Relativistic effects)

Z	Bambynek et al., 1972 ω_L table	Bambynek et al., 1972 ω_L fit	% diff. fit from table	Mitchell & Barfoot 1981 ω_L table	Mitchell & Barfoot 1981 ω_L fit	% Diff. fit from table	Cohen 1987 ECPSSR ω_L table	Cohen 1987 ECPSSR ω_L fit	% Diff. fit from table
23 V	.00235	.00588	+150.2%	.0038	.0154	+305.%			
24 Cr									
25 Mn	.00295	.00685	132.2	.0052	.0170	+227.			
26 Fe									
27 Co									
28 Ni				.0081	.0201	+148.	.0092	.0111	+20.7%
29 Cu	.0056	.00951	69.8	.0093	.0213	129.	.0105	.0121	15.2
30 Zn				.011	.0226	105.	.0117	.0132	12.8
31 Ga	.0064	.0113	76.6	.012	.0240	100.	.0131	.0145	10.7
32 Ge				.014	.0255	82.1	.0145	.0158	9.0
33 As				.015	.0271	80.7	.0161	.0172	6.8
34 Se				.017	.0289	70.0	.0177	.0188	6.2
35 Br				.019	.0308	62.1	.0198	.0204	3.0
36 Kr				.022	.0329	49.5	.0219	.0222	1.4
37 Rb	.010	.0192	92.0	.024	.0351	46.3	.0241	.0242	.42
38 Sr				.027	.0375	38.9	.0262	.0263	.38
39 Y	.0315	.0229	-27.3	.030	.0401	33.7	.0288	.0285	-1.0
40 Zr				.033	.0428	29.7	.0313	.0309	-6.4
41 Nb				.036	.0458	27.2	.0344	.0335	-2.6
42 Mo				.040	.0489	22.3	.0374	.0363	-2.9
43 Tc				.043	.0522	21.4	.0406	.0393	-3.2
44 Ru				.047	.0558	18.7	.0438	.0425	-3.0
45 Rh				.052	.0596	14.6	.0471	.0459	-2.5
46 Pd				.056	.0637	13.8	.0503	.0495	-1.6
47 Ag	.0518	.0458	-11.6	.061	.0679	11.3	.0544	.0534	-1.8
48 Cd				.066	.0726	10.0	.0584	.0575	-1.5
49 In				.071	.0774	9.0	.0629	.0618	-1.7
50 Sn				.077	.0825	7.1	.0673	.0665	-1.2
51 Sb				.082	.0879	7.2	.0724	.0714	-1.4
52 Te				.089	.0936	5.2	.0774	.0765	-1.2
53 I				.096	.0996	3.8	.0828	.0820	-.97
54 Xe	.107	.0804	-24.9	.102	.106	3.9	.0882	.0877	-.57
55 Cs	.089	.0867	-2.6	.110	.113	2.7	.102	.0938	-8.04
56 Ba	.093	.0934	+4.43	.117	.119	1.7	.101	.100	-.99
57 La	.101	.101	0	.125	.127	1.6	.108	.107	-.93
58 Ce				.133	.134	.75	.115	.114	-.87
59 Pr	.123	.116	5.7	.141	.142	.71	.123	.121	-1.63
60 Nd	.131	.124	-5.3	.150	.150	0	.130	.129	-.77
61 Pm				.158	.159	.63	.138	.137	-.73
62 Sm				.168	.168	0	.145	.145	0
63 Eu	.142	.151	+6.3	.177	.177	0	.154	.153	-.65
64 Gd				.187	.187	0	.162	.163	-.62
65 Tb	.194	.172	-11.3	.197	.196	-.51	.172	.172	0
66 Dy	.14	.182	+30.0	.207	.206	-.48	.181	.182	-.55
67 Ho				.217	.217	0	.191	.192	-.52
68 Er				.228	.227	-.44	.201	.202	-.50
69 Tm				.239	.238	-.42	.210	.212	+0.95
70 Yb				.250	.249	-.40	.220	.223	1.36
71 Lu				.261	.261	0	.231	.234	1.30
72 Hf				.272	.272	0	.242	.245	1.24
73 Ta	.225	.266	+18.2	.284	.284	0	.255	.257	.78
74 W				.296	.296	0	.267	.269	.75
75 Re				.308	.308	0	.280	.281	.56
76 Os				.320	.320	0	.293	.293	0
77 Ir	.30	.320	6.7	.332	.332	0	.305	.305	0
78 Pt	.32	.334	4.4	.344	.344	0	.318	.318	0
79 Au	.398	.348	-12.6	.356	.357	+.28	.332	.331	-.30
80 Hg	.38	.362	-4.7	.369	.369	0	.345	.343	-.58
81 Tl	.43	.376	-12.6	.381	.382	+.26	.359	.356	-.84
82 Pb	.36	.390	8.3	.393	.394	.25	.372	.369	-.81
83 Bi	.40	.403	.7	.406	.406	0	.385	.382	-.78

