

# Probing optical and electronic properties of individual defects through scanning transmission electron microscopy and first-principles theory

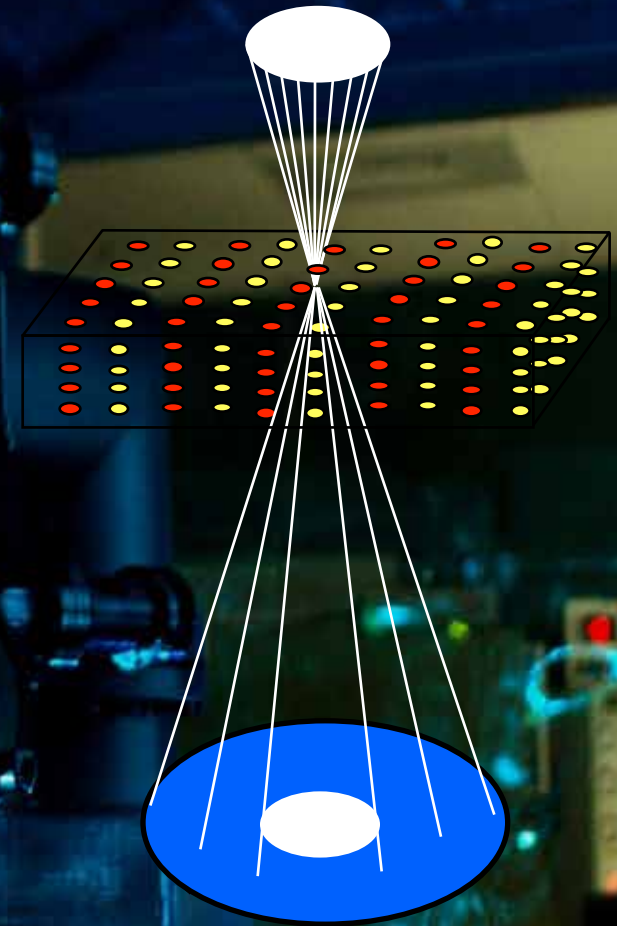
Stephen J. Pennycook

Departments of Materials Science and  
Engineering,

National University of Singapore

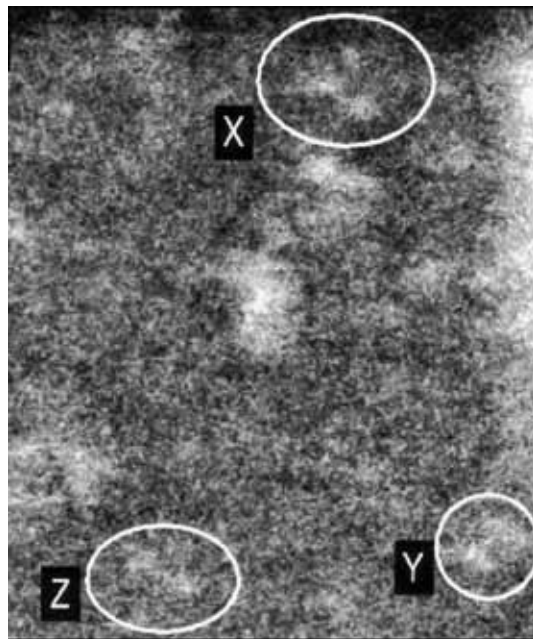
and

University of Tennessee



# Advances in aberration correction

## Pt on $\gamma\text{-Al}_2\text{O}_3$

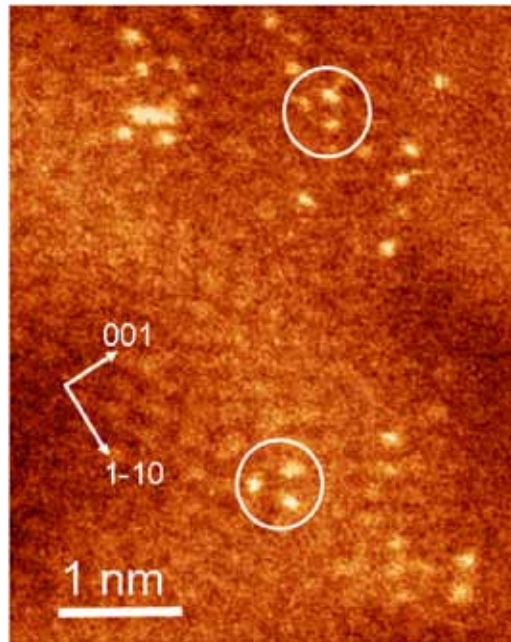


1 nm

**P.D. Nellist 1996**

**Uncorrected**

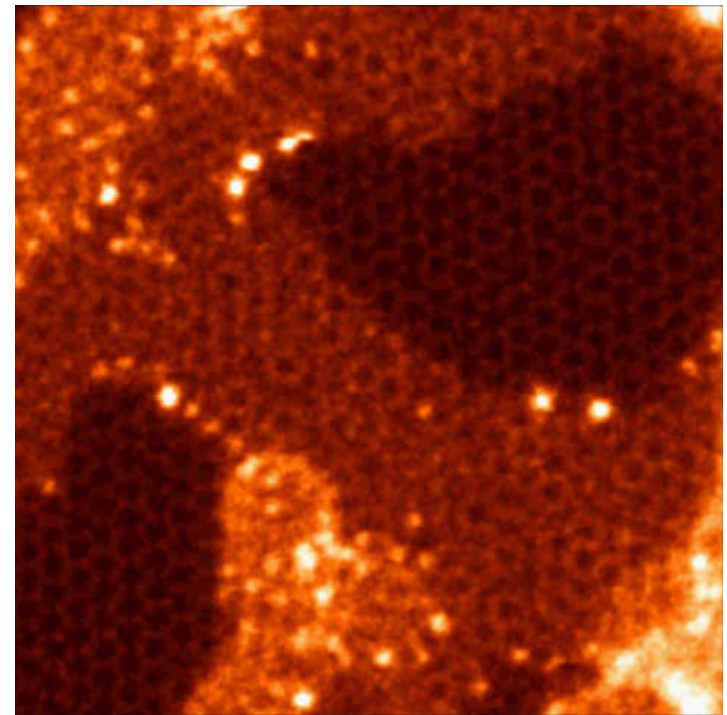
VG Microscopes HB603U 300 kV



**A. Y. Borisevich  
2004**

**3<sup>rd</sup> order  
corrected**

## Pt on graphene

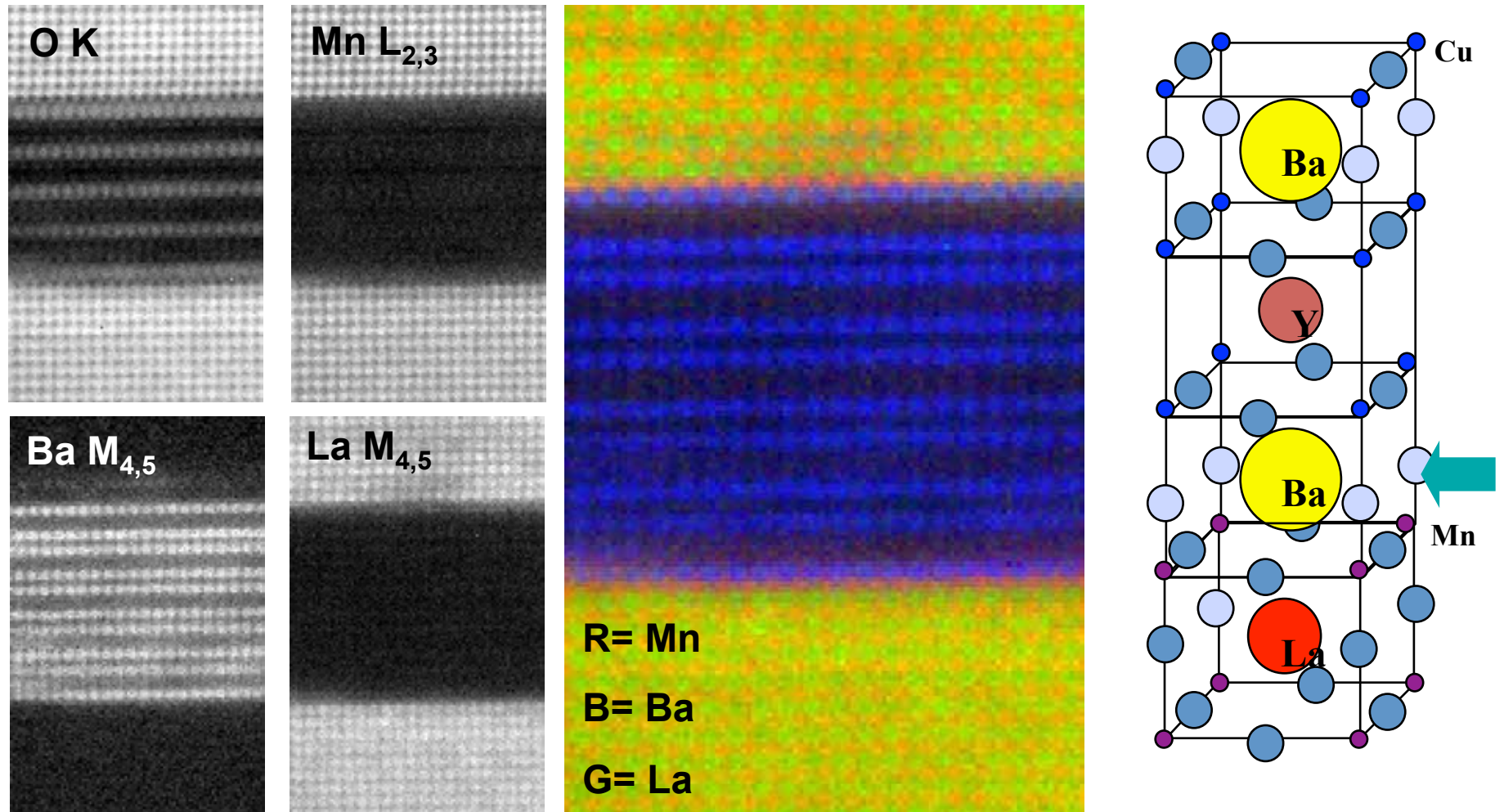


**Wu Zhou 2012**

**5<sup>th</sup> order  
corrected**

Nion UltraSTEM 60 kV

# Stacking Sequence at Interfaces by STEM/EELS

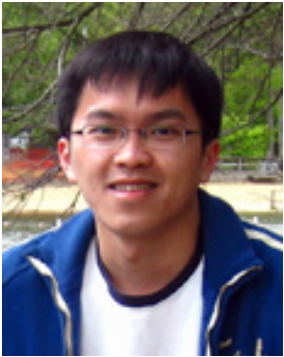


S. J. Pennycook and M. Varela, *J Electron Microsc*, 60, S213–S223 (2011).

# Outline

- **Point defects in graphene**
  - Configuration
  - Bonding, electronic structure
  - Dynamics
- **Optical properties**
  - “Plasmons” in graphene
  - Atomic resolution valence EELS: optical properties defect by defect!
- **Future directions**
  - meV resolution EELS
  - 3D atomic resolution

# Graphene team



Wu Zhou



Juan-Carlos Idrobo



Myron Kapetanakis



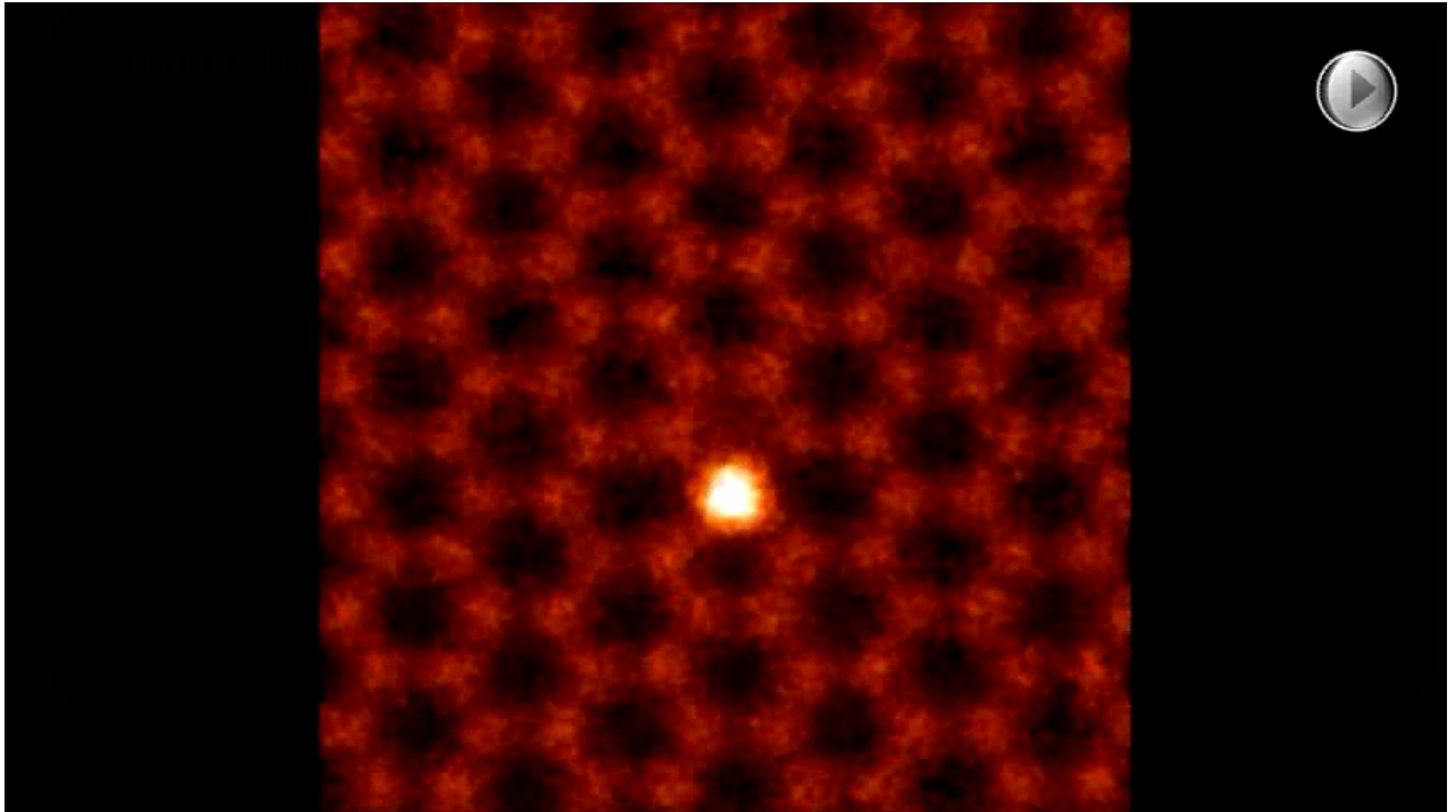
Jaekwang Lee

**Single atom microscopy:  
Bonding, Dynamics and  
Optical Properties**



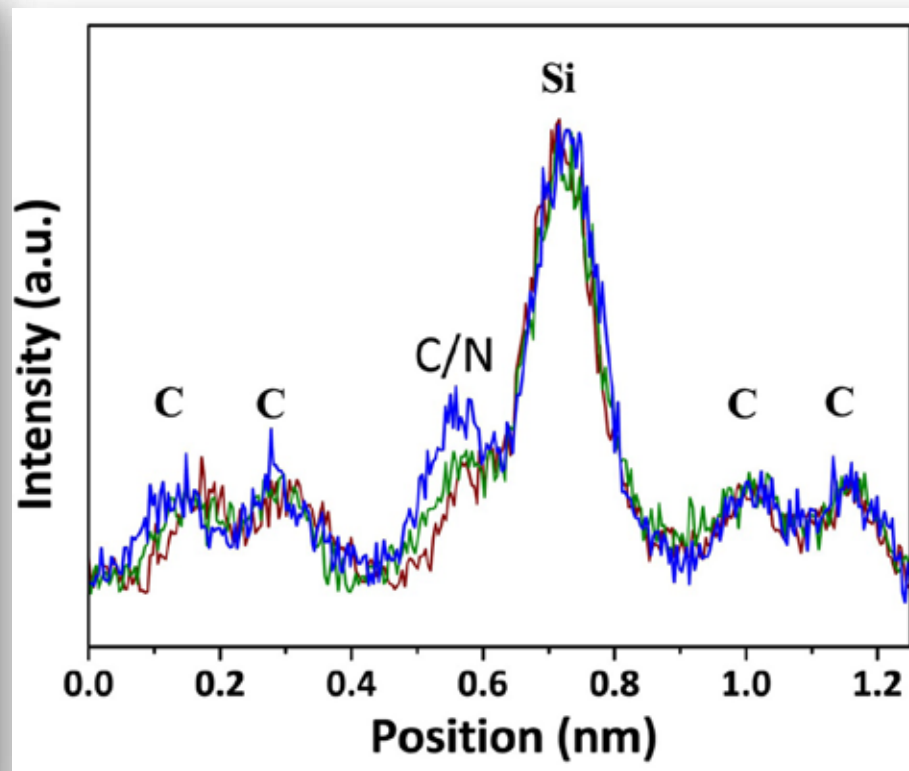
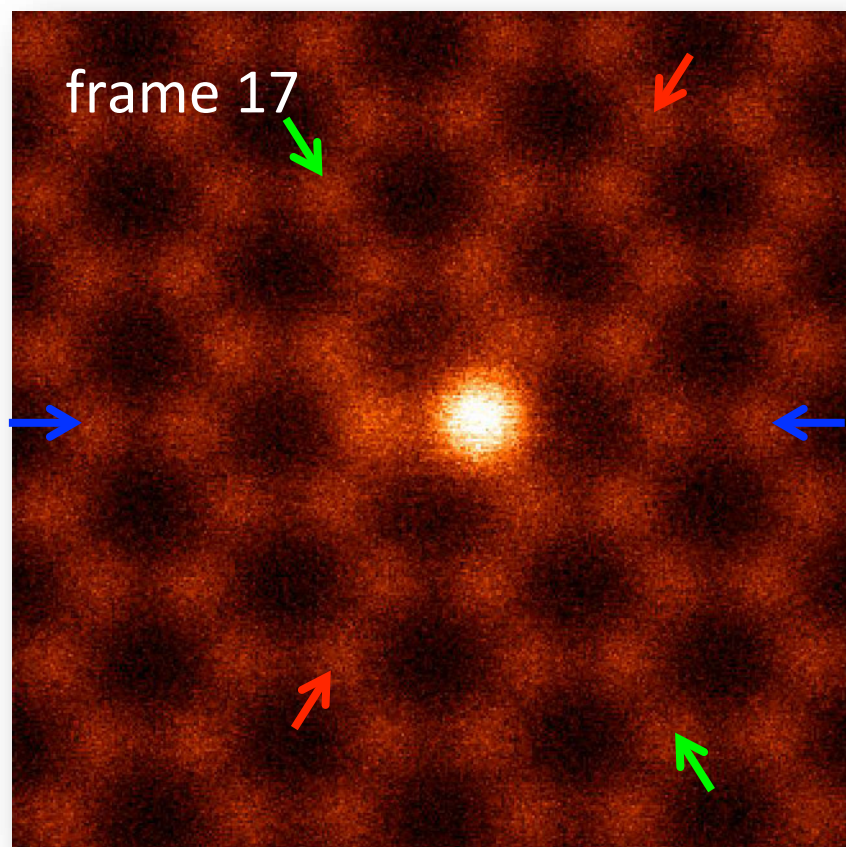
Sok Pantelides

# Tracking point defect dynamics

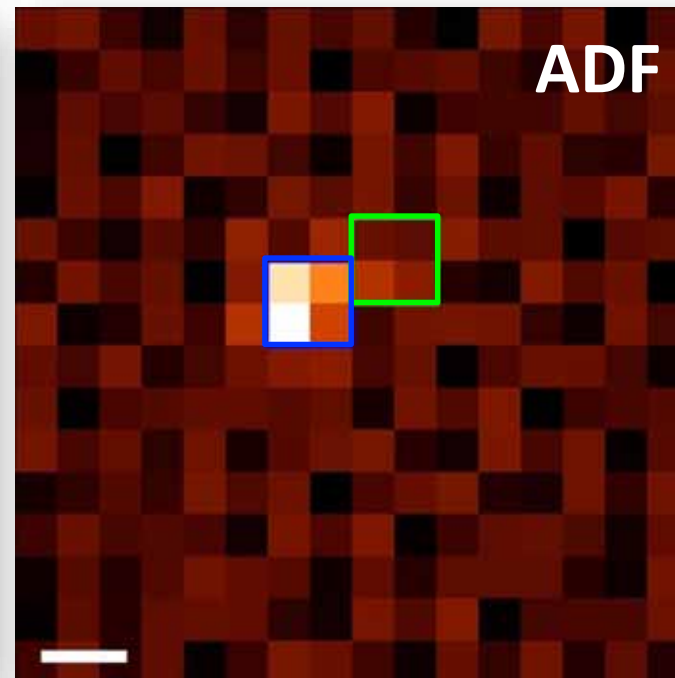
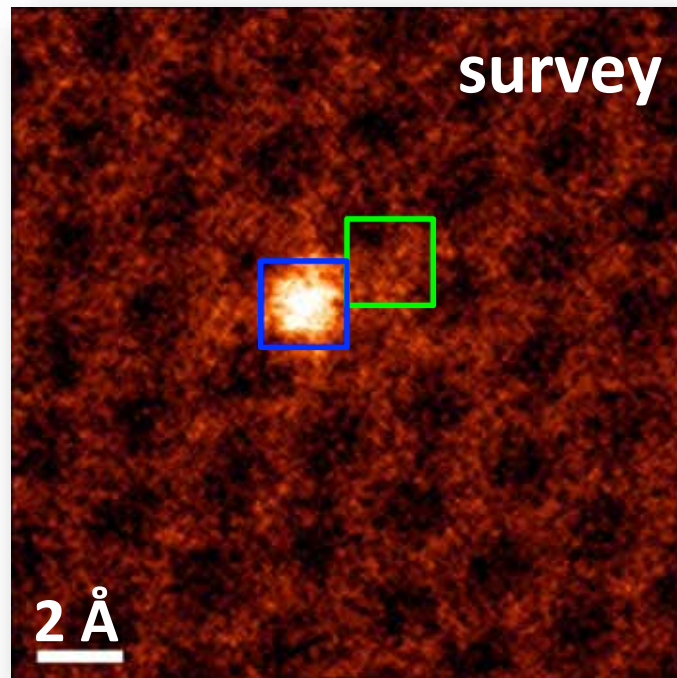


Zhou et al. *Microsc. Microanal.* (2012)

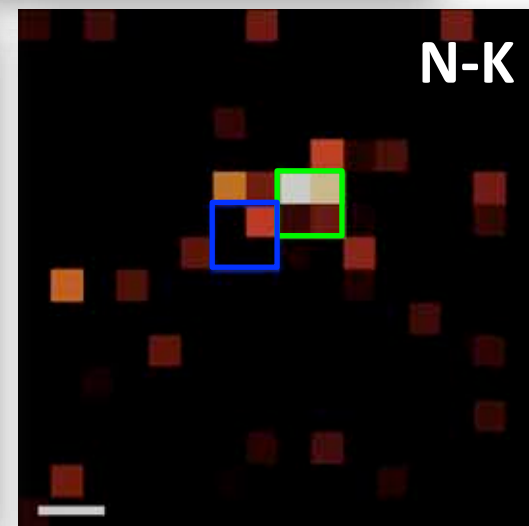
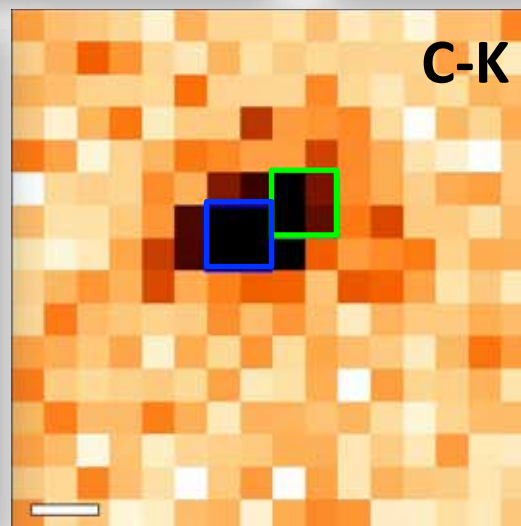
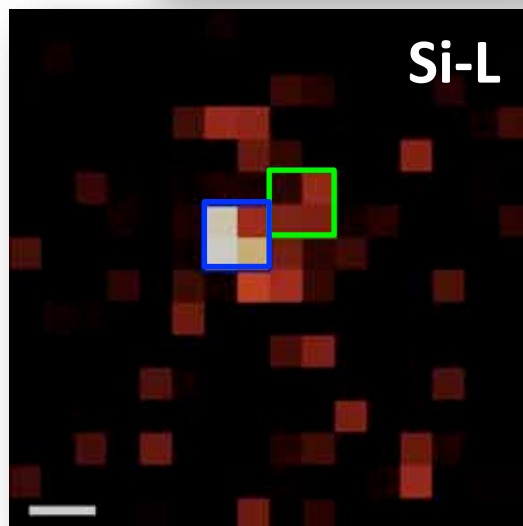
# Direct identification of species from image intensity



# Atom by atom spectroscopy

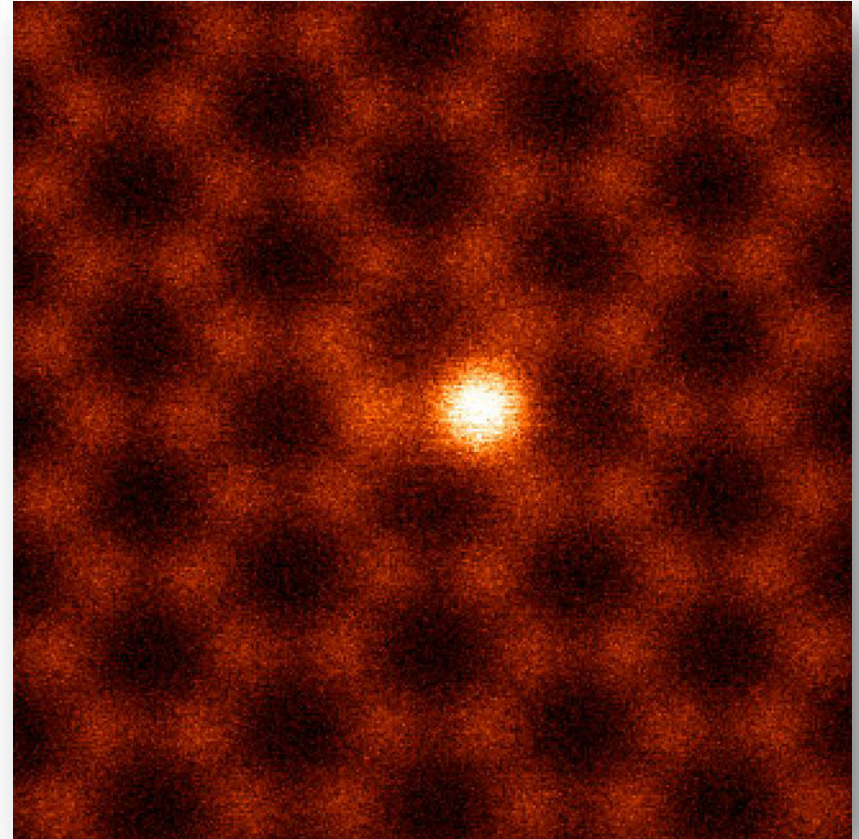
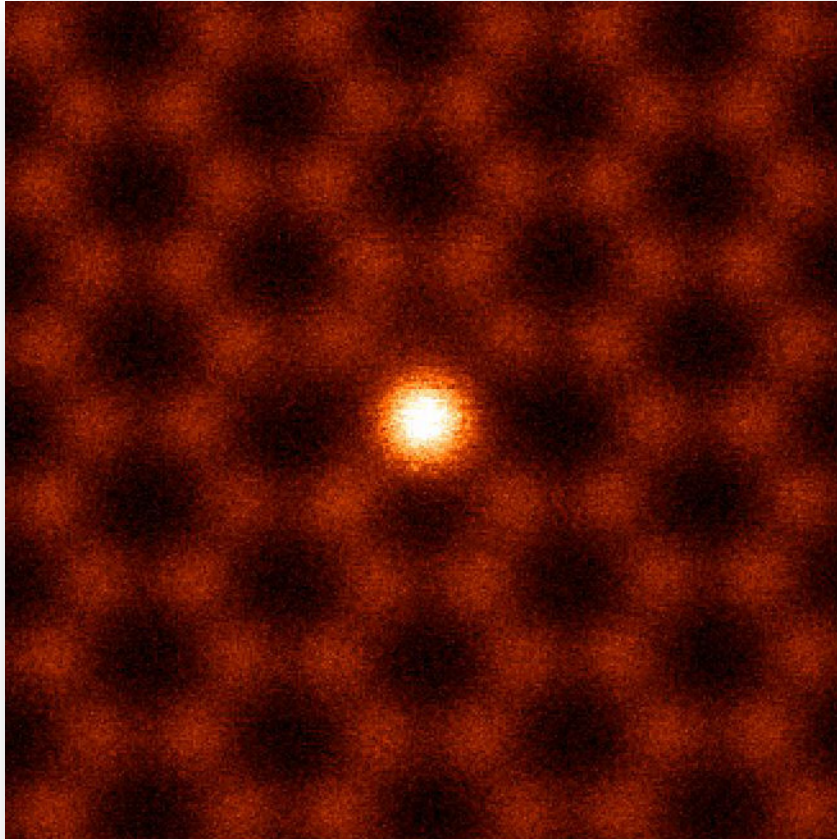


□ N  
□ Si



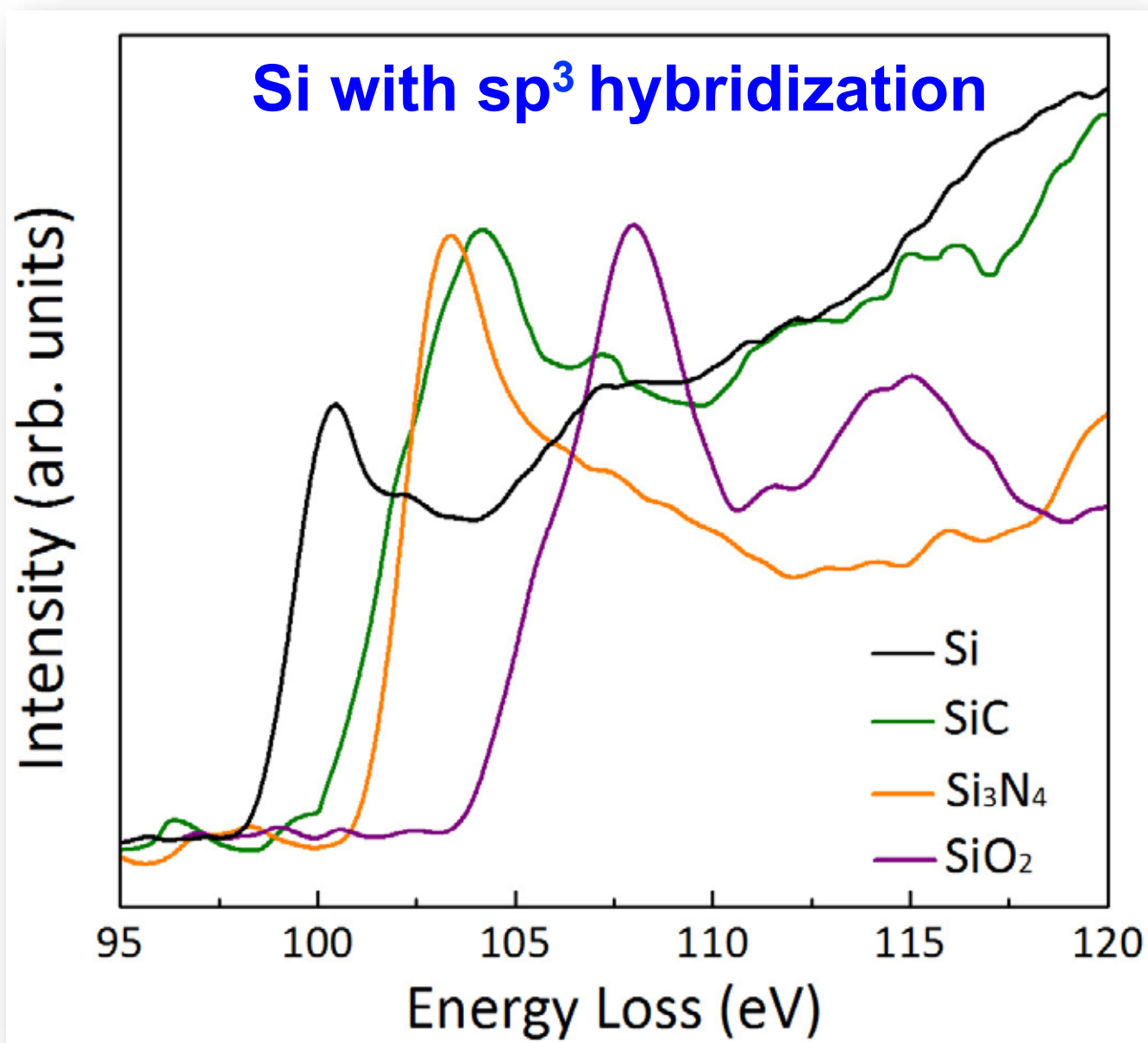


# Revealing the nature of chemical bonding

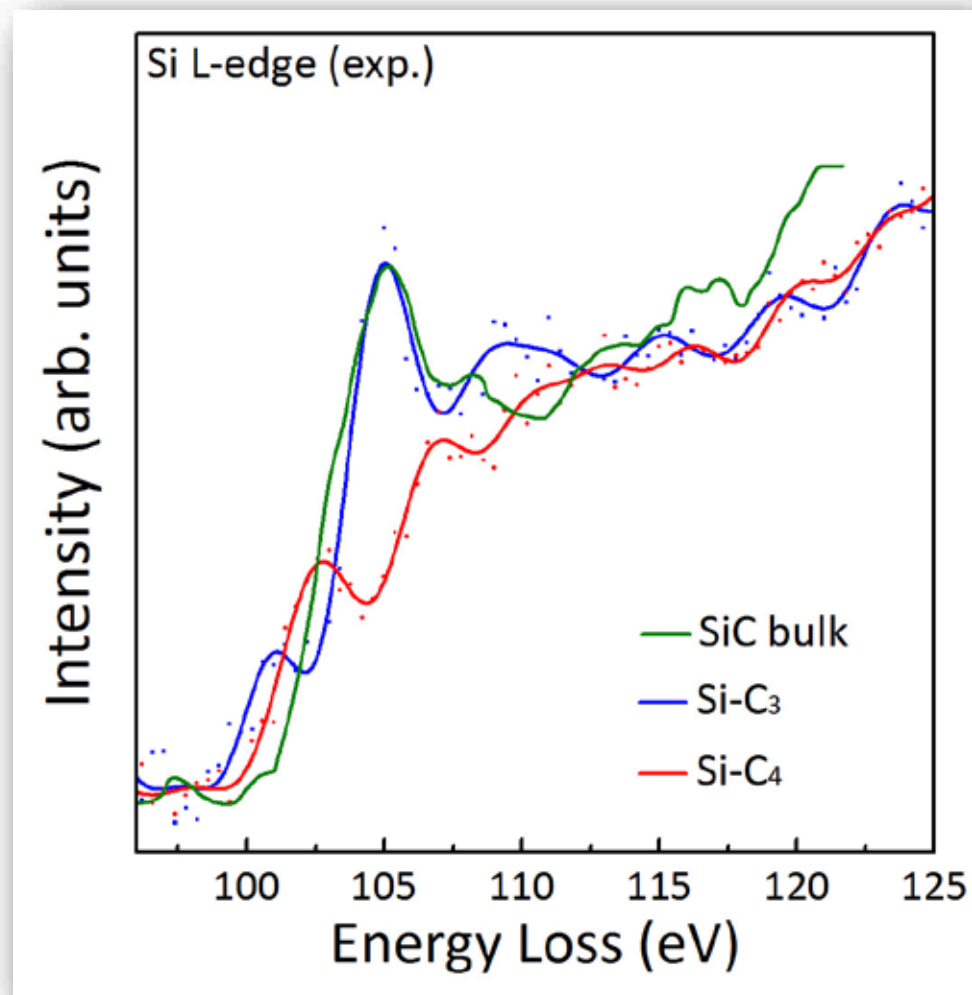
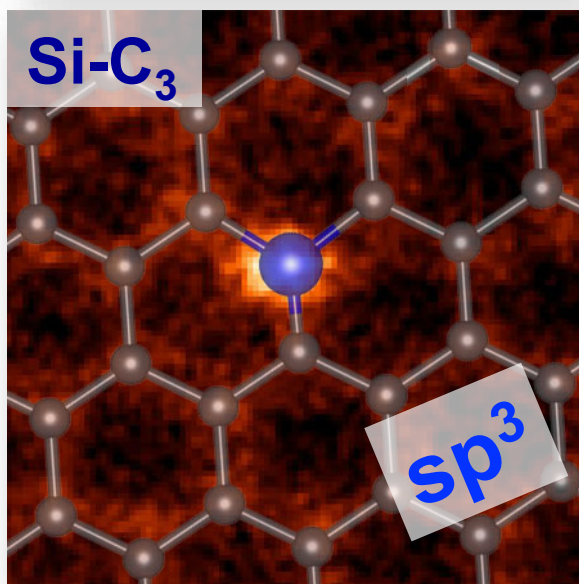
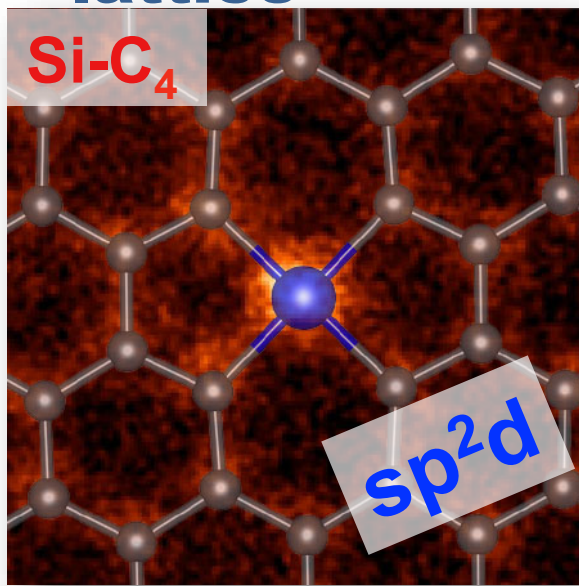


**How do the Si atoms bond in the graphene lattice?**

## Reference spectra for the Si L-edge



# Bonding of single Si atoms in the graphene lattice



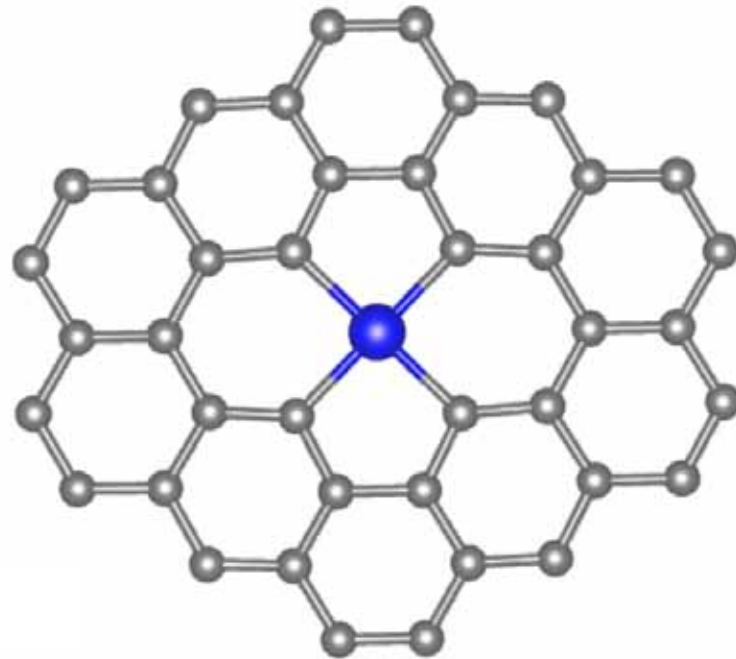
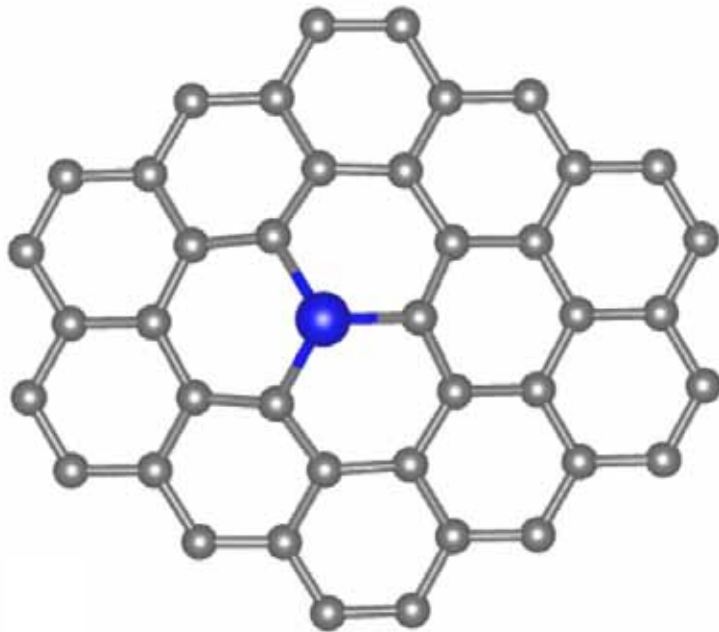
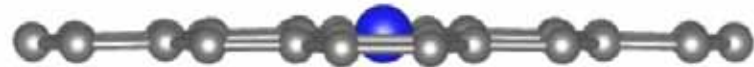
**Some 3d states in Si-C<sub>4</sub> structure are missing!**