TABLE 2. Summary and analysis of Cohen (87Co01) ω_L review, fits, and tabulations, including ECPSSR (Energy loss, Coulomb deflection, Perturbed Stationary State, Relativistic effects) — Continued

Z	Bambynek <i>et al.</i> , 1972 ω_L table	Bambynek <i>et al.</i> , 1972 ω_L fit	% diff. fit from table	Mitchell & Barfoot 1981 ω_L table	Mitchell & Barfoot 1981 ω_L fit	% Diff. fit from table	Cohen 1987 ECPSSR ω_L table	Cohen 1987 ECPSSR ω_L fit	% Diff. fit from table
84 Po							.398	.395	-.75
85 At							.411	.409	-.49
86 Rn							.423	.422	-.24
87 Fr							.436	.435	-.23
88 Ra	.451	.472	4.6				.448	.448	0
89 Ac							.461	.461	0
90 Th	.488	.498	2.0				.475	.474	-.21
91 Pa	.51	.511	.2				.487	.486	-.21
92 U	.51	.524	2.7	.515	.514	-.19	.499	.499	0
93 Np	.575	.537	-6.6				.510	.511	+.20
94 Pu	.581	.549	-5.5				.522	.524	.38
95 Am							.535	.536	19
96 Cm	.531	.572	+ 7.7				.547	.548	18
97 Bk							.560		
98 Cf							.572		
99 Es							.583		
100 Fm							.595		

TABLE 3. Earlier (89Hu01) analysis of ω_M data

Z	Measurements						Theory		Fit
	Lay 1934	Jaffe 1954	Jopson <i>et al.</i> 1965	Konstantinov & Sazonova 1968	Hribar <i>et al.</i> 1982	Shatendra <i>et al.</i> 1984	Chen <i>et al.</i> (McGuire recipe)	Sarkar <i>et al.</i> 1981 (derived Chen 1980)	
64 Gd								.0072	.0087
66 Dy								.0095	.0102
67 Ho								.0110	.0110
70 Yb							.0128	.0154	.0136
72 Hf								.0190	.0156
73 Ta								.0209	.0167
74 W							.0194	.0220	.0179
76 Os		.013±.003							.0203
78 Pt								.0257	.0230
79 Au		.024±.005		.023±.001		.025±.004		.0268	.0245
80 Hg							.0274		.0260
82 Pb		.026±.005		.029±.002	.032±.003	.028±.004			.0292
83 Bi	.037±.007	.030±.006		.035±.002			.0318		.0310
88 Ra							.0400		.0408
90 Th						.044±.004			.0453
92 U	.06					.051±.005	.0470		.0502
96 Cm							.0520		.0612
100 Fm							.0578		.0739