# 3D structure from DFT calculations

$sp^3$

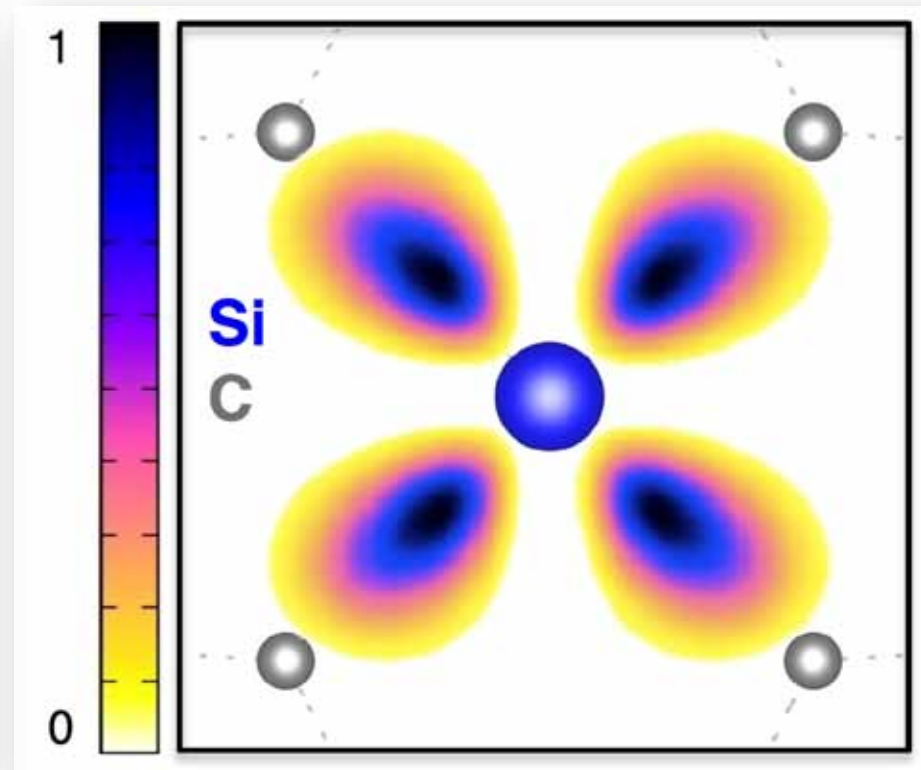
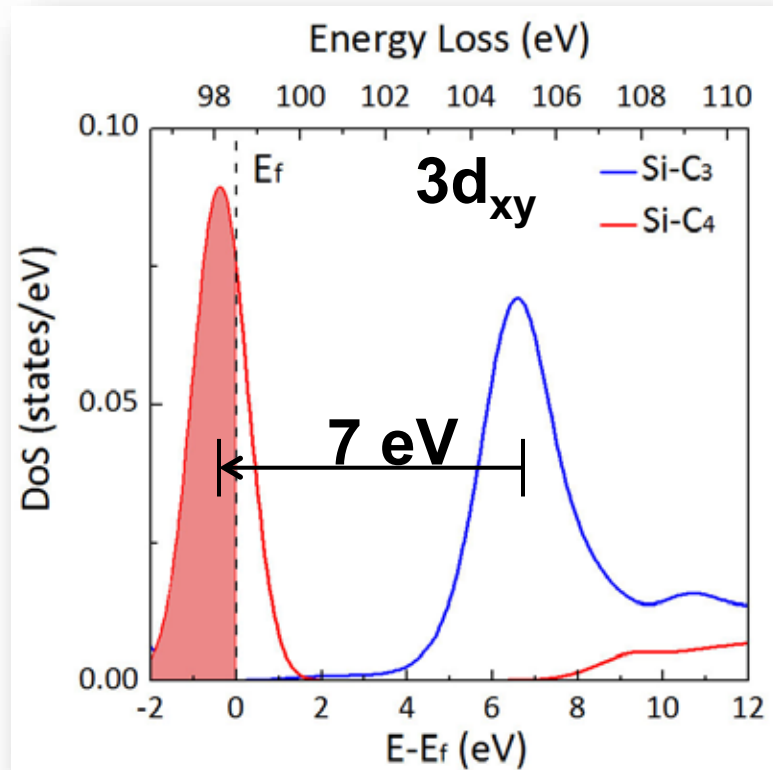


$sp^2d$



0.54 Å out-of-plane for Si-C<sub>3</sub>  $\implies$   $sp^3$ -like hybridization for Si-C<sub>3</sub>  
0.63 Å out-of-plane for SiC bulk

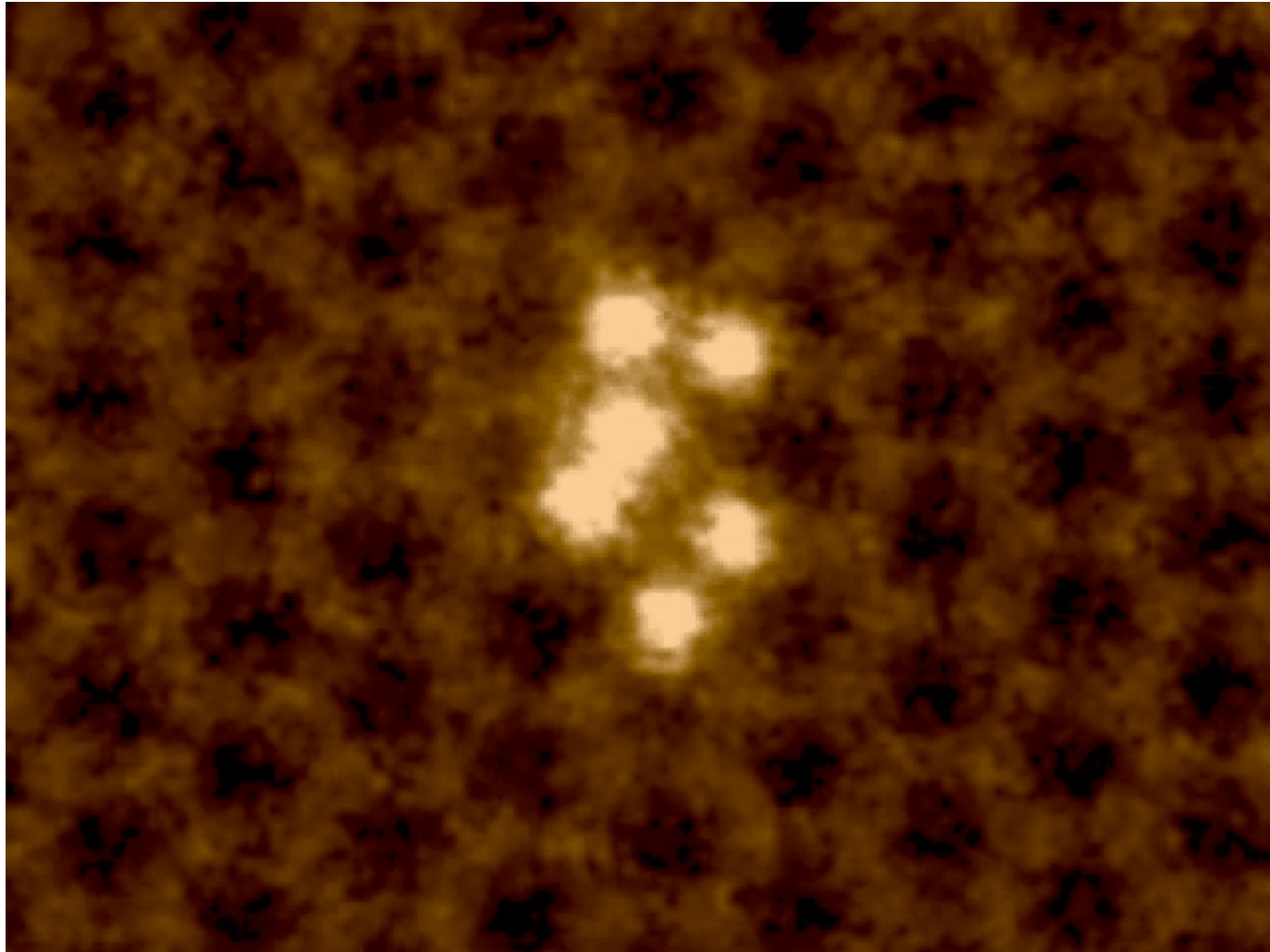
# Participation of $d_{xy}$ states in chemical bonding of Si-C<sub>4</sub>



**Si  $3d_{xy}$  states participate in the chemical bonding of Si-C<sub>4</sub>  
 $sp^2d$  hybridization for planar 4-fold coordination**

W. Zhou, M. D. Kapetanakis, M. P. Prange, S. T. Pantelides, S. J. Pennycook, and J.-C. Idrobo *Phys Rev Lett*, 109, 206803 (2012).

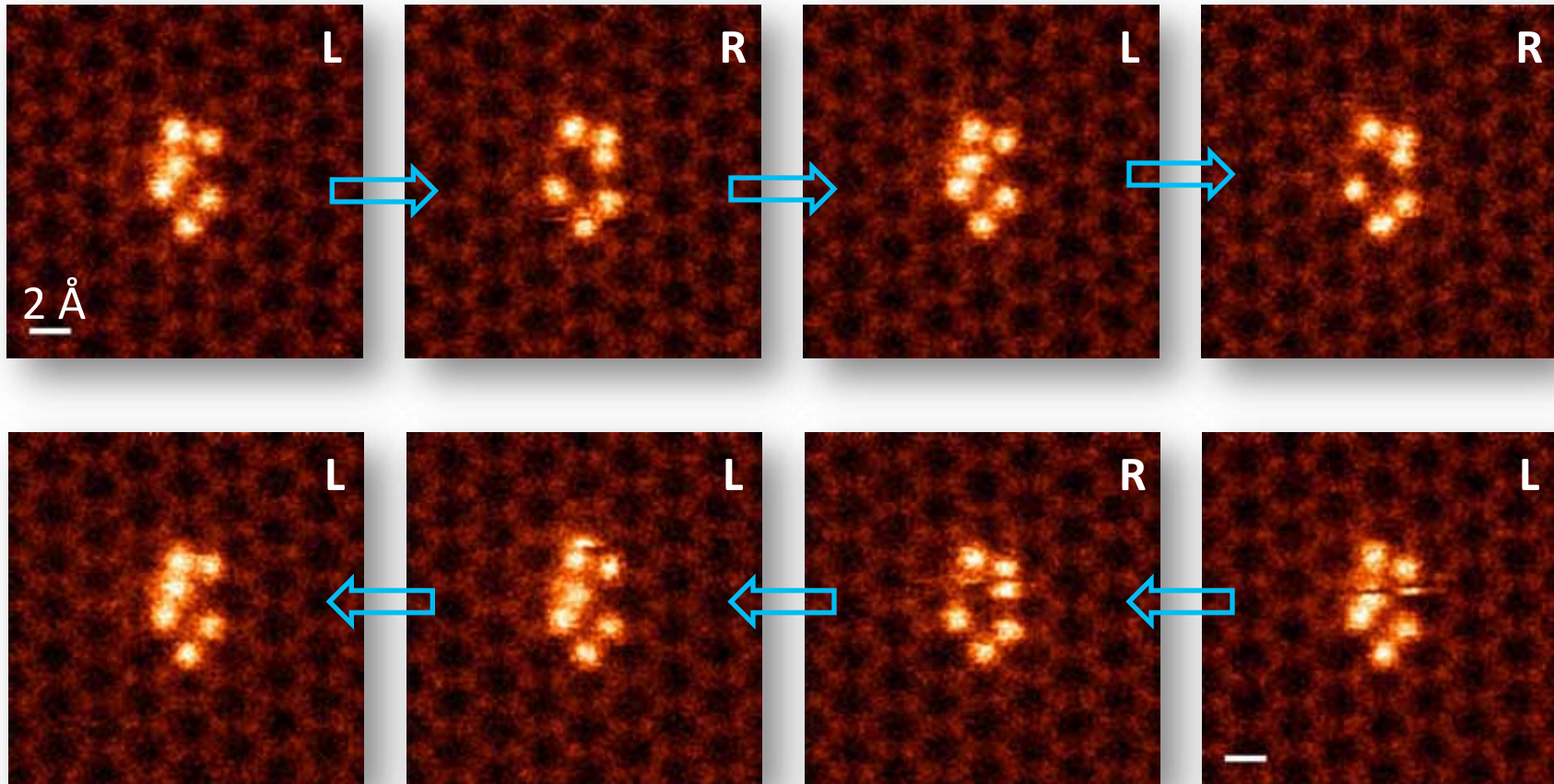
# **Si<sub>6</sub> magic cluster**



**UltraSTEM  
at 60 kV**

**W. Zhou, J-C Idrobo: Si on graphene**

# Reversible dynamics at the atomic scale

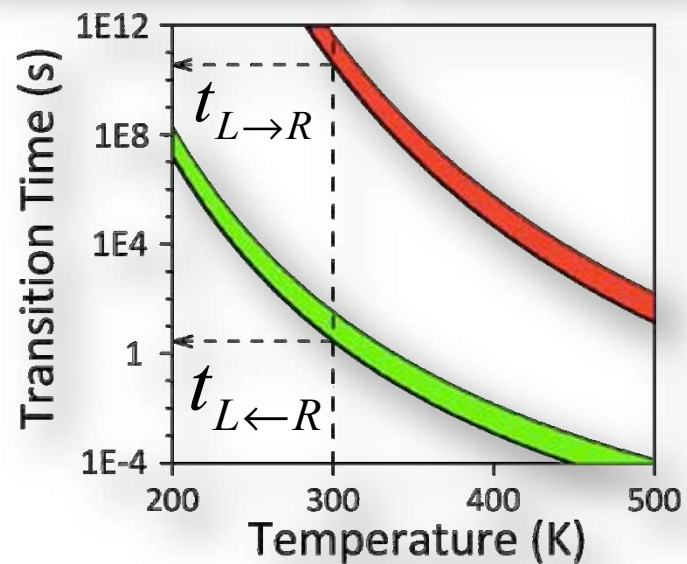
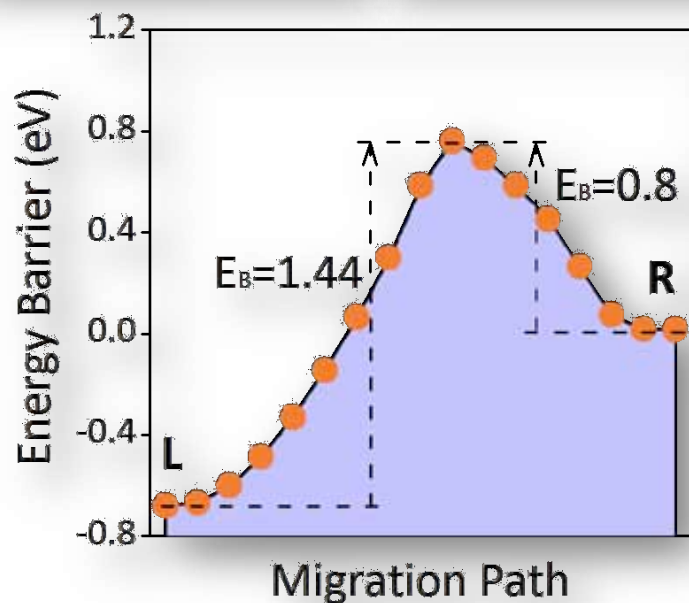
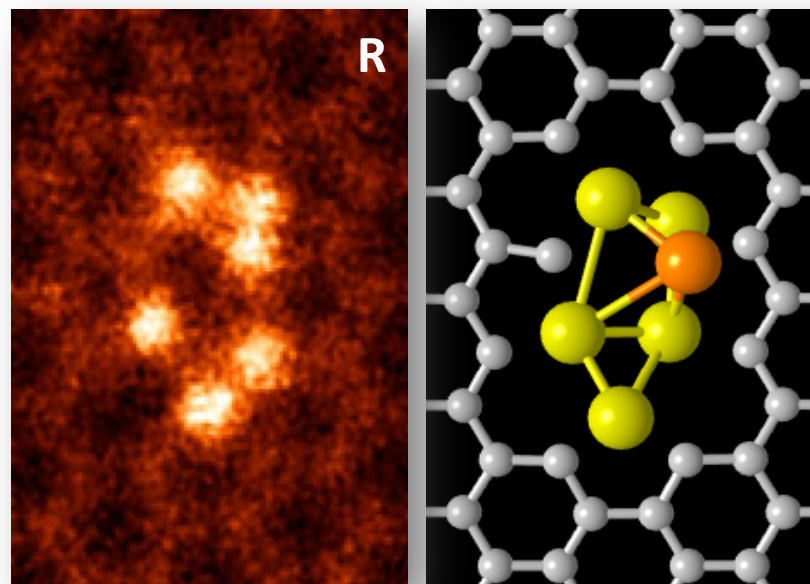
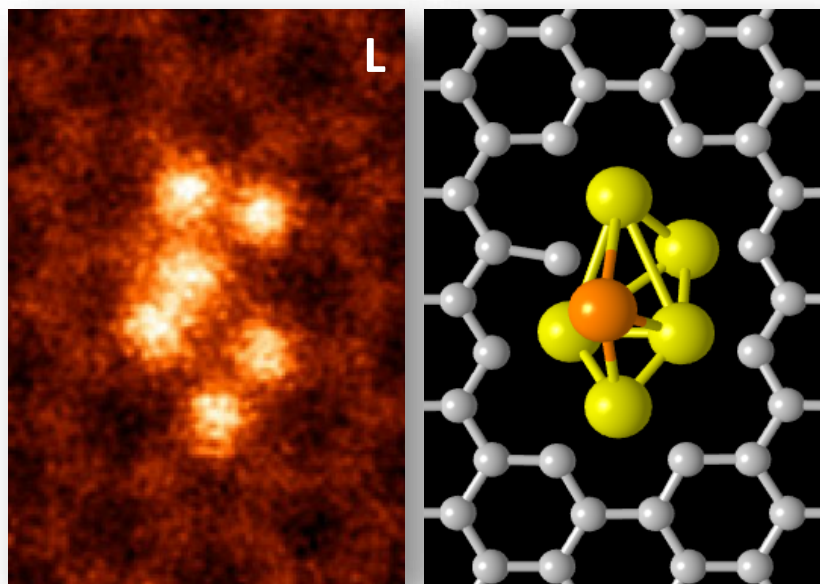


6 s/frame



Turn to theory: Jaekwang Lee

# Atomic scale molecular dynamics



J. Lee, et al. *Nature Commun.* (2012)



## Optical properties from EELS

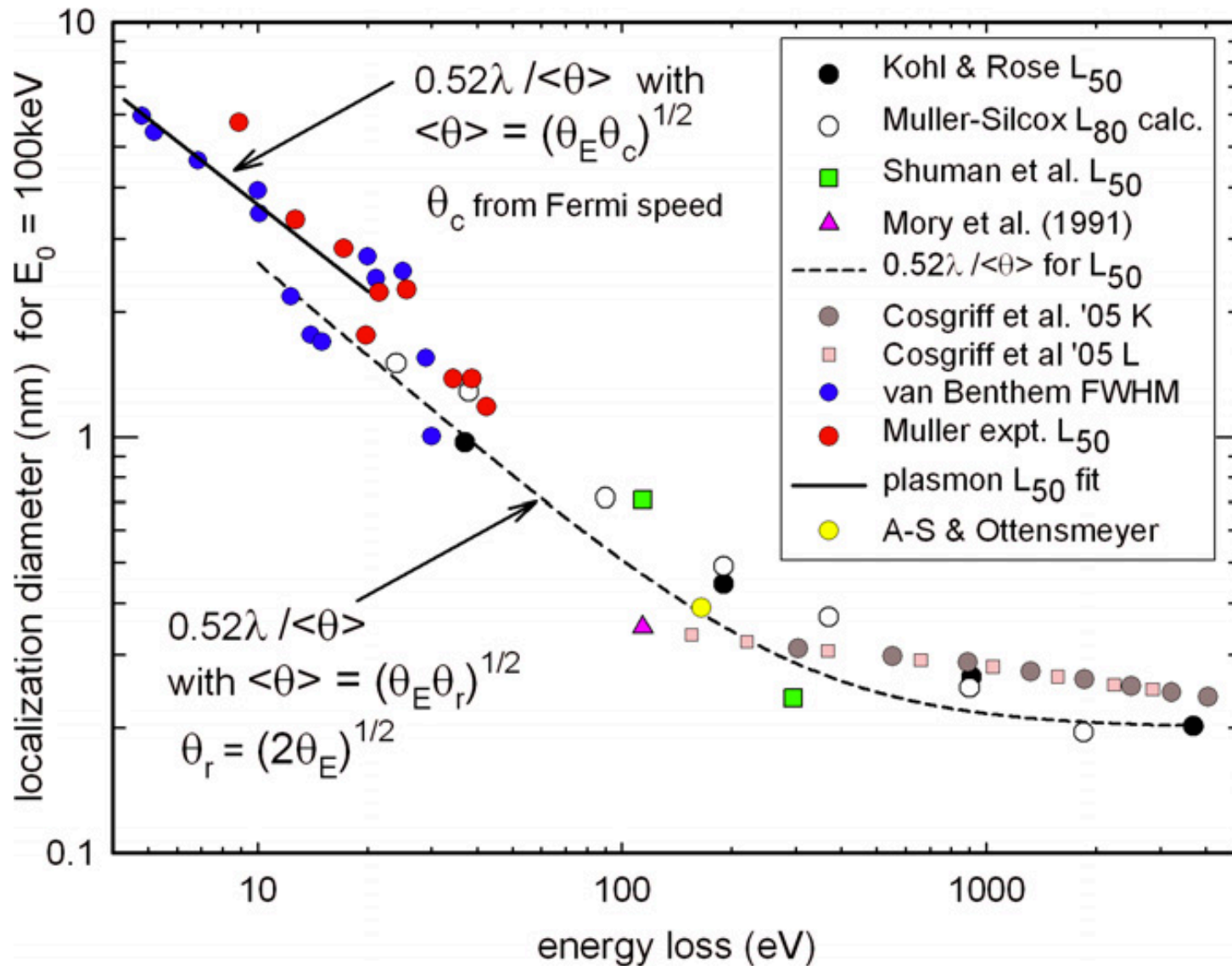
Energy loss:  $\chi(E) = \text{Im} \left( \frac{1}{\varepsilon(E)} \right)$

$$\chi(E) \rightarrow \varepsilon(E) \rightarrow \varepsilon_1(E), \varepsilon_2(E)$$

↑  
Optical absorption

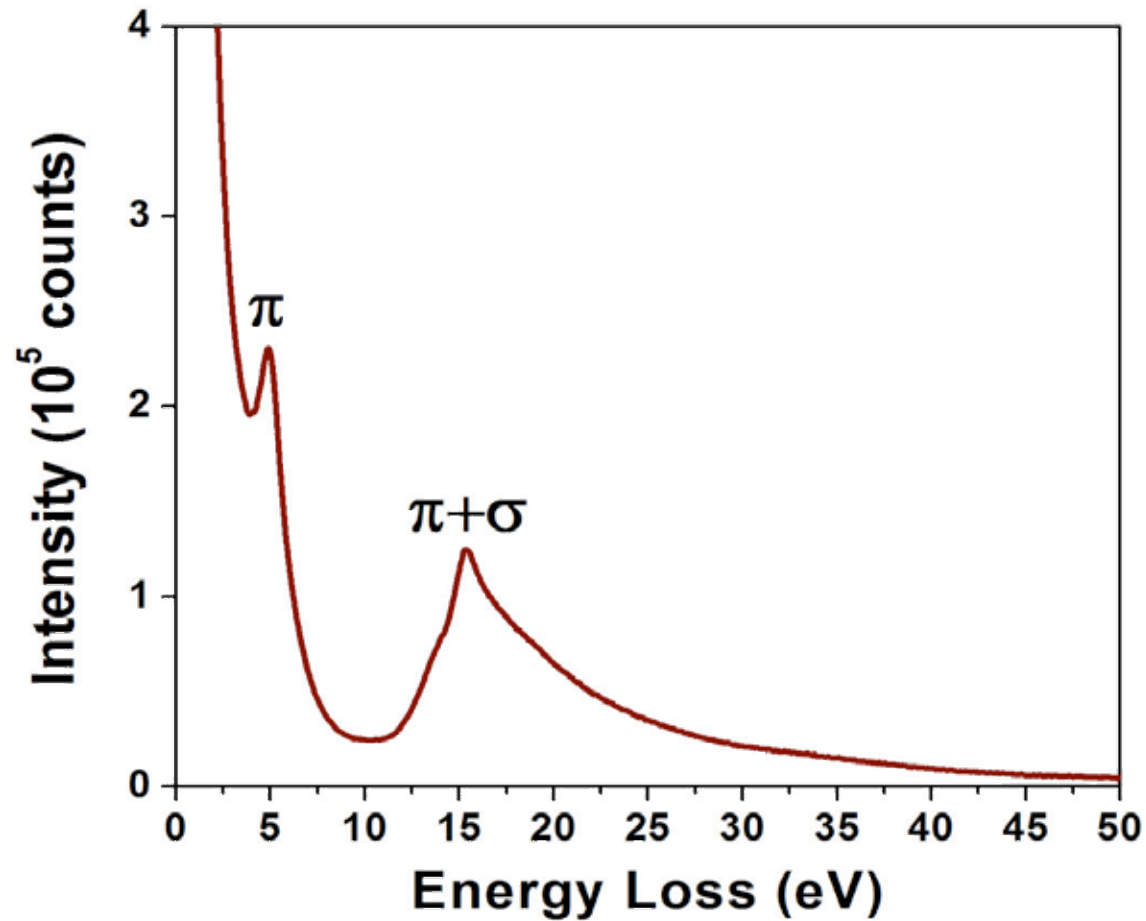
**STEM-EELS: Optical absorption with atomic resolution**

# Delocalization is a problem?



R. F. Egerton, Rep. Prog. Phys. 72, 6502 (2009).

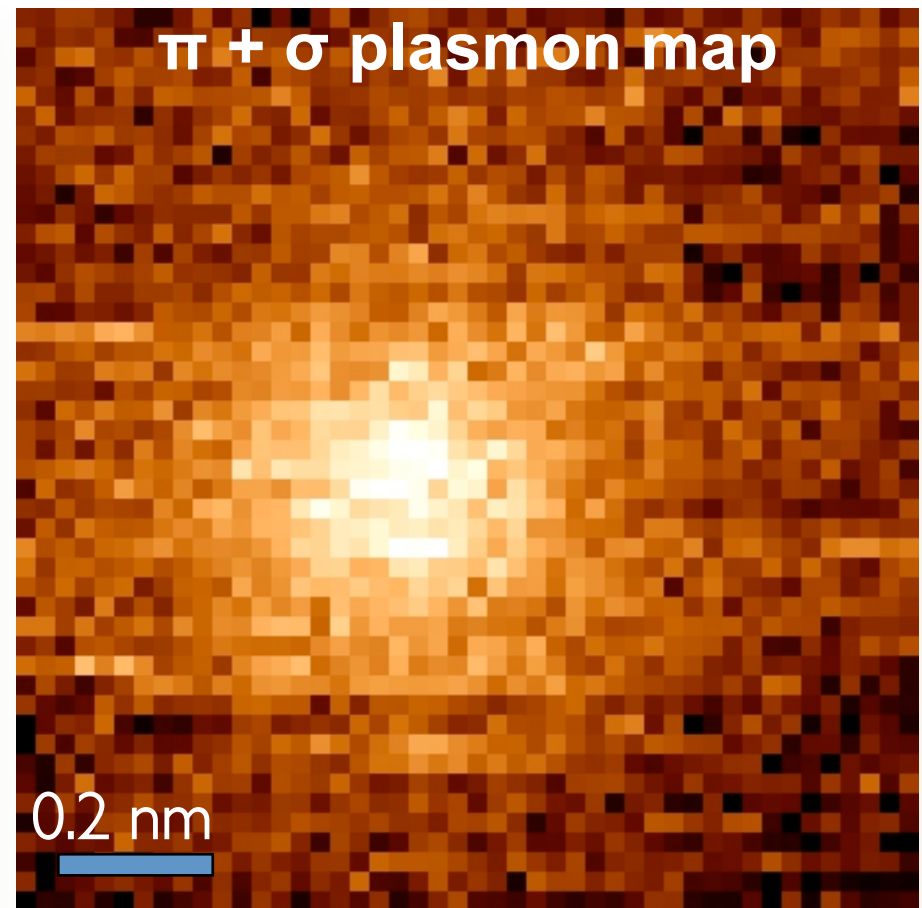
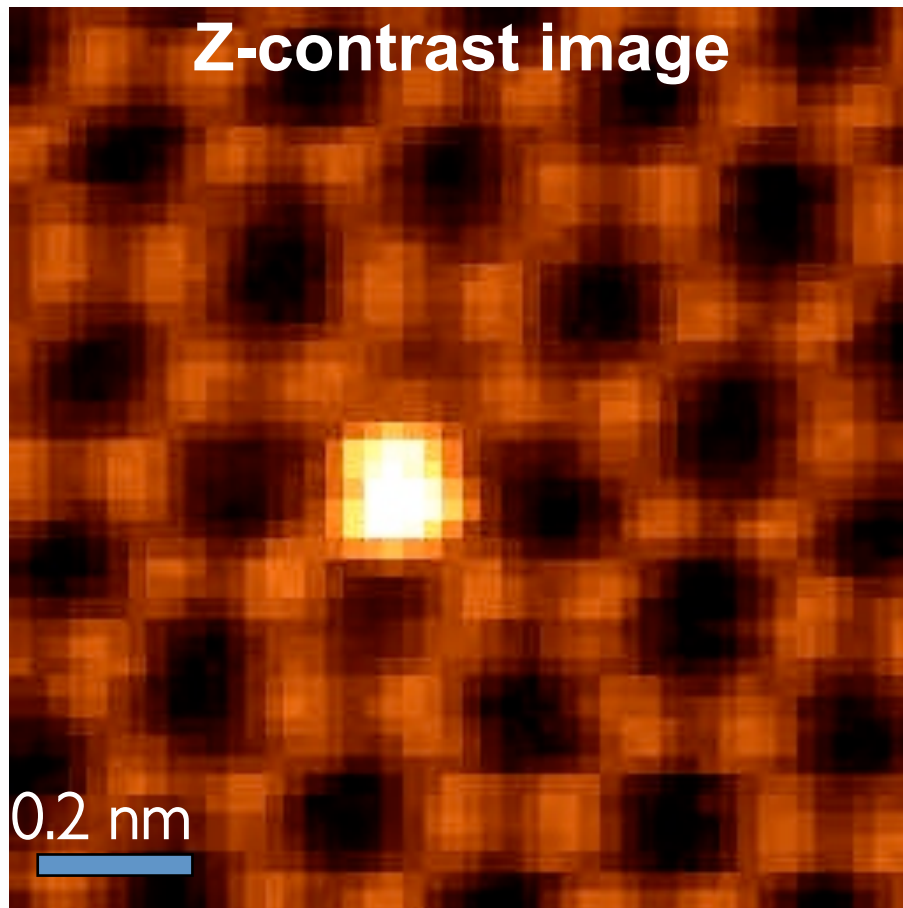
# Typical “plasmon” peaks of monolayer graphene



How do point defects affect the plasmon response?

# Atomically localized plasmon enhancement

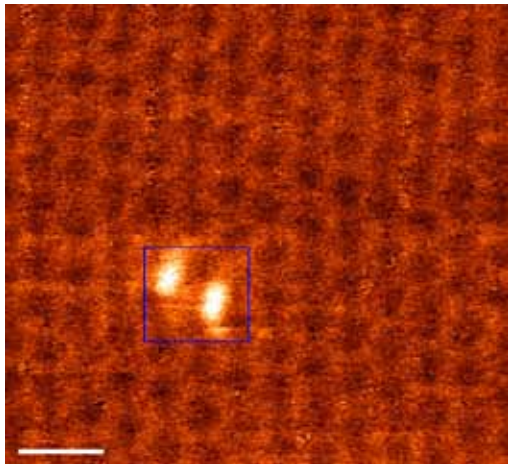
**FWHM for enhanced plasmon:  $0.43 \pm 0.05$  nm ( $< \lambda/200$ )**



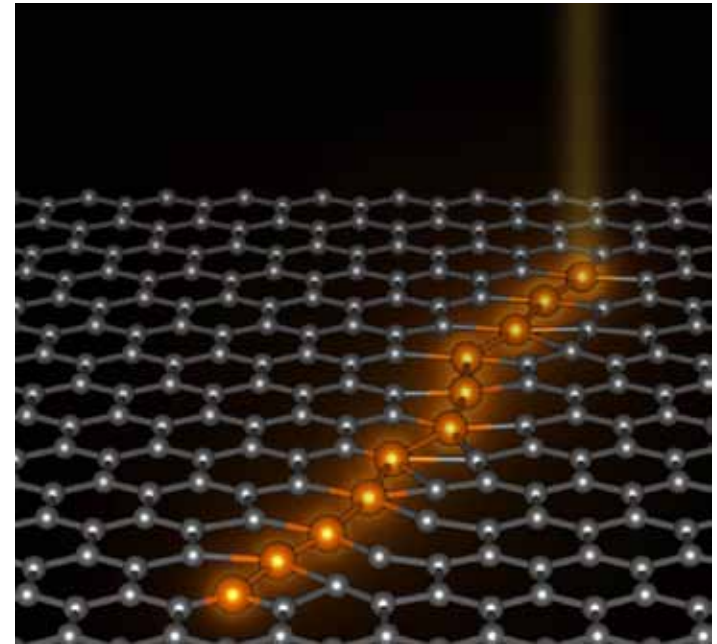
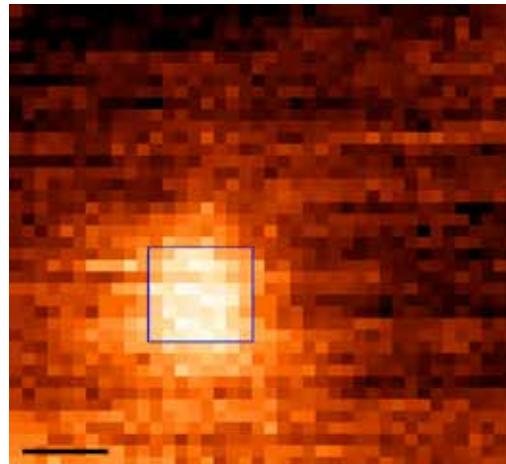
W. Zhou, J. Lee, J. Nanda, S.T. Pantelides, S.J. Pennycook and J.C. Idrobo  
Nature Nanotechnology 7, 161 (2012).

# Atomically localized plasmon resonance

Z-contrast



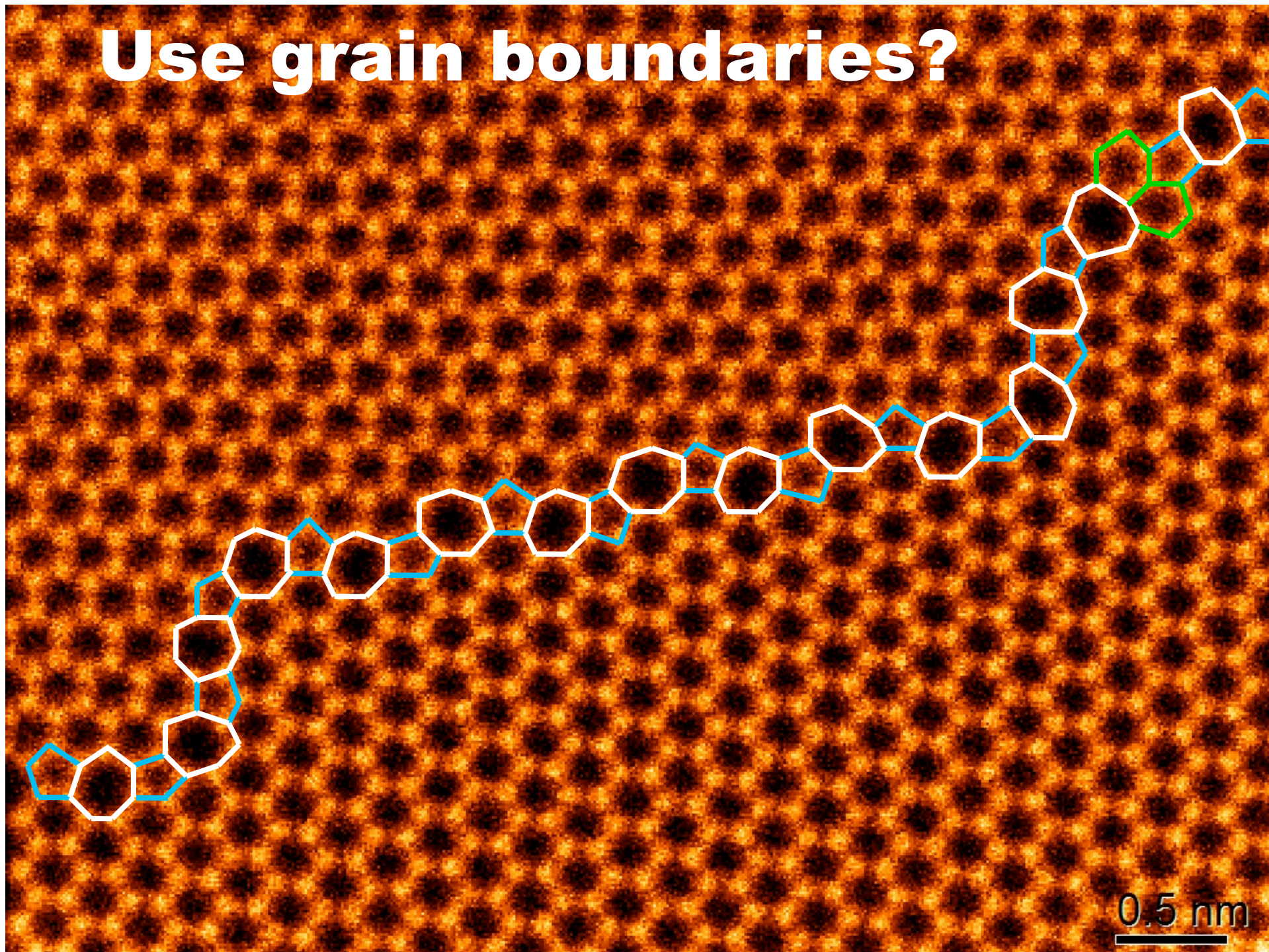
plasmon (11-18 eV)



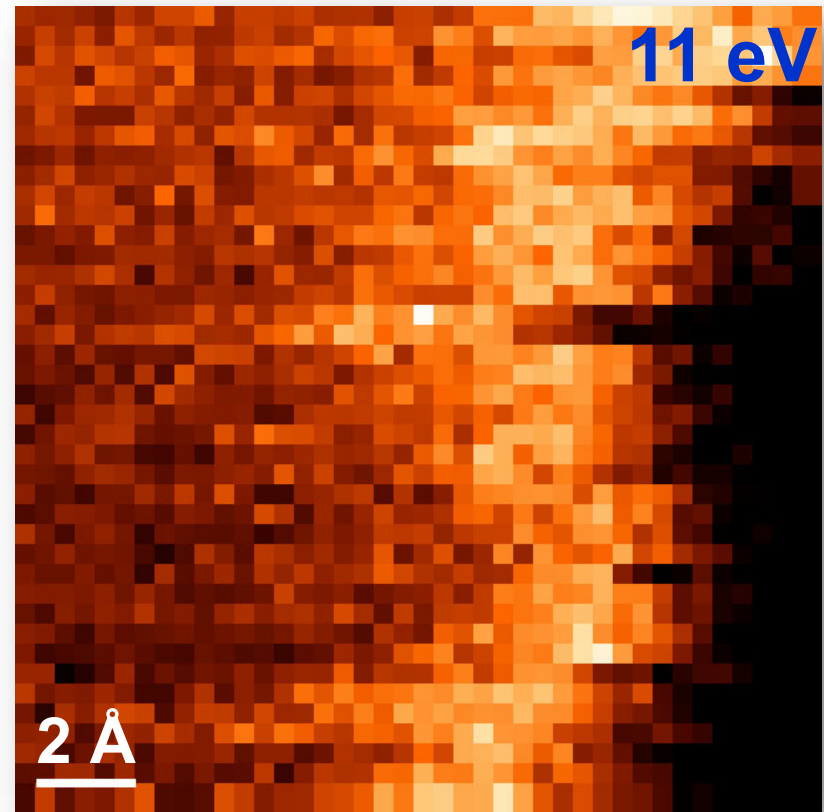
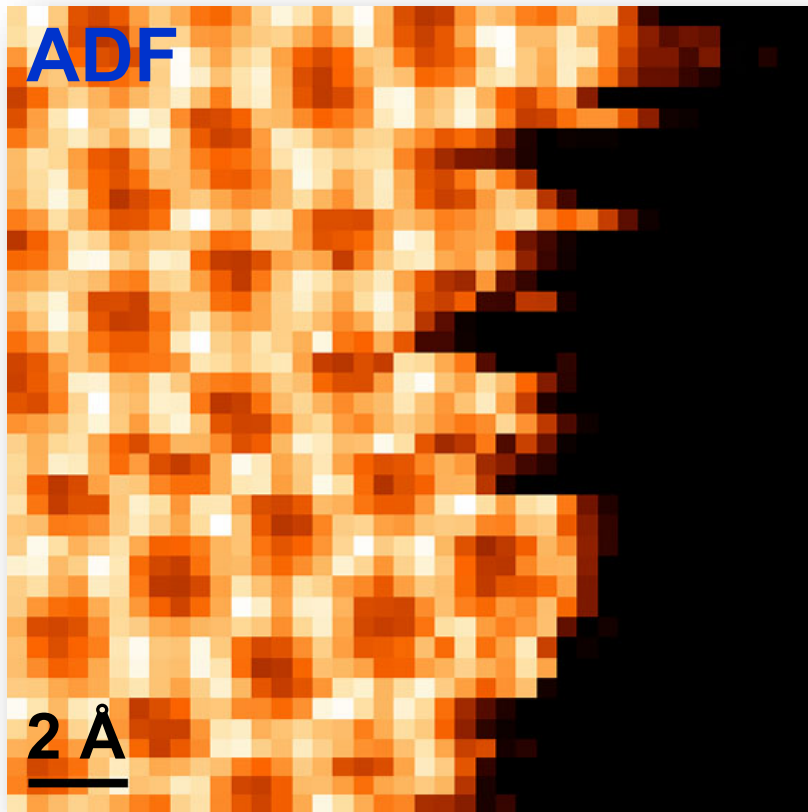
High frequency ( $10^{15}$  Hz) signals could be transmitted along atomically confined paths by assembling single atoms on graphene