TABLE 4. K-shell fluorescence yields for the elements 11 ≤ Z ≤ 99

Z	$\omega_K(\text{Exp.})^a$	Ref.	Average values	This work Fitted values	Krause 79Kr01	89Hu01 84Ba01
11 Na	0.021±0.002	88Ra01	—	0.021	0.023	0.0213
12 Mg	0.027±0.003	88Ra01	—	0.026	0.030	0.0291
13 Al	0.034±0.003 0.027±0.005 0.030±0.003	88Ra01 81Ku01	0.033±0.039	0.0387	0.039	0.0387
14 Si	0.048±0.005 [0.0481±0.0014] ^a	88Ra01 87Br01	—	0.043	0.050	0.0504
16 S	0.070±0.008	88Ra01	—	0.071	0.078	0.0804
17 Cl	0.089±0.009 [0.101±0.004] ^a	88Ra01 78Es01	—	0.089	0.097	0.0989
19 K	0.134±0.020 0.131±0.003 [0.144±0.004]	88Ra01 81Bh01 93So01	0.132±0.003	0.132	0.140	0.143
20 Ca	0.151±0.003 0.156±0.005 0.127±0.013 [0.164±0.004] ^a	85Ga02 81Bh01 81Ku01 93So01	0.151±0.002	0.147	0.163	0.169
21 Sc	0.211±0.006	90Si02	—	0.183	0.188	0.196
22 Ti	0.205±0.005 0.216±0.008	85Ga02 81Bh01	0.208±0.004	0.218	0.214	0.226
23 V	0.249±0.006 0.252±0.020	87Ku01 81Ku01	0.249±0.005	0.253	0.243	0.256
24 Cr	0.281±0.006 [0.2901±0.0025] ^a	90Si02 78Ma01	—	0.286	0.275	0.289
25 Mn	0.321±0.007 0.310±0.023 [0.312±0.003] ^a [0.340±0.017] ^a [0.283±0.002] ^a [0.28 ± 0.02] ^a	87Ku01 81Ku01 89Ko01 88Ge01 82Sm01 81Gu01	0.320±0.007	0.319	0.308	0.321
26 Fe	0.336±0.006 0.335±0.011 0.350±0.019 [0.352±0.004] ^a	90Si02 81Bh01 92Pi01 92So01	0.336±0.005	0.351	0.340	0.355
27 Co	0.368±0.007	87Ku01	—	0.382	0.373	0.388
28 Ni	0.418±0.011 0.394±0.016	90Si02 81Ar02	0.410±0.009	0.412	0.406	0.421
29 Cu	0.448±0.010 0.441±0.018 0.440±0.018 0.425±0.021 [0.452±0.003] ^a	85Ga02 81Ar02 81Bh01 92Pi01 94So01	0.442±0.007	0.441	0.440	0.454
30 Zn	0.482±0.009 0.490±0.020 0.478±0.018 0.471±0.025	85Ga02 81Ar02 81Bh01 92Pi01	0.481±0.007	0.469	0.474	0.486
31 Ga	0.543±0.011	90Si02	—	0.496	0.507	0.517

TABLE 4. K-shell fluorescence yields for the elements 11 ≤ Z ≤ 99 — Continued

Z	ω_K (Exp.) ^a	Ref.	Average values	This work Fitted values	Krause 79Kr01	89Hu01 84Ba01
32 Ge	0.529±0.010 0.538±0.029 0.549±0.011 [0.532±0.016] ^a	85Ga02 92Pi01 84Ca01 87Br01	0.539±0.009	0.523	0.535	0.546
33 As	0.590±0.024 0.579±0.025 0.574±0.029	81Ar02 88Si01 84Si01	0.581±0.017	0.549	0.562	0.575
34 Se	0.591±0.011 0.537±0.022	85Ga02 81Ar02	0.580±0.009	0.574	0.589	0.602
35 Br	0.586±0.011 0.586±0.023 [0.626±0.012] ^a	87Ku01 81Ar02 78Eg01	0.586±0.009	0.598	0.618	0.628
36 Kr	0.660±0.007	86Ko01	—	0.621	0.643	0.652
37 Rb	0.635±0.013 [0.673±0.008] ^a	87Ku01 78Th01	—	0.643	0.667	0.674
38 Sr	0.687±0.023 0.697±0.034	90Si02 81Bh01	0.690±0.02	0.665	0.690	0.696
39 Y	0.668±0.030	85Ga02	—	0.685	0.710	0.716
40 Zr	0.700±0.028 0.725±0.039	81Ar02 81Bh01	0.708±0.23	0.705	0.730	0.734
41 Nb	0.722±0.044 0.738±0.030	90Si02 81Ar02	0.733±0.25	0.724	0.747	0.751
42 Mo	0.746±0.041 0.740±0.01 0.804±0.032 0.758±0.043 0.792±0.013	87Al01 87Ku01 81Ar02 92Pi01 91Ca01	0.773±0.009	0.742	0.765	0.767
45 Rh	0.829±0.058	90Si02	—	0.792	0.808	0.809
46 Pd	0.846±0.059	90Si02	—	0.807	0.820	0.820
47 Ag	0.843±0.046 0.856±0.025 0.857±0.034 0.861±0.072 0.826±0.005 [0.84±0.02] ^a [0.847±0.013] ^a	87Al01 85Ga02 81Ar02 81Bh01 80Ta02 79Pi01 89Eg01	0.828±0.004	0.822	0.831	0.831
48 Cd	0.874±0.048 0.856±0.075 0.853±0.013	87Al01 81Bh01 91Ca01	0.854±0.012	0.836	0.843	0.842
49 In	0.900±0.049 0.864±0.025 0.849±0.036 0.843±0.005	87Al01 87Ku01 81Ar02 80Ta02	0.844±0.004	0.848	0.853	0.851
50 Sn	0.854±0.047 0.872±0.025 0.890±0.036 0.874±0.013	87Al01 85Ga02 81Ar02 91Ca01	0.874±0.011	0.861	0.862	0.860