\*W. Zhou et al., Nat. Nanotech. 7, 161 (2012)

# Use grain boundaries?



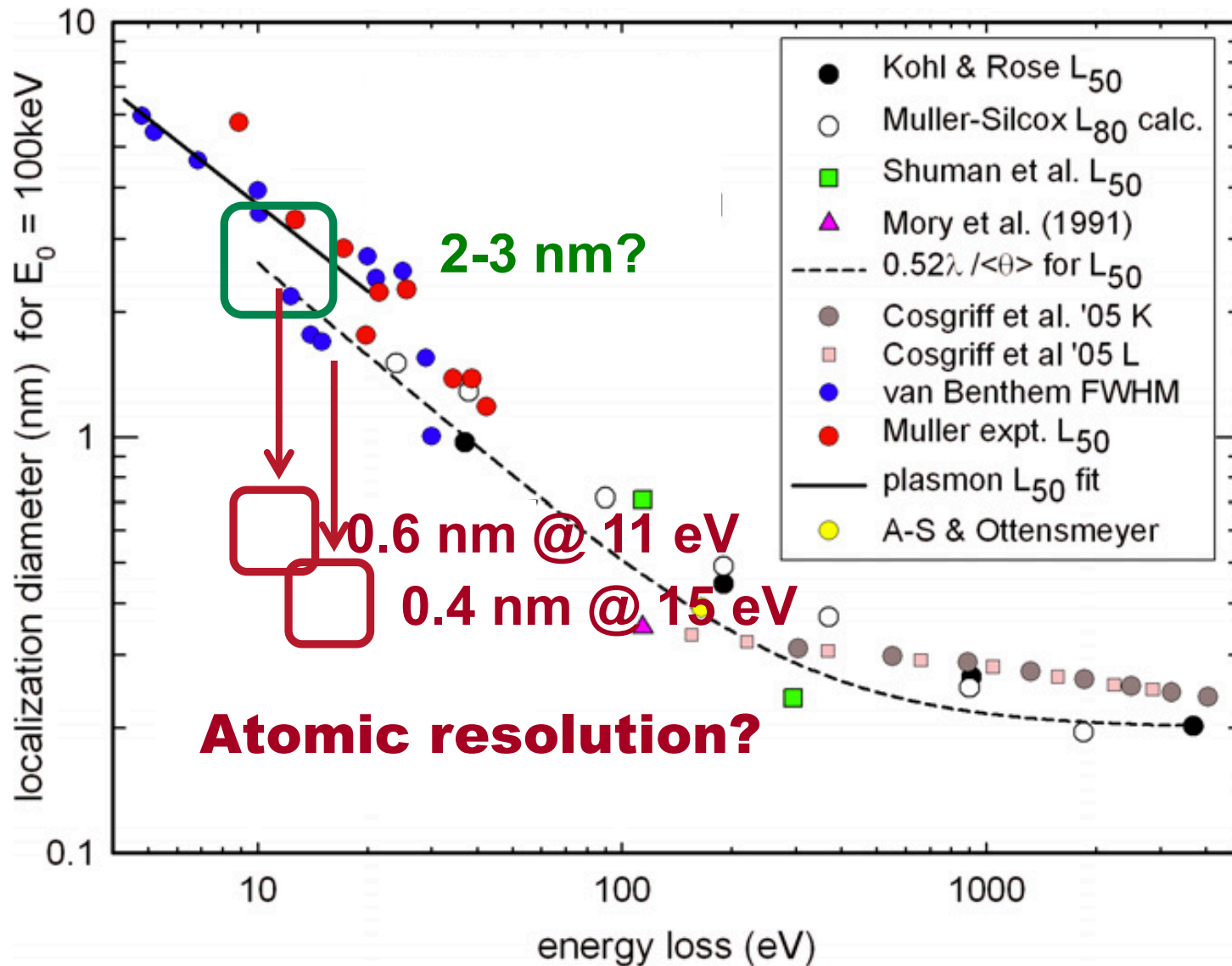
# 1-D edge plasmon on graphene with 6Å localization



**Localization depends both on the energy loss and the specific electron excitation mode contributing to the energy loss**

W. Zhou et al. Ultramicroscopy (2012)

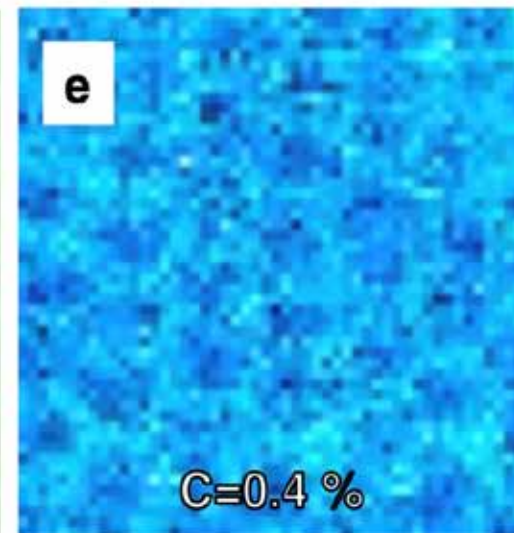
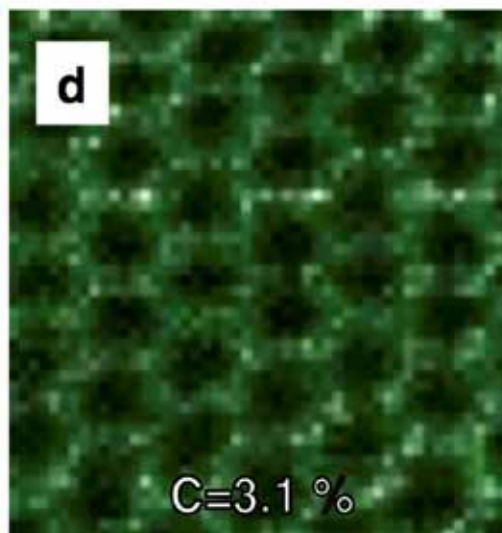
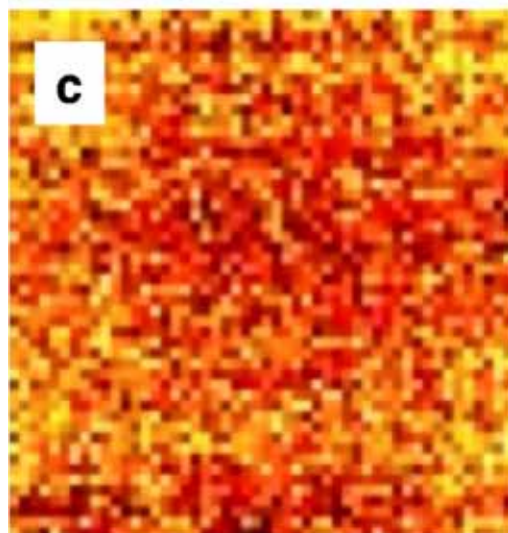
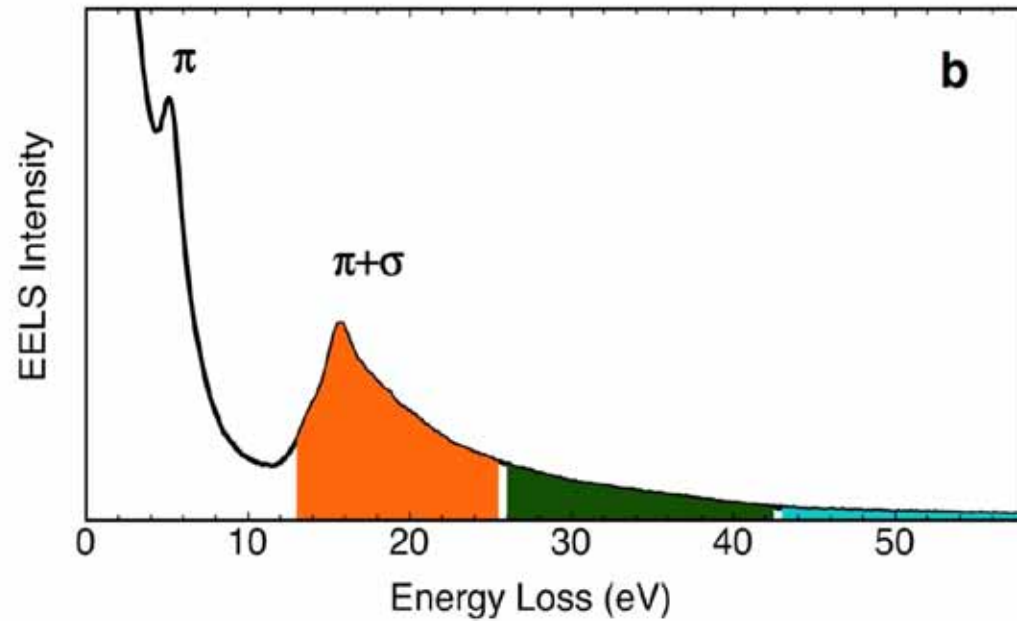
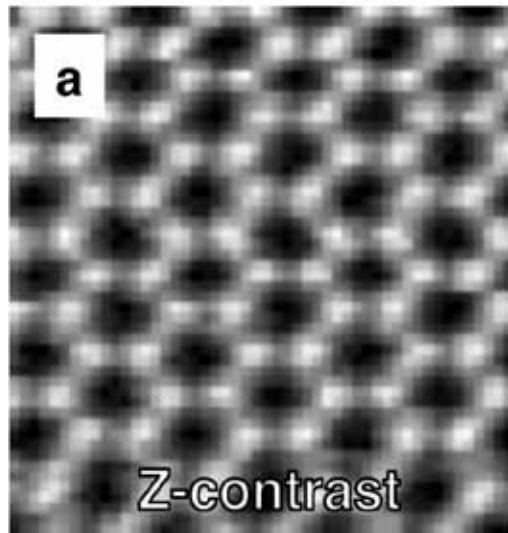
# Surprisingly localized



R. F. Egerton, Rep. Prog. Phys. 72, 6502 (2009).

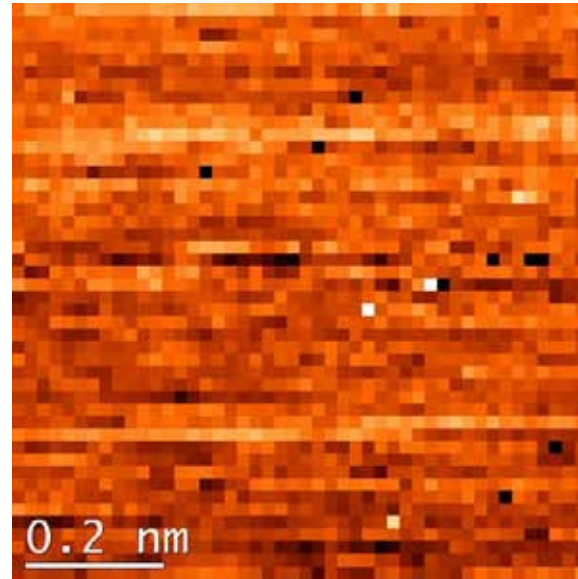
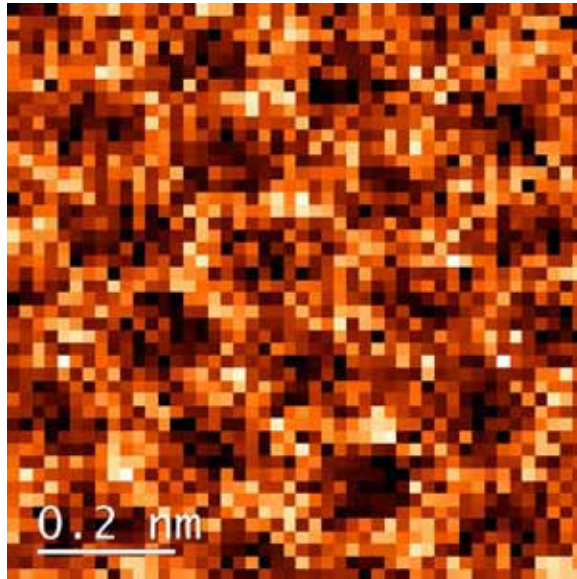


# STEM-VEELS SPECTRA AND MAPS -- GRAPHENE



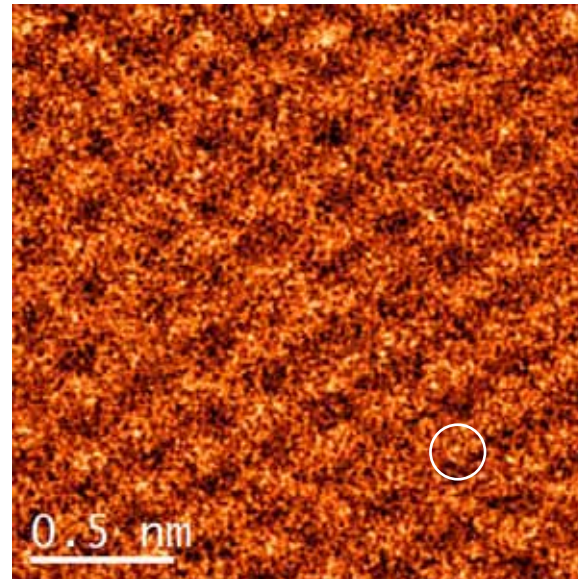
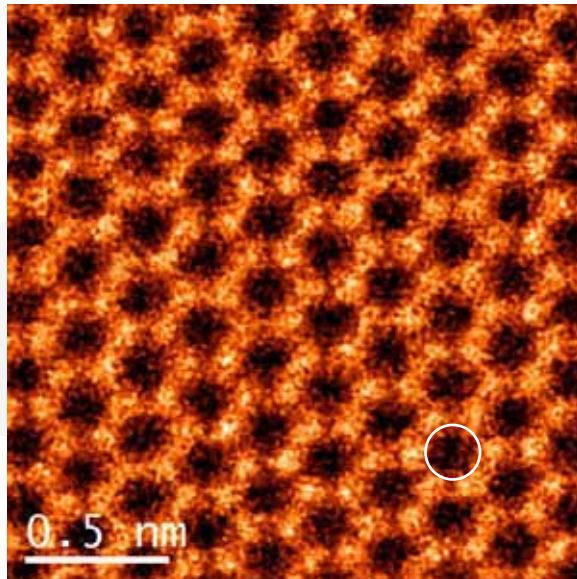
# NOT preservation of elastic contrast

HAADF



Zero loss

HAADF  
at optimum  
focus



Bright field  
shows  
reverse  
contrast

# Definition of Plasmons

David Pines & David Bohm, *Physical Review* **85**, 338 (1952)

$$\text{Im} \left( -\frac{1}{\epsilon(\omega)} \right) \epsilon(\omega) \rightarrow 0$$

$$\frac{\epsilon_2(\omega)}{\epsilon_1(\omega)^2 + \epsilon_2(\omega)^2}$$

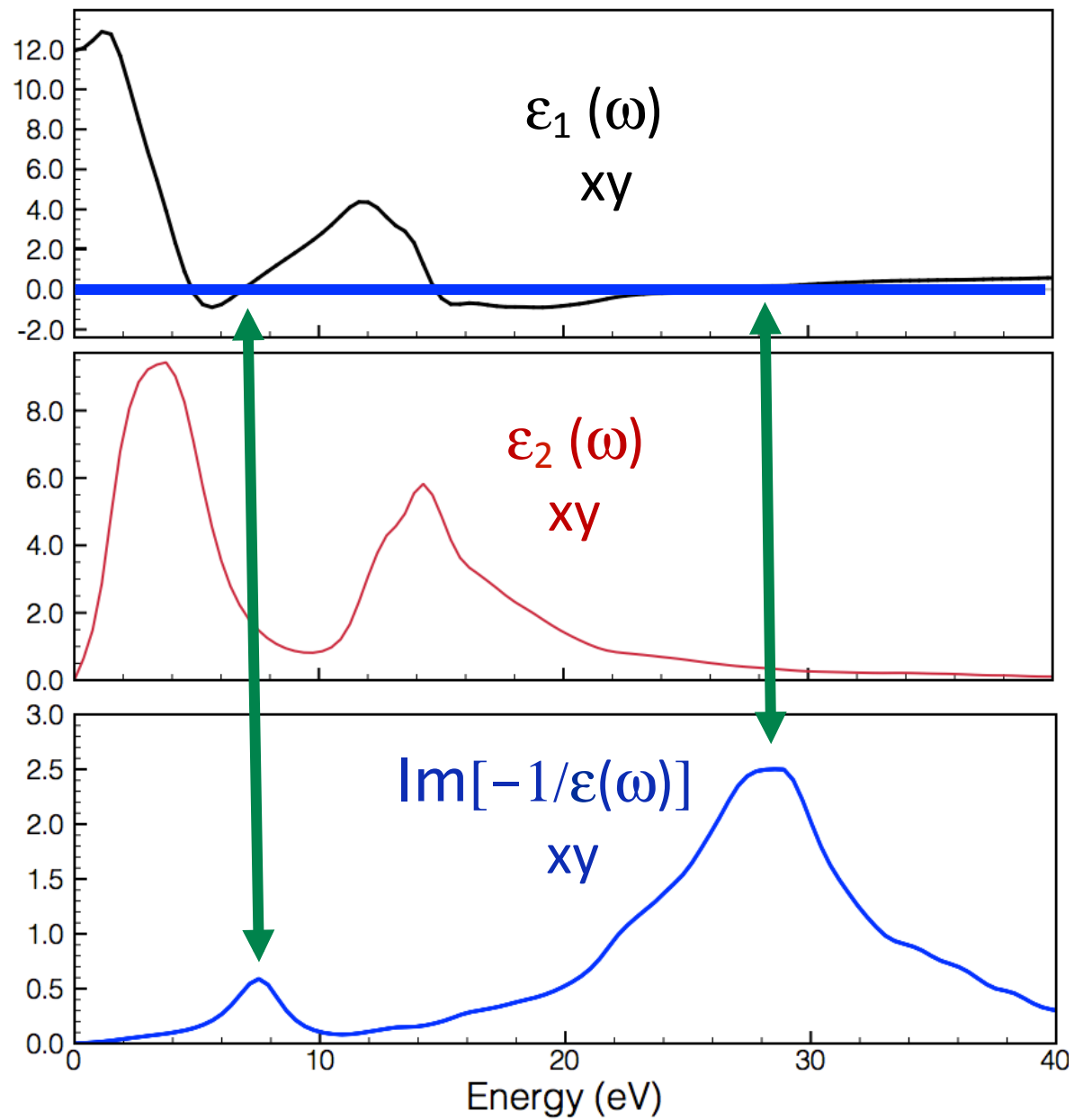


$$\epsilon_1(\omega) = 0$$

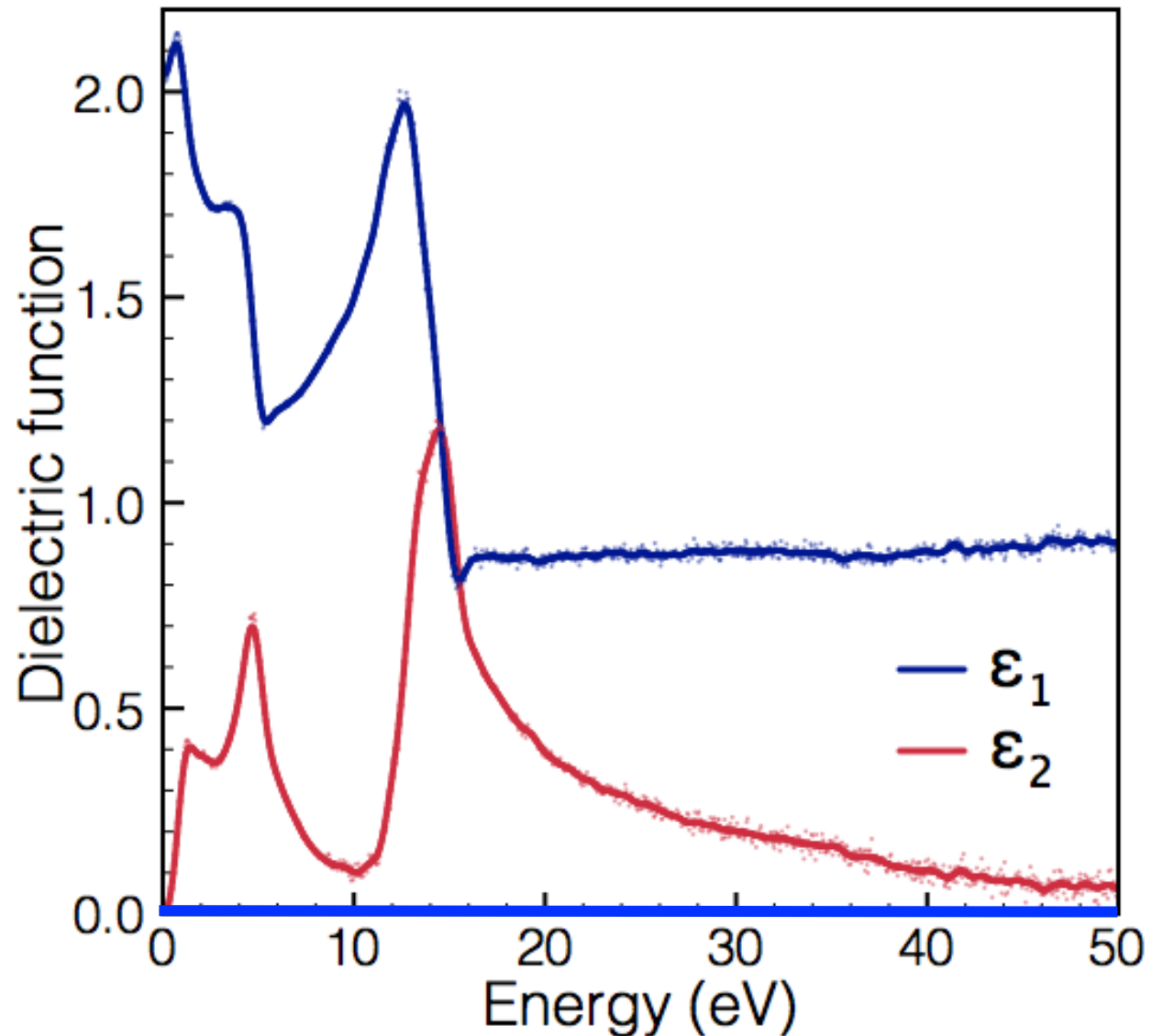


$$\epsilon_2(\omega) \approx 0$$

# Plasmons in graphite



**“Plasmons”  
in  
graphene  
are  
*interband*  
transitions  
some of  
which  
are  
localized**

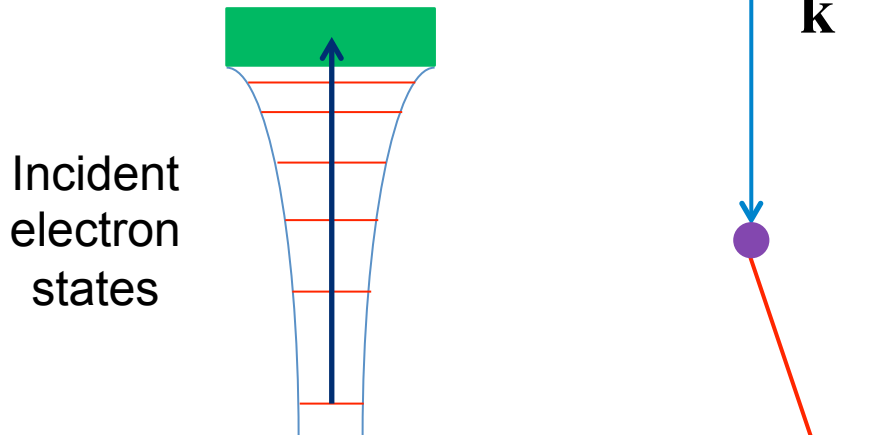


F. J. Nelson, J.-C. Idrobo, J. D. Fite, Z. L. Mišković, S. J. Pennycook, S. T. Pantelides, J. U. Lee, and A. C. Diebold, "Electronic Excitations in Graphene in the 1–50 eV Range: The  $\pi$  and  $\pi + \sigma$  Peaks Are Not Plasmons," *Nano Lett*, 14, 3827–3831 (2014).

# EELS simulation at atomic resolution

## Dynamical Electron Scattering

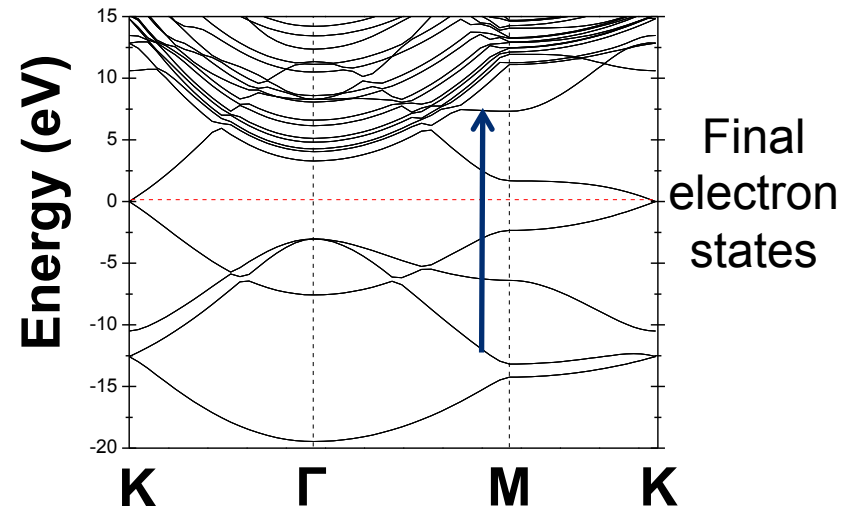
$$\psi(\mathbf{r}) = \sum_{\mathbf{g}} \Psi_{\mathbf{g}} e^{2\pi i(\mathbf{k}+\mathbf{g})\cdot\mathbf{r}}$$



- **NO SOLID STATE BONDING**
- Detailed description of electron propagation (dynamical diffraction, channeling)
- Probe position dependence

## Density Functional Theory


$$\psi = e^{2\pi i\mathbf{k}\cdot\mathbf{r}}$$




- **FULL SOLID STATE BONDING**
- Incident electron is a plane wave- magically appears at the atom
- No incident beam direction
- No detector
- No spatial information

M. P. Prange, M. P. Oxley, M. Varela, S. J. Pennycook, and S. T. Pantelides, "Simulation of Spatially Resolved Electron Energy Loss Near-Edge Structure for Scanning Transmission Electron Microscopy," *Phys Rev Lett*, 109, 246101 (2012).

# Combining Electron Scattering & DFT

core state  final state 

$$M_{if}(\mathbf{Q}_g) = \langle \varphi_i(\mathbf{r}) | \exp(2\pi i \mathbf{Q}_g \cdot \mathbf{r}) | \varphi_f(\mathbf{r}) \rangle$$

  $\mathbf{Q}_g = \mathbf{q} + \mathbf{g} = (\mathbf{k} + \mathbf{g}) - \mathbf{k}'$  momentum transfer

**We measure intensities:**  $|M_{if}(\mathbf{Q}_g)|^2 = M_{if}(\mathbf{Q}_g) M_{if}^*(\mathbf{Q}_h)$

**Mixed Dynamical Form Factor (MDFF) – needed for STEM probe**

$$S^{i,f}(\mathbf{Q}_g, \mathbf{Q}_h, E_{\text{loss}}) = \langle i | \exp(2\pi i \mathbf{Q}_g \cdot \mathbf{r}) | f \rangle \langle f | \exp(-2\pi i \mathbf{Q}_h \cdot \mathbf{r}') | i \rangle$$

**Dynamical Form Factor (DFF) – only applies for plane wave**

$$S^{i,f}(\mathbf{q}, E_{\text{loss}}) = \left| \langle i | \exp(2\pi i \mathbf{q} \cdot \mathbf{r}) | f \rangle \right|^2$$

# EELS: Inelastic Scattering Potential

Fourier Component detector and beam direction

$$\mu_{\mathbf{h},\mathbf{g}}^{i,f} \propto \int_{\text{detector}} k' \frac{S^{i,f}(\mathbf{Q}_{\mathbf{g}}, \mathbf{Q}_{\mathbf{h}}, E_{\text{loss}})}{|\mathbf{Q}_{\mathbf{g}}|^2 |\mathbf{Q}_{\mathbf{h}}|^2} d\Omega_{k'}$$

The full inelastic scattering potential:

$$W^{i,f}(\mathbf{r}, \mathbf{r}', E_{\text{loss}}) \propto \sum_{\mathbf{h},\mathbf{g}} \mu_{\mathbf{h},\mathbf{g}}^{i,f} e^{-2\pi i \mathbf{g} \cdot \mathbf{r}} e^{2\pi i \mathbf{h} \cdot \mathbf{r}'}$$

$$\text{DFF} \Rightarrow \mu_{0,0}^{i,f} \propto \int_{\text{detector}} k' \frac{|\langle i | \exp(2\pi i \mathbf{q} \cdot \mathbf{r}) | f \rangle|^2}{|\mathbf{q}|^4} d\Omega_{k'}$$

Mean value of the inelastic potential



# The Inelastic STEM Image

The inelastic image as a function of probe position  $\mathbf{R}$ .

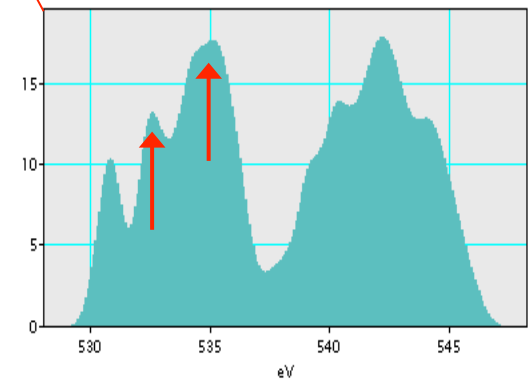
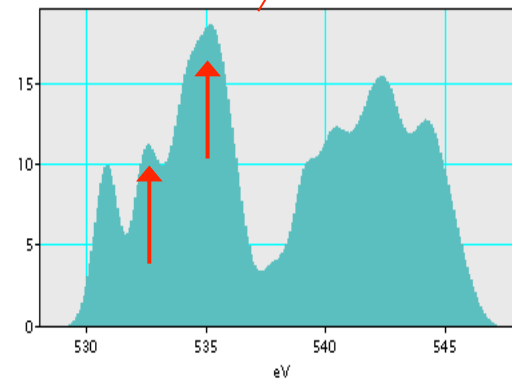
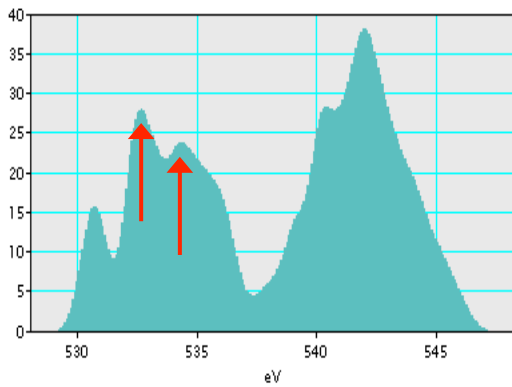
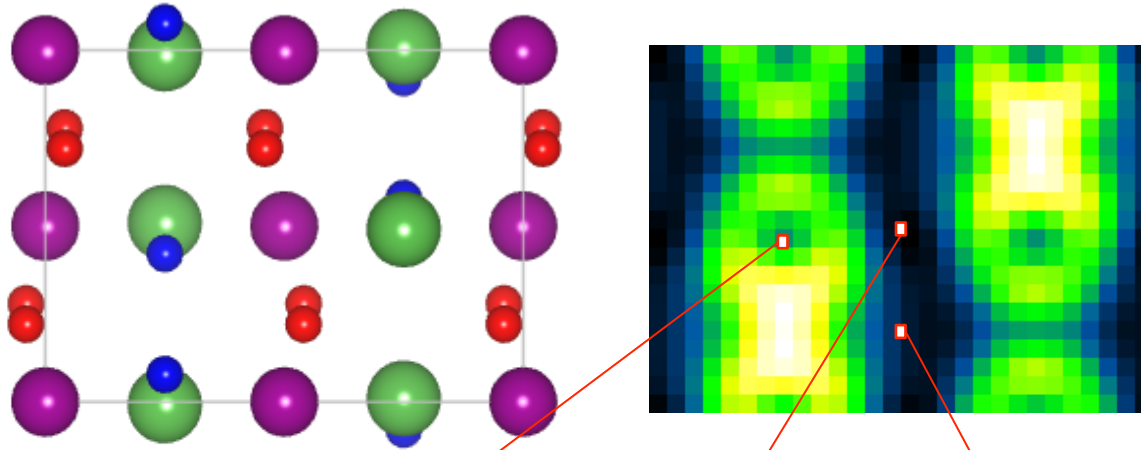
$$I^{i,f}(\mathbf{R}, t) \propto \int_0^t \sum_{\mathbf{h}, \mathbf{g}} \Psi_{\mathbf{h}}^*(\mathbf{R}, \mathbf{r}'_{\perp}, z) \Psi_{\mathbf{g}}(\mathbf{R}, \mathbf{r}_{\perp}, z) \mu_{\mathbf{h}, \mathbf{g}}^{i,f} dz$$

$$\text{DFF} \Rightarrow \mu_{0,0}^{i,f} \Rightarrow I \propto |\Psi_0|^2 \mu_{0,0}^{i,f}$$

The DFF contains no information about the probe position or propagation

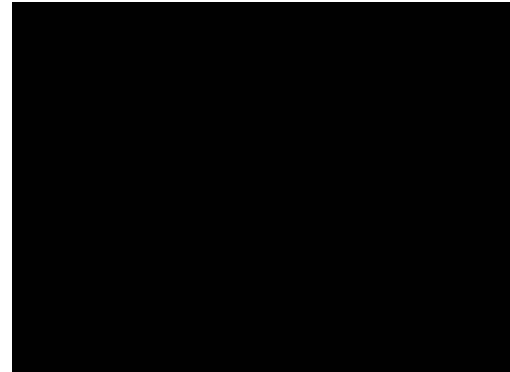
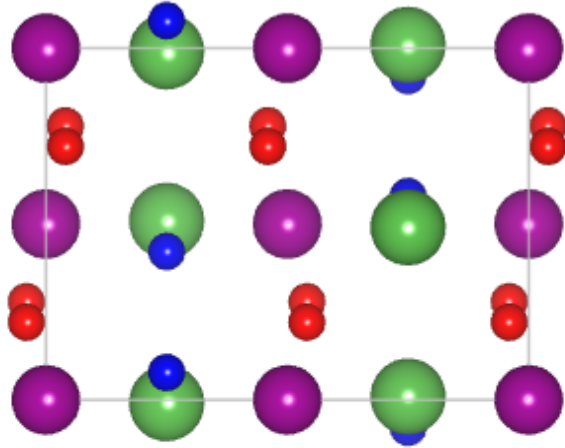
For a given transition, fine structure will not change shape, only intensity

# Near Edge Structure as a Function of Probe Position: O1 signal in LMO

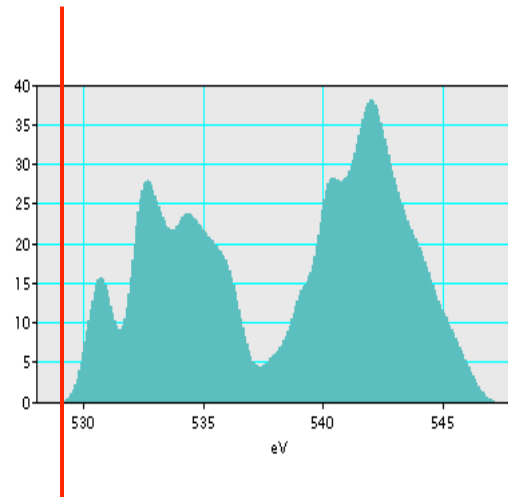


**Fine structure varies with probe position – need full simulation**

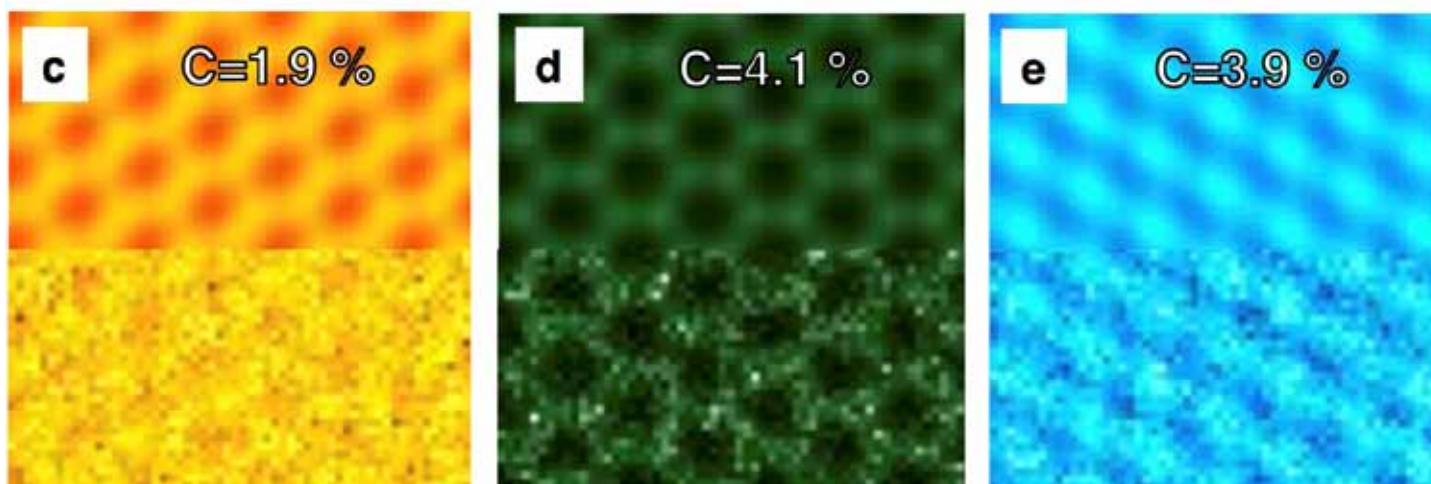
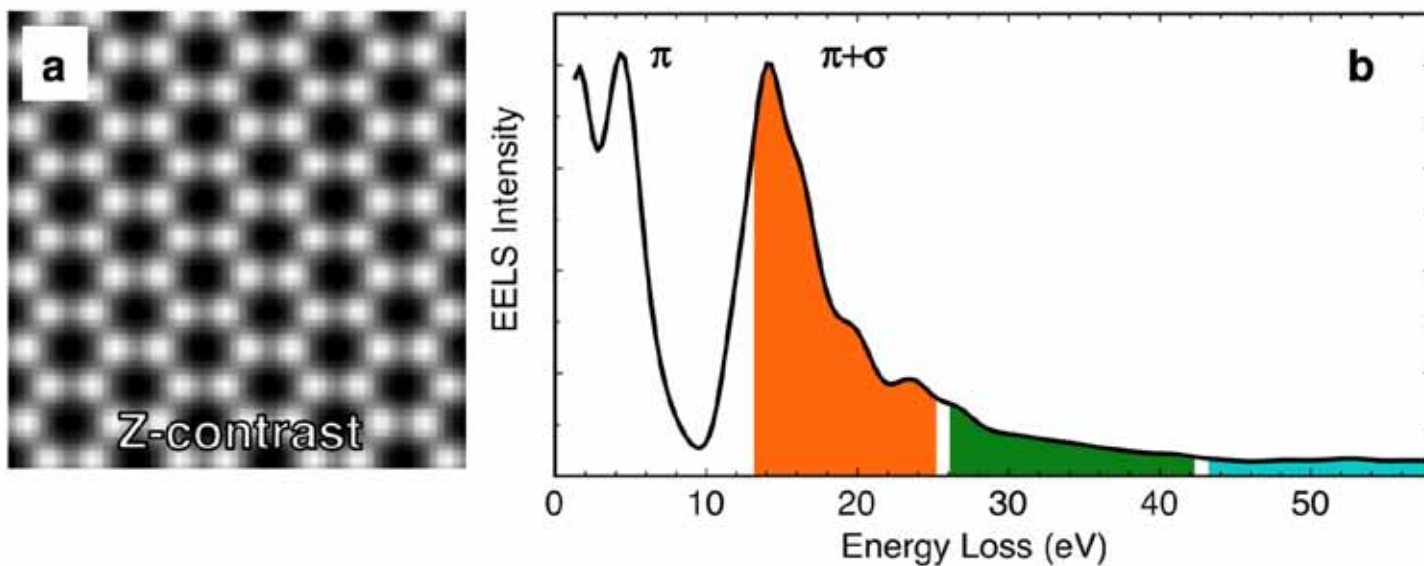
# Near Edge Structure as a Function of Energy: O1 signal



- Fine structure varies with energy
- Delocalization changes for transitions  $< 5$  eV apart
- DFF or simple  $d_{50}$  type expressions fail
- Need full simulation