TABLE 4. K-shell fluorescence yields for the elements 11 $\leq Z \leq 99$ — Continued

Z	ω_K (Exp.) ^a	Ref.	Average values	This work Fitted values	Krause 79Kr01	89Hu01 84Ba01
51 Sb	0.896 \pm 0.049 0.866 \pm 0.025	87Al01 87Ku01	0.872 \pm 0.022	0.872	0.870	0.868
52 Te	0.823 \pm 0.073	90Si02	—	0.883	0.877	0.875
53 I	0.846 \pm 0.024 0.831 \pm 0.033	87Ku01 81Ar02	0.841 \pm 0.019	0.894	0.884	0.882
54 Xe	0.889 \pm 0.010	77Hr01	—	0.903	0.891	0.888
55 Cs	0.902 \pm 0.026 0.899 \pm 0.015 0.896 \pm 0.016 [0.904 \pm 0.018] ^a	87Ku01 88Si01 83Si01 89Eg01	0.898 \pm 0.013	0.912	0.897	0.894
56 Ba	0.920 \pm 0.051 0.934 \pm 0.027	87Al01 85Ga02	0.931 \pm 0.024	0.920	0.902	0.900
57 La	0.913 \pm 0.050	87Al01	—	0.928	0.907	0.905
59 Pr	0.930 \pm 0.023	88Si01	—	0.941	0.917	0.914
60 Nd	[0.917 \pm 0.020] ^a	79Ch02	—			0.918
61 Pm	[0.918 \pm 0.017] ^a	89Eg01	—			0.922
62 Sm	[0.913 \pm 0.028] ^a	85Se02	—			0.923
63 Eu	0.957 \pm 0.030 0.933 \pm 0.019	88Si01 85Si01	0.939 \pm 0.16	0.962	0.932	0.929
66 Dy	0.975 \pm 0.027	88Si01	—	0.972	0.941	0.938
69 Tm	0.983 \pm 0.028	88Si01	—	0.979	0.949	0.945
70 Yb	[0.954 \pm 0.023] ^a	89Eg01	—			0.947
71 Lu	0.951 \pm 0.030	88Si01	—	0.981	0.953	0.949
73 Ta	0.955 \pm 0.011	88Si01	—	0.983	0.957	0.952
80 Hg	0.980 \pm 0.009	88Si01	—	0.980	0.965	0.962
93 Np	0.972 \pm 0.003	79Ah01	—	0.969	0.973	0.971
94 Pu	0.972 \pm 0.003	79Ah01	—	0.969	0.973	0.971
96 Cm	0.971 \pm 0.006	79Ah01	—	0.971	0.974	0.972
97 Bk	[0.971 \pm 0.006] ^a	79Ah01	—			0.972
98 Cf	0.973 \pm 0.004 [0.976 \pm 0.005] ^a	79Ah01 77Fr01	—	0.974	0.975	0.972
99 Es	0.972 \pm 0.004	79Ah01	—	0.976	0.975	0.973

^aValues in square brackets [] have been added in proof at the suggestion of W. Bambynek, and have not been included in the average values in column 4 nor in the fitting to obtain column 5.

TABLE 5. Average *L*-shell fluorescence yields for the elements 26 $\leq Z \leq 92$

Average L-Shell fluorescence yields ω_L							
Z	Experimental values	Ref.	Average values	This work fitted values	RHDS 93Pu04	ECPSSR 87Co01	Hubbell 89Hu01
26 Fe	0.0063 \pm 0.0010	91Mc01	—	0.0064	0.0052	0.0063	0.00614
28 Ni	0.0091 \pm 0.0014	91Mc01	0.0087 \pm 0.001	0.0088	0.0085	0.0092	0.00819
	0.0083 \pm 0.0016	85Du01					
29 Cu	0.0105 \pm 0.001	91Mc01	0.0102 \pm 0.0012	0.0100	0.0100	0.0105	0.00939
	6.0098 \pm 0.0019	85Du01					
30 Zn	0.0117 \pm 0.0018	91Mc01	—	0.0113	0.0103	0.0117	0.0107
31 Ga	0.0129 \pm 0.0019	91Mc01	—	0.0128	0.0121	0.0131	0.0122
32 Ge	0.0139 \pm 0.0021	91Mc01	0.0140 \pm 0.0016	0.0141	0.0140	0.0145	0.0138
33 As	0.0156 \pm 0.0023	85Du01	—	0.0156	0.0160	0.0161	0.0155
36 Kr	0.0210 \pm 0.002	79Sp01	—	0.0211	0.0209	0.0219	0.0218
37 Rb	0.0186 \pm 0.0028	85Du01	—	0.0232	0.0234	0.0241	0.0242
38 Sr	0.0213 \pm 0.0032	85Du01	—	0.0256	0.0260	0.0262	0.0263
39 Y	0.0246 \pm 0.0036	80Se01	0.0245 \pm 0.0024	0.0282	0.0288	0.0288	0.0285
40 Zr	0.0330 \pm 0.0049	85Du01	—	0.0310	0.0318	0.0313	0.0319
	0.0282 \pm 0.0014	83Si02					
41 Nb	0.037 \pm 0.003	92Ga01	—	0.0342	0.0361	0.0344	0.0335
	0.029 \pm 0.0014	83Si02					
42 Mo	0.0380 \pm 0.003	92Ga01	0.0380 \pm 0.0029	0.0376	0.0396	0.0374	0.0363
	0.0316 \pm 0.0016	83Si02					
45 Rh	0.051 \pm 0.005	92Ga01	—	0.0499	0.0517	0.0471	0.0459
46 Pd	0.054 \pm 0.005	92Ga01	0.0498 \pm 0.0040	0.0547	0.0557	0.0503	0.0495
	0.039 \pm 0.007	85Du01					
47 Ag	0.057 \pm 0.005	92Ga01	—	0.0599	0.0599	0.0544	0.0534
	0.0556 \pm 0.002	883Si02					
48 Cd	0.066 \pm 0.005	92Ga01	—	0.0656	0.0652	0.0584	0.0575
	0.0569 \pm 0.002	883Si02					
49 In	0.075 \pm 0.005	92Ga01	—	0.0717	0.0705	0.0629	0.0618
	0.0571 \pm 0.0029	83Si02					
50 Sn	0.079 \pm 0.006	92Ga01	—	0.0782	0.0757	0.0673	0.0665
	0.081 \pm 0.012	80Se01					
51 Sb	0.083 \pm 0.006	92Ga01	—	0.0852	0.0813	0.0724	0.0714
52 Te	0.093 \pm 0.007	92Ga01	—	0.0934	0.0873	0.0774	0.0765
53 I	0.077 \pm 0.004	83Si02	—	0.096	0.092	0.083	0.082
56 Ba	0.110 \pm 0.003	90Si01	—	0.110	0.114	0.101	0.100
57 La	0.118 \pm 0.003	90Si01	0.117 \pm 0.003	0.116	0.121	0.108	0.107
	0.108 \pm 0.008	90Ma02					
58 Ce	0.121 \pm 0.004	90Si01	0.120 \pm 0.004	0.123	0.129	0.115	0.114
	0.108 \pm 0.008	90Ma02					

TABLE 5. Average *L*-shell fluorescence yields for the elements 26 $\leq Z \leq 92$ — Continued