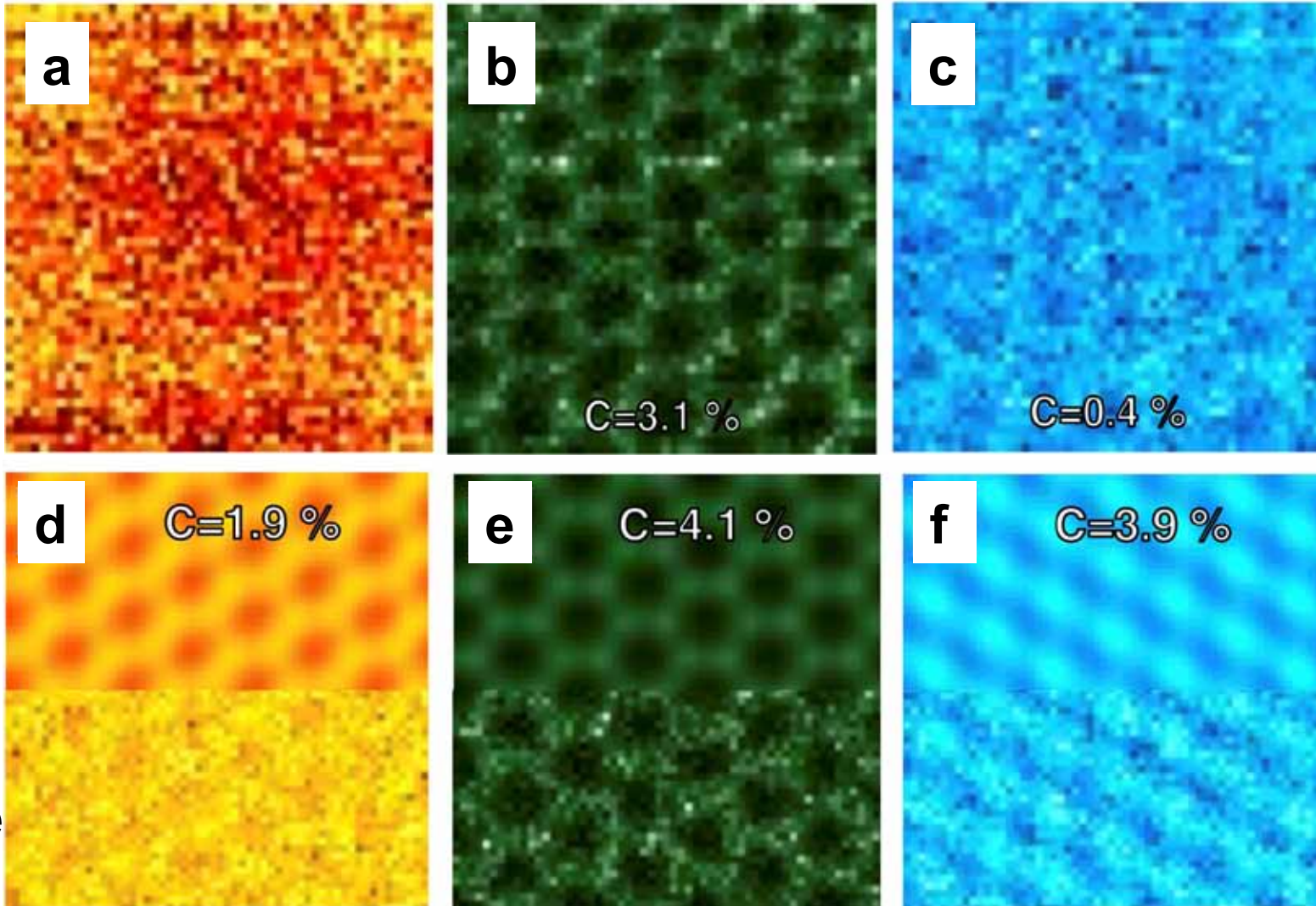
# THEORETICAL VEELS SPECTRA AND MAPS -- GRAPHENE



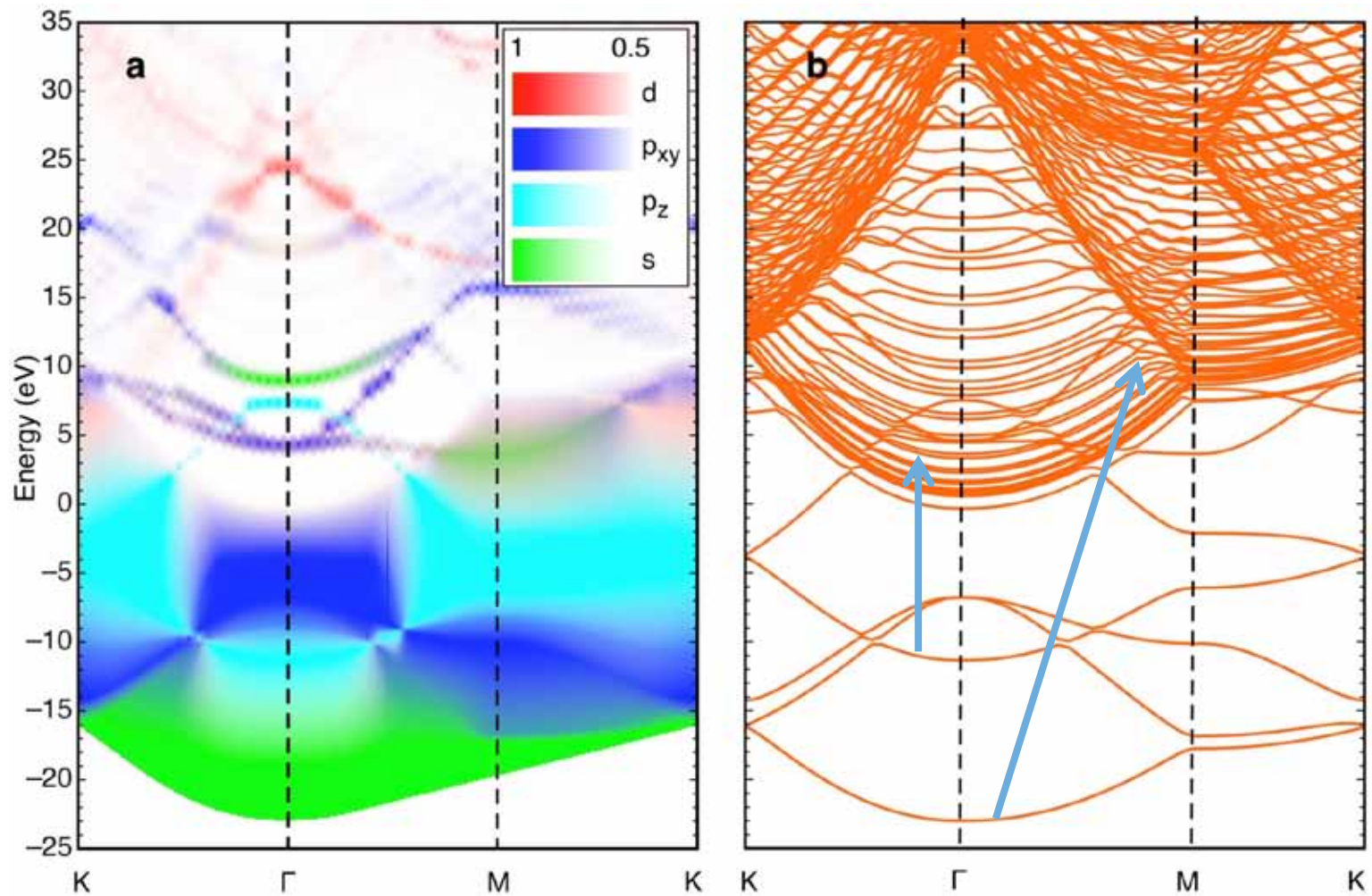
+1%  
noise

COMBINES DFT EXCITATION MATRIX ELEMENTS  
WITH DYNAMICAL SCATTERING

# Experiment



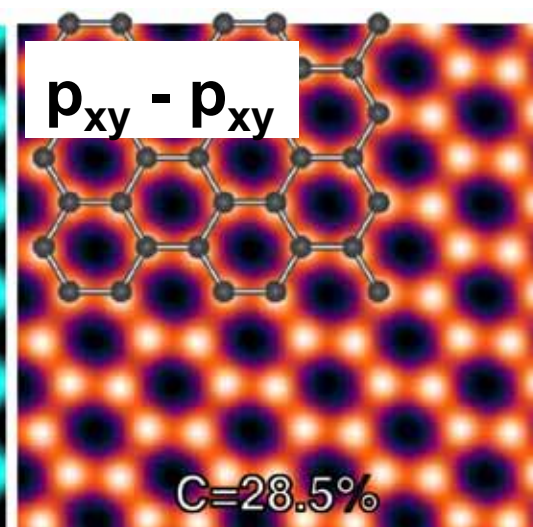
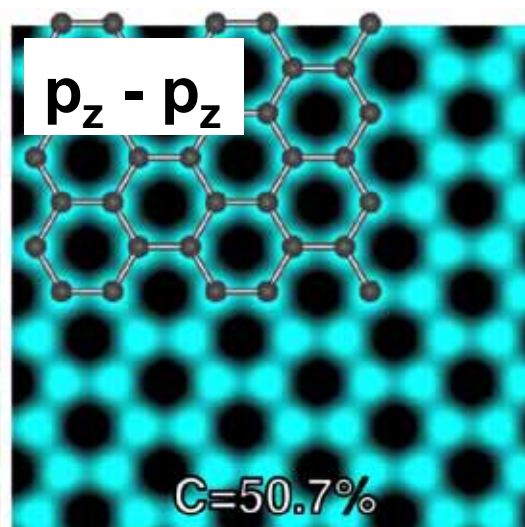
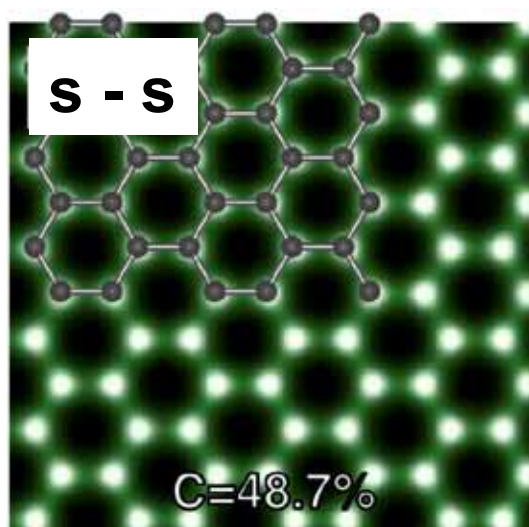
# THE ATOMIC CHARACTER OF GRAPHENE STATES



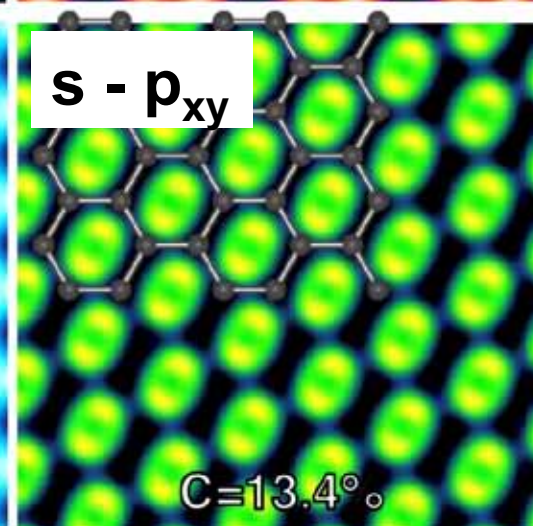
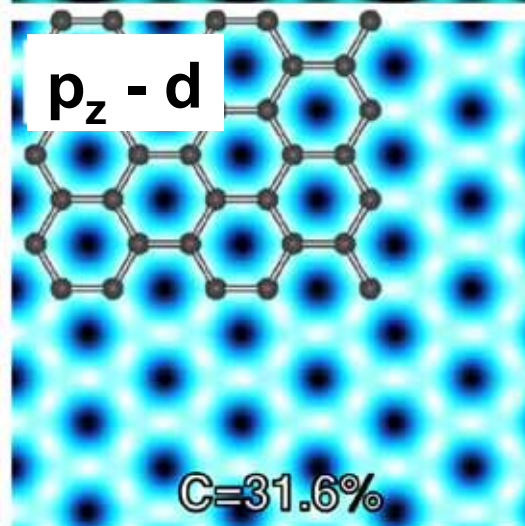
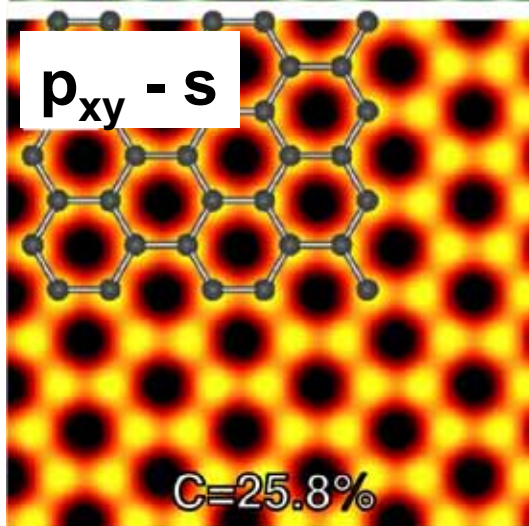
# SELECT ELECTRONIC EXCITATIONS AT $k = 0$

Images from states with maximum atomic character

non-dipole allowed



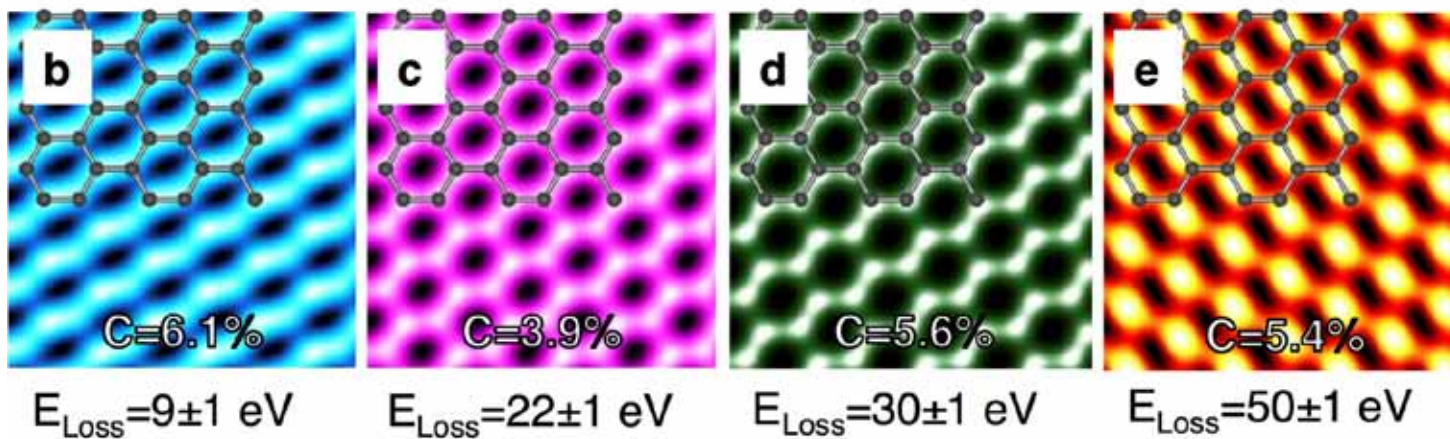
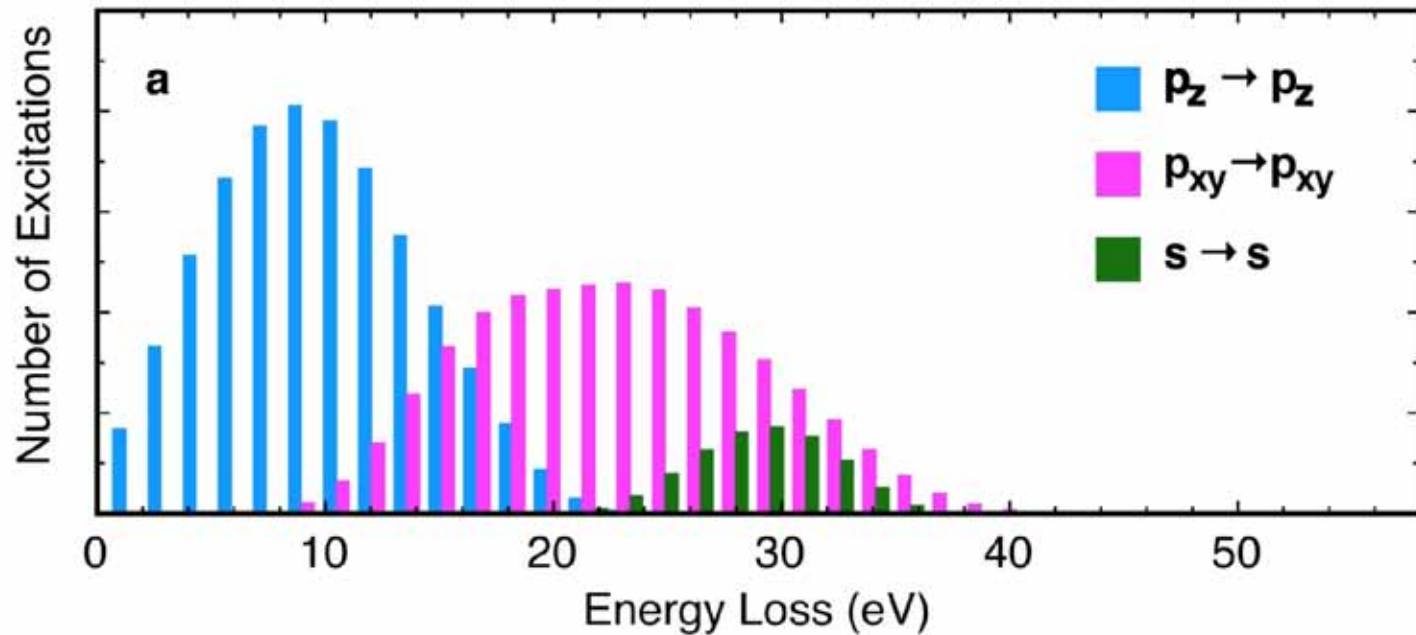
dipole allowed



Reverse contrast

# THE ORIGIN OF GRAPHENE-LIKE IMAGES

All k points





## PROPOSED APPLICATIONS

- **THEORY: DEMONSTRATE THE POWER OF**  
**ATOMICALLY-RESOLVED VEELS**
  - **POINT DEFECTS**
  - **INTERFACES**
  - **MAGNETIC DICHROISM**
- **EXPT: NEED BETTER S/N RATIO AND ENERGY RESOLUTION**

# Bilayer graphene: oriented BLG

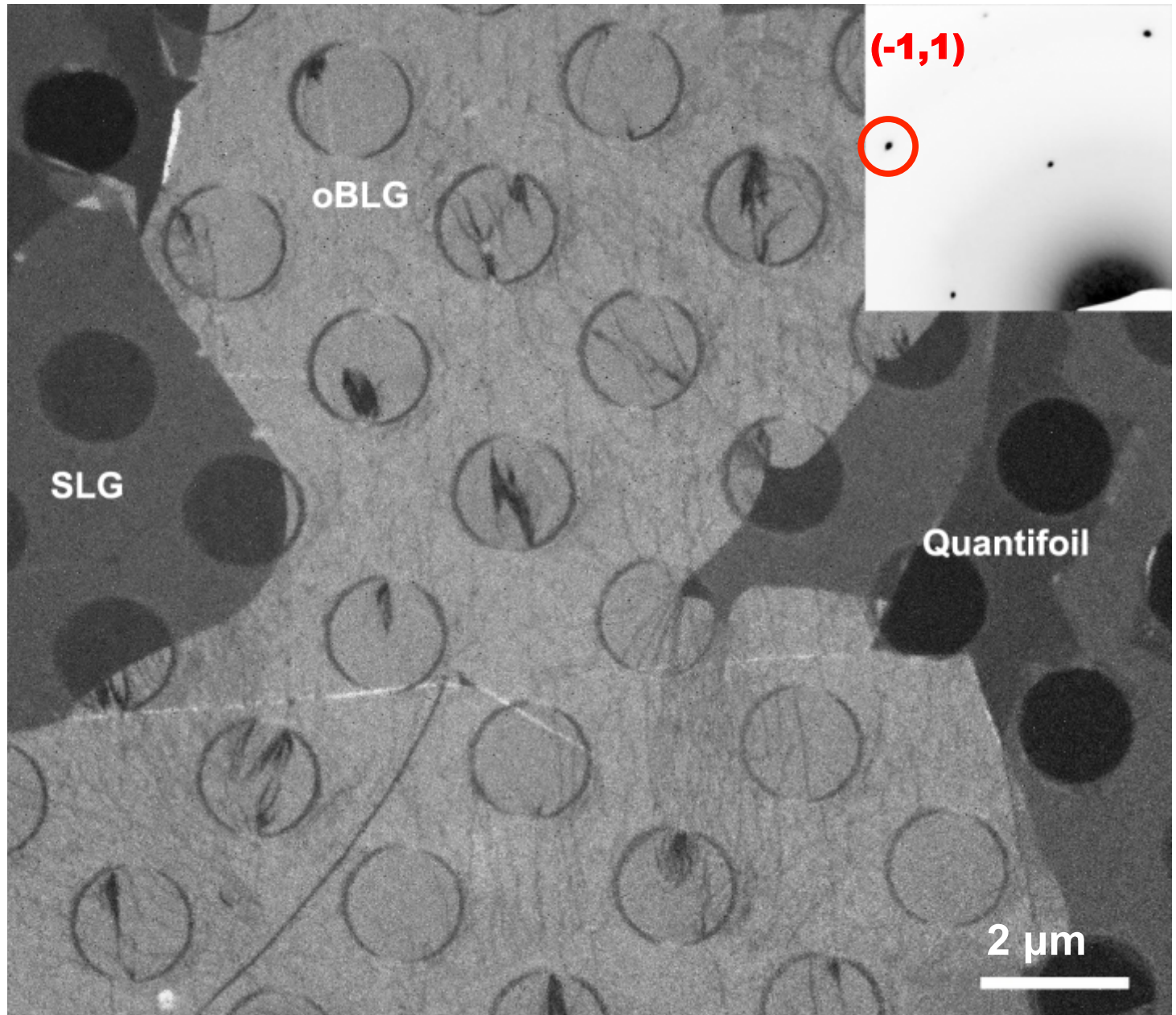
Tilt:  $0^\circ$   
Boundaries  
show dark