Z	Average L-Shell fluorescence yields w_L						
	Experimental values	Ref.	Average values	This work fitted values	RHDS 93Pu04	ECPSSR 87Co01	Hubbell 89Hu01
59 Pr	0.132±0.004 0.127±0.009	90Si01 90Ma02	0.131±0.003	0.130	0.138	0.123	0.121
60 Nd	0.143±0.004 0.131±0.009	90Si01 90Ma02	0.141±0.003	0.138	0.146	0.130	0.129
62 Sm	0.161±0.005 0.149±0.010 0.144±0.005	90Si01 90Ma02 92St01	0.152±0.003	0.155	0.164	0.145	0.145
63 Eu	0.164±0.005 0.148±0.010	90Si01 90Ma02	0.161±0.004	0.165	0.173	0.154	0.153
64 Gd	0.184±0.005 0.165±0.010	90Si01 90Ma02	0.180±0.004	0.174	0.184	0.162	0.163
65 Tb	0.192±0.006 0.168±0.010	90Si01 90Ma02	0.186±0.005	0.184	0.194	0.172	0.172
66 Dy	0.199±0.006 0.175±0.010	90Si01 90Ma02	0.192±0.005	0.194	0.204	0.181	0.182
67 Ho	0.217±0.006 0.193±0.010 0.267±0.010	90Si01 90Ma02 86Bh01	0.222±0.004	0.205	0.214	0.191	0.192
68 Er	0.223±0.007 0.205±0.010	90Si01 90Ma02	0.217±0.006	0.215	0.223	0.201	0.202
69 Tm	0.228±0.007	90Si01	—	0.226	0.231	0.210	0.212
70 Yb	0.239±0.007 0.228±0.010	90Si01 90Ma02	0.235±0.006	0.236	0.241	0.220	0.223
71 Lu	0.246±0.007 0.235±0.010	90Si01 90Ma02	0.242±0.006	0.247	0.252	0.231	0.234
72 Hf	0.255±0.007	90Si01	—	0.258	0.264	0.242	0.245
73 Ta	0.274±0.008 0.254±0.012 20.273±0.008 0.316±0.013 0.280±0.020	90Si01 90Ma0 85Sh01 86Bh01 85Si02	0.267±0.005	0.269	0.277	0.255	0.257
74 W	0.285±0.008 0.272±0.013 0.296±0.021 0.290±0.020	90Si01 90Ma02 85Sh01 85Si02	0.283±0.006	0.280	0.290	0.267	0.269
75 Re	0.286±0.008	90Si01	—	0.292	0.301	0.280	0.281
77 Ir	0.326±0.010	90Si01	—	0.314	0.322	0.305	0.305
78 Pt	0.328±0.010	90Si01	—	0.326	0.332	0.318	0.318
79 Au	0.330±0.010 0.338±0.016 0.336±0.023 0.345±0.014 0.360±0.020	90Si01 90Ma02 85Sh01 86Bh01 85Si02	0.338±0.007	0.337	0.342	0.332	0.331

TABLE 5. Average *L*-shell fluorescence yields for the elements 26 $\leq Z \leq 92$ — Continued

Average L-Shell fluorescence yields w_L							
Z	Experimental values	Ref.	Average values	This work fitted values	RHDS 93Pu04	ECPSSR 87Co01	Hubbell 89Hu01
80 Hg	0.346±0.017	90Ma02	0.349±0.010	0.348	0.352	0.345	0.343
	0.323±0.020	85Sh01					
	0.380±0.020	85Si02					
81 Tl	0.354±0.010	90Si01	0.356±0.007	0.360	0.363	0.359	0.356
	0.349±0.017	90Ma02					
	0.337±0.023	85Sh01					
	0.365±0.015	86Bh01					
	0.390±0.030	85Si02					
82 Pb	0.374±0.010	790Si01	0.376±0.008	0.371	0.374	0.372	0.369
	0.361±0.018	90Ma02					
	0.391±0.027	85Sh01					
	0.395±0.019	86Bh01					
	0.380±0.030	85Si02					
	0.401±0.000	81Ko01					
83 Bi	0.374±0.010	90Si01	0.328±0.007	0.383	0.385	0.385	0.382
	0.367±0.017	90Ma02					
	0.410±0.023	85Sh01					
	0.411±0.015	86Bh01					
90 Th	0.473±0.010	90Si01	0.464±0.018	0.468	0.470	0.475	0.474
	0.407±0.017	85Sh01					
	0.456±0.023	86Bh01					
	0.490±0.015	85Si02					
92 U	0.489±0.010	90Si01	0.500±0.019	0.495	0.492	0.499	0.499
	0.609±0.042	85Sh01					
	0.492±0.025	86Bh01					
	0.600±0.040	85Si02					

TABLE 6. Average M-shell fluorescence yields for the elements $71 \leq Z \leq 92$

Z	Average M-Shell fluorescence yields ω_M						
	Experimental values	Ref.	Average values	This work fitted values	Chen 80Ch02 83Ch01	McGuire 72Mc01	Hubbell 89Hu01
71 Lu	0.0154±0.0015	93Pu01		-0.0156	0.0172	0.0160	0.0146
72 Hf	0.0176±0.0017	93Pu01		0.0173	0.0183	0.0186	0.0156
73 Ta	0.0190±0.0019	93Pu01		0.0189	0.0193	0.0208	0.0167
77 Ir	0.0276±0.0022	93Pu01		0.0257	0.0240	0.0236	0.0216
78 Pt	0.0285±0.0023	93Pu01		0.0274	0.0254	0.0247	0.0230
79 Au	0.0264±0.0021 0.0300±0.0024	90Ma01 93Pu01	0.0279±0.0015	0.0292	0.0268	0.0270	0.0245
81 Tl	0.0332±0.0020	91Ga01	-	0.0328	0.0298	0.0305	0.0275
82 Pb	0.0362±0.0024 0.0311±0.0025 0.0334±0.0027	91Ga01 90Ma01 93Pu01	0.0336±0.0014	0.0346	0.0313	0.0320	0.0292
83 Bi	0.0384±0.0020 0.0356±0.0025	91Ga01 93Pu01	0.0373±0.0015	0.0365	0.329	0.0334	0.0310
90 Th	0.0525±0.0036 0.0537±0.0037 0.0512±0.0035	91Ga01 90Ma01 93Pu01	0.0524±0.002	0.0501	0.0451	0.0543	0.0453
92 U	0.0539±0.0037 0.0535±0.003 0.0514±0.0031	91Ga01 790Ma01 93Pu01	0.0527±0.0019	0.0541	0.0491	-	0.0502

TABLE 7. K-, L-, and M-shell fluorescence yields fitted to the polynomials $\sum_n a_n Z^n$ as a function of atomic number (Z)

Parameter	Range of Z	Fitting coefficient				
		a_0	a_1	a_2	a_3	a_4
ω_K	11-19	1.4340×10^{-1}	-2.5606×10^{-2}	1.3163×10^{-3}	-	-
	20-99	-7.6388×10^{-1}	5.4070×10^{-2}	-4.0544×10^{-4}	-1.4348×10^{-6}	-1.8252×10^{-8}
ω_L	26-51	-9.2521×10^{-2}	8.7531×10^{-3}	-2.8087×10^{-4}	3.4823×10^{-6}	-
	52-92	4.2193	-2.3520×10^{-1}	4.7911×10^{-3}	-4.1549×10^{-5}	-1.3564×10^{-7}
ω_M	71-92	-4.587×10^{-2}	1.208×10^{-4}	1.051×10^{-5}	-	-

TABLE 8. Fluorescence yield values ω_K , ω_L , and ω_M for $3 \leq Z \leq 100$, generated from the Bambynek (84Ba01), Cohen (87Co01) and Burhop (55Bu01) fitting functions extended and modified (in part) by Hubbell (89Hu01)