Width  $\sim 10$  nm

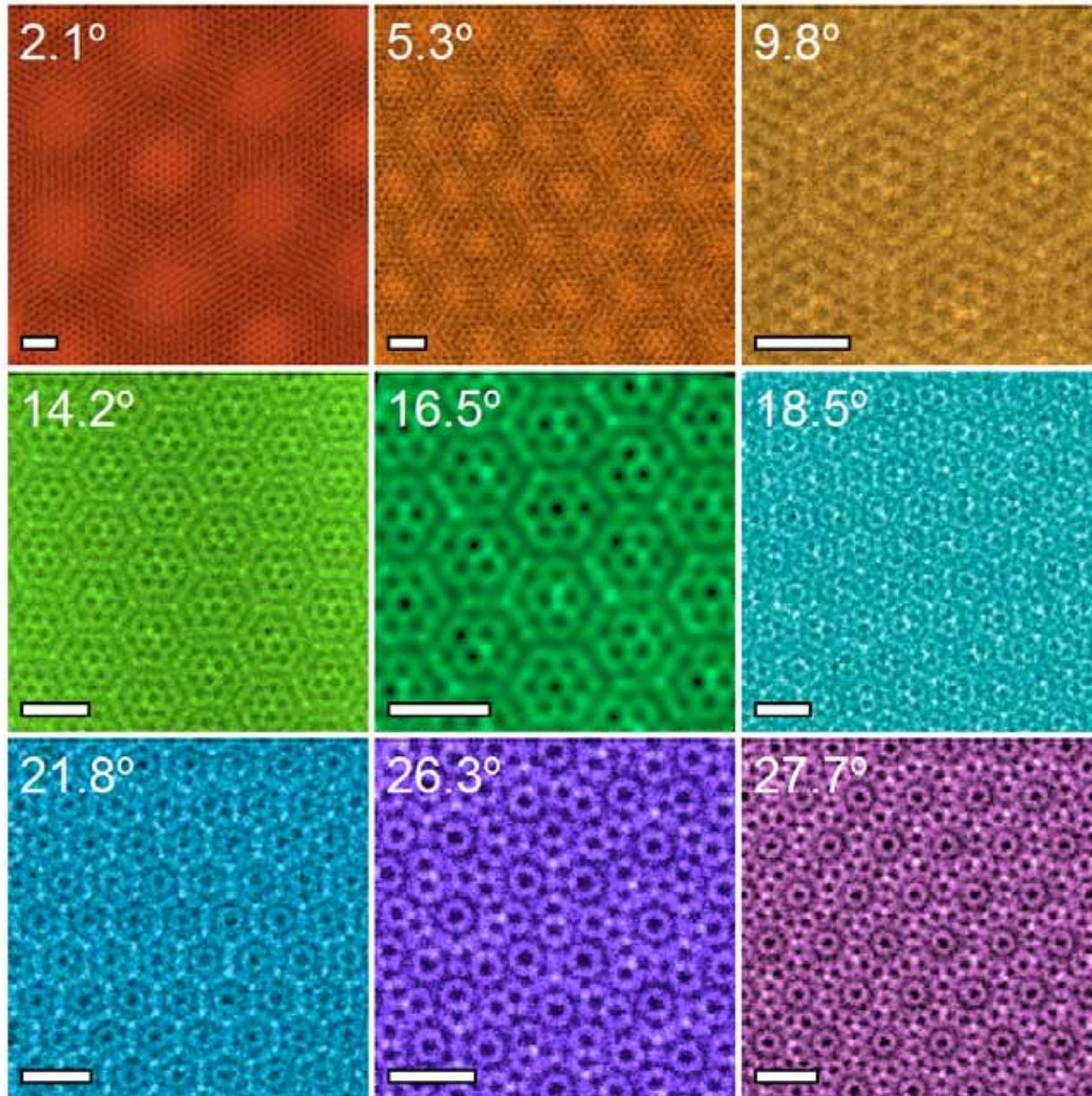
Junhao Lin



Sok Pantelides

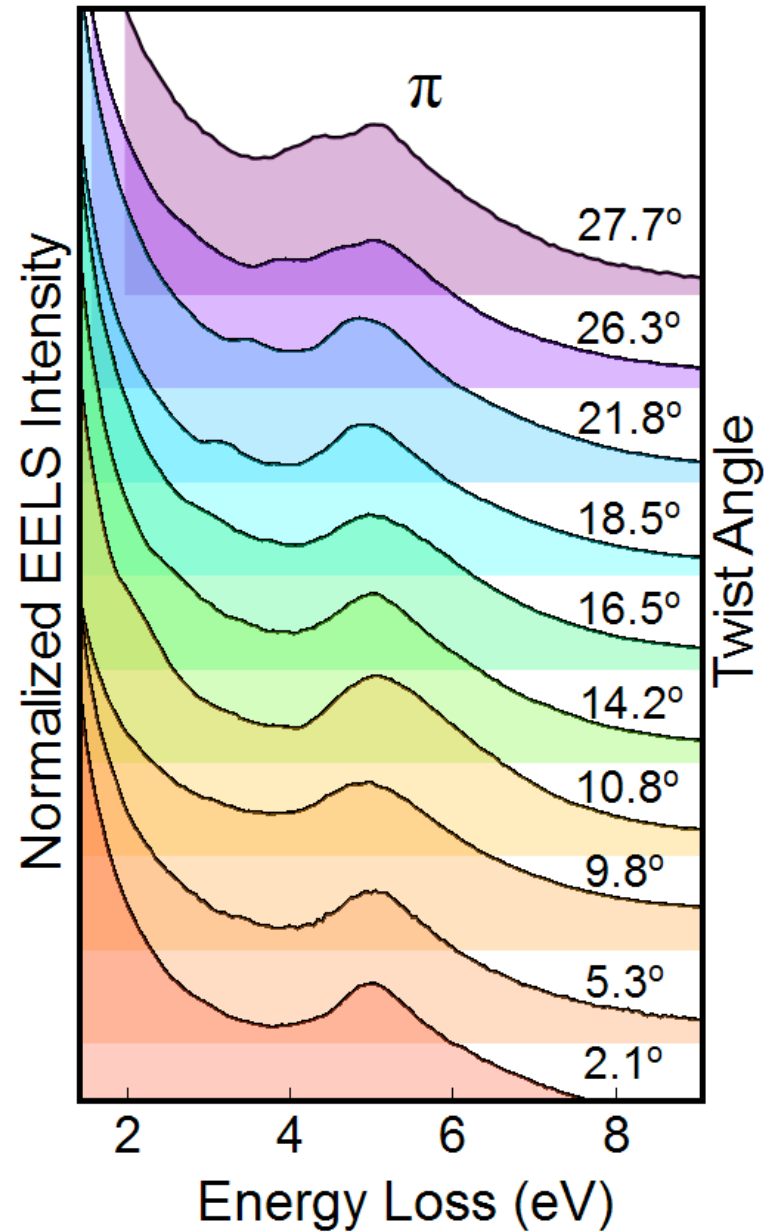


# STEM images of Twisted Bilayer Graphene



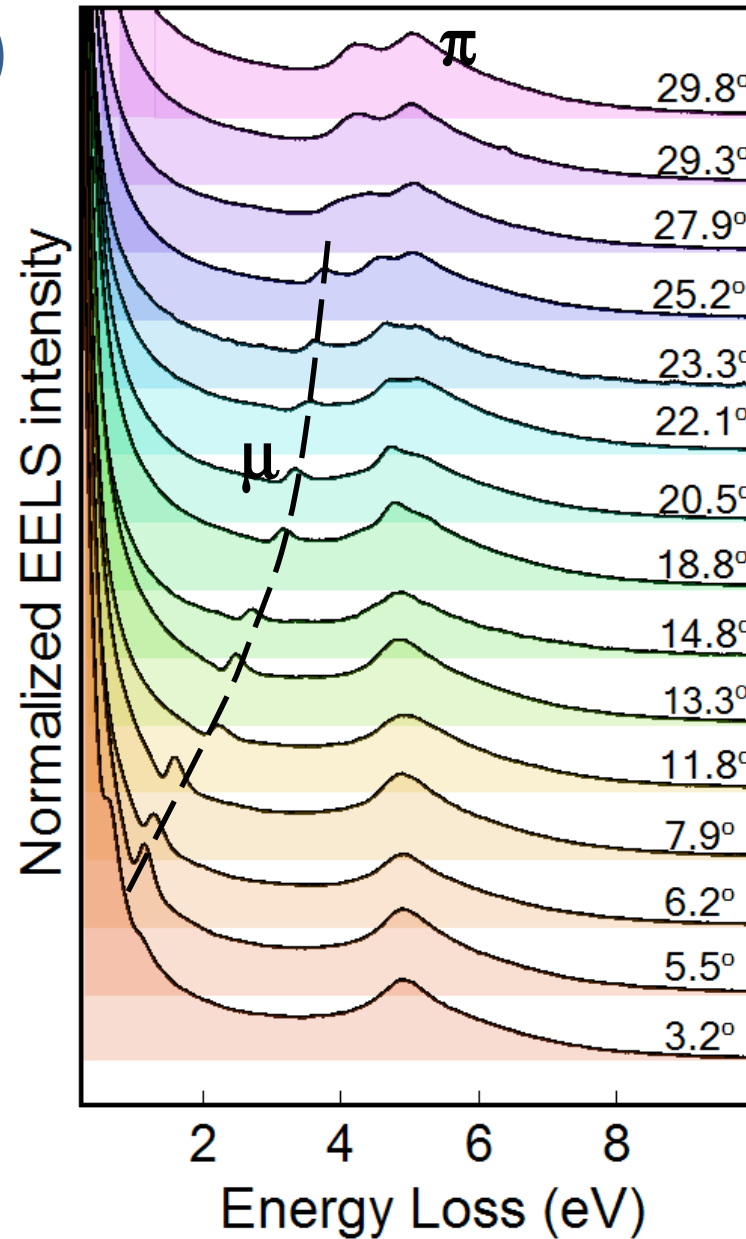
Scale bars  
are 1 nm

# EEL spectra of BLG (350 meV energy)



Experiments @ ORNL

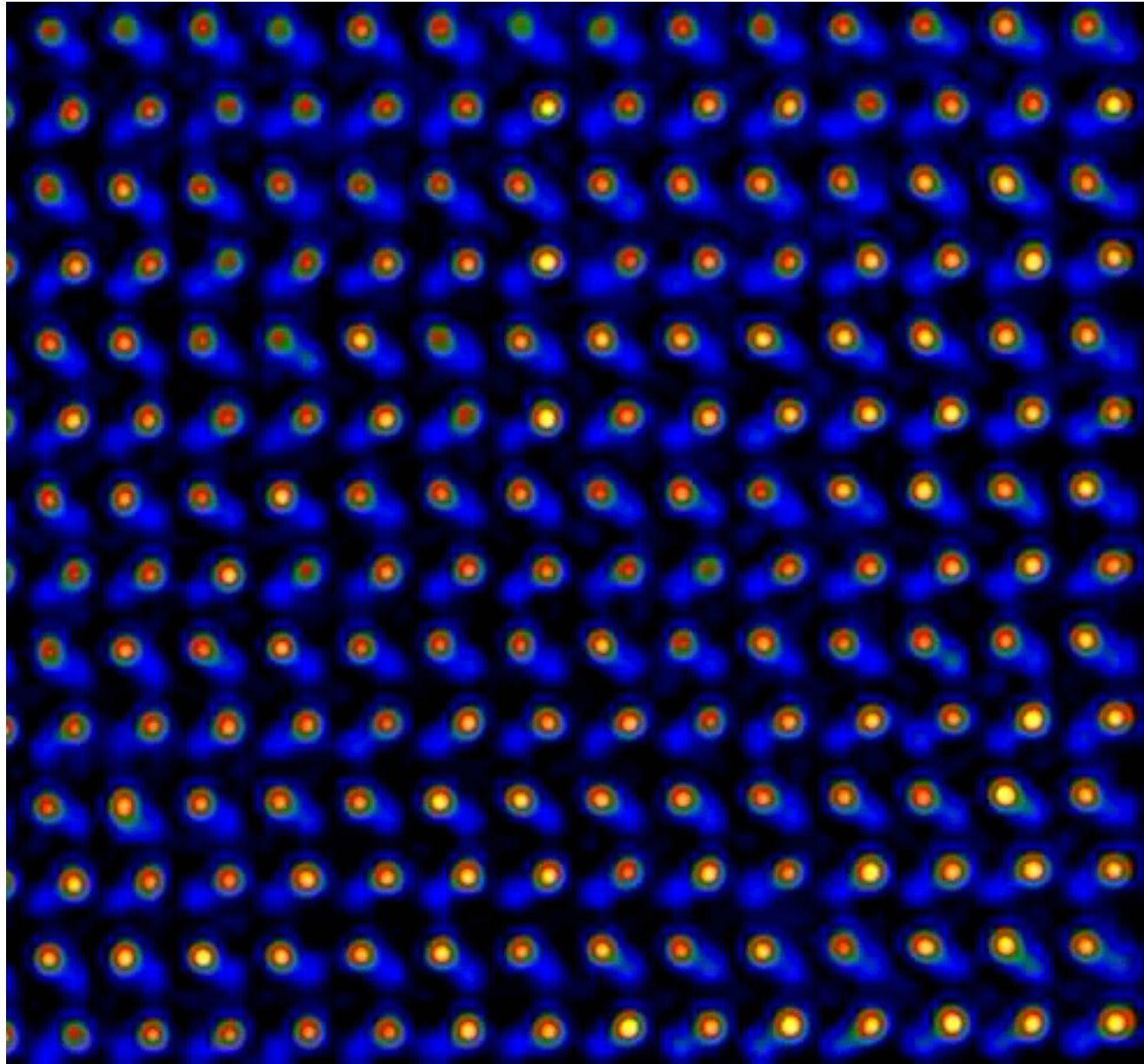
# EEL spectra of BLG (60 meV energy resolution)



Experiments @ ASU

# Dynamics of point defects...

Ce  
in  
AlN

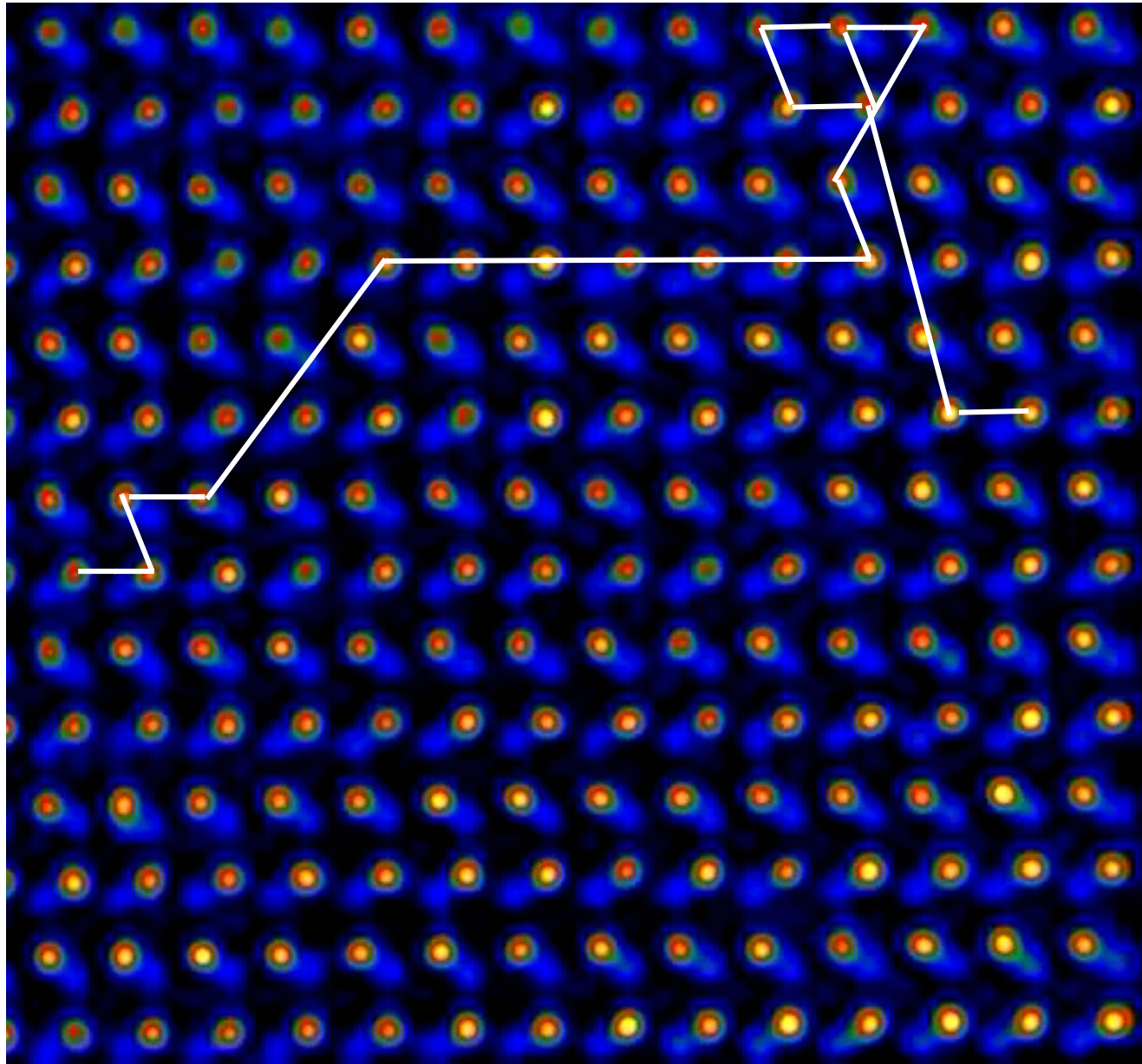


R. Ishikawa and  
A. R. Lupini  
UltraSTEM200

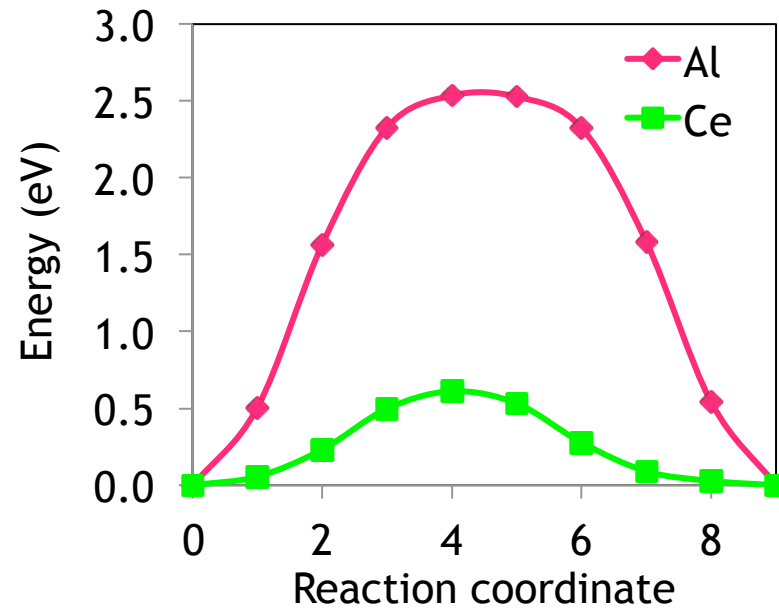
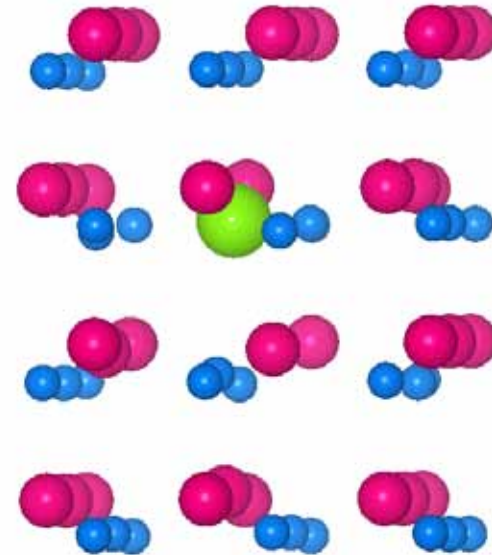