Z	ω_K	ω_L	ω_M
3 Li	2.928(-4)		
4 Be	6.929(-4)		
5 B	.001409		
6 C	.002575		
7 N	.004349		
8 O	.006909		
9 F	.01045		
10 Ne	.01519		
11 Na	.02133	2.17(-4)	
12 Mg	.02911	3.04(-4)	
13 Al	.03872	4.15(-4)	
14 Si	.05037	5.53(-4)	
15 P	.06422	7.24(-4)	
16 S	.08038	9.30(-4)	
17 Cl	.09892	.00118	
18 Ar	.1199	.00147	
19 K	.1432	.00181	1.67(-6)
20 Ca	.1687	.00221	3.10(-6)
21 Sc	.1962	.00268	5.28(-6)
22 Ti	.2256	.00321	8.46(-6)
23 V	.2564	.00381	1.29(-5)
24 Cr	.2885	.00450	1.89(-5)
25 Mn	.3213	.00527	2.67(-5)
26 Fe	.3546	.00614	3.68(-5)
27 Co	.3880	.00711	4.96(-5)
28 Ni	.4212	.00819	6.53(-5)
29 Cu	.4538	.00939	8.45(-5)
30 Zn	.4857	.0107	1.08(-4)
31 Ga	.5166	.0122	1.35(-4)
32 Ge	.5464	.0138	1.68(-4)
33 As	.5748	.0155	2.06(-4)
34 Se	.6019	.0174	2.51(-4)
35 Br	.6275	.0195	3.02(-4)
36 Kr	.6517	.0218	3.61(-4)
37 Rb	.6744	.0242	4.28(-4)
38 Sr	.6956	.0263	5.04(-4)
39 Y	.7155	.0285	5.90(-4)
40 Zr	.7340	.0309	6.86(-4)
41 Nb	.7512	.0335	7.93(-4)
42 Mo	.7672	.0363	9.12(-4)
43 Tc	.7821	.0393	.00104
44 Ru	.7958	.0425	.00119
45 Rh	.8086	.0459	.00135
46 Pd	.8204	.0495	.00153
47 Ag	.8313	.0534	.00172
48 Cd	.8415	.0575	.00193
49 In	.8508	.0618	.00217
50 Sn	.8595	.0665	.00242
51 Sb	.8676	.0714	.00269

TABLE 8. Fluorescence yield values ω_K , ω_L , and ω_M for $3 \leq Z \leq 100$, generated from the Bambynek (84Ba01), Cohen (87Co01) and Burhop (55Bu01) fitting functions extended and modified (in part) by Hubbell (89Hu01) — Continued

Z	ω_K	ω_L	ω_M
52 Te	.8750	.0765	.00298
53 I	.8819	.0820	.00330
54 Xe	.8883	.0877	.00365
55 Cs	.8942	.0938	.00401
56 Ba	.8997	.100	.00441
57 La	.9049	.107	.00484
58 Ce	.9096	.114	.00529
59 Pr	.9140	.121	.00578
60 Nd	.9181	.129	.00629
61 Pm	.9220	.137	.00685
62 Sm	.9255	.145	.00744
63 Eu	.9289	.153	.00806
64 Gd	.9320	.163	.00873
65 Tb	.9349	.172	.00943
66 Dy	.9376	.182	.0102
67 Ho	.9401	.192	.0110
68 Er	.9425	.202	.0118
69 Tm	.9447	.212	.0127
70 Yb	.9467	.223	.0136
71 Lu	.9487	.234	.0146
72 Hf	.9505	.245	.0156
73 Ta	.9522	.257	.0167
74 W	.9538	.269	.0179
75 Re	.9553	.281	.0191
76 Os	.9567	.293	.0203
77 Ir	.9580	.305	.0216
78 Pt	.9592	.318	.0230
79 Au	.9604	.331	.0245
80 Hg	.9615	.343	.0260
81 Tl	.9625	.356	.0275
82 Pb	.9634	.369	.0292
83 Bi	.9643	.382	.0310
84 Po	.9652	.395	.0328
85 At	.9659	.409	.0347
86 Rn	.9667	.422	.0366
87 Fr	.9674	.435	.0387
88 Ra	.9680	.448	.0408
89 Ac	.9686	.461	.0430
90 Th	.9691	.474	.0453
91 Pa	.9696	.486	.0477
92 U	.9701	.499	.0502
93 Np	.9706	.511	.0528
94 Pu	.9710	.524	.0555
95 Am	.9713	.536	.0583
96 Cm	.9717	.548	.0612
97 Bk	.9720	.560	.0642
98 Cf	.9722	.572	.0673
99 Es	.9725	.583	.0706
100 Fm	.9727	.595	.0739

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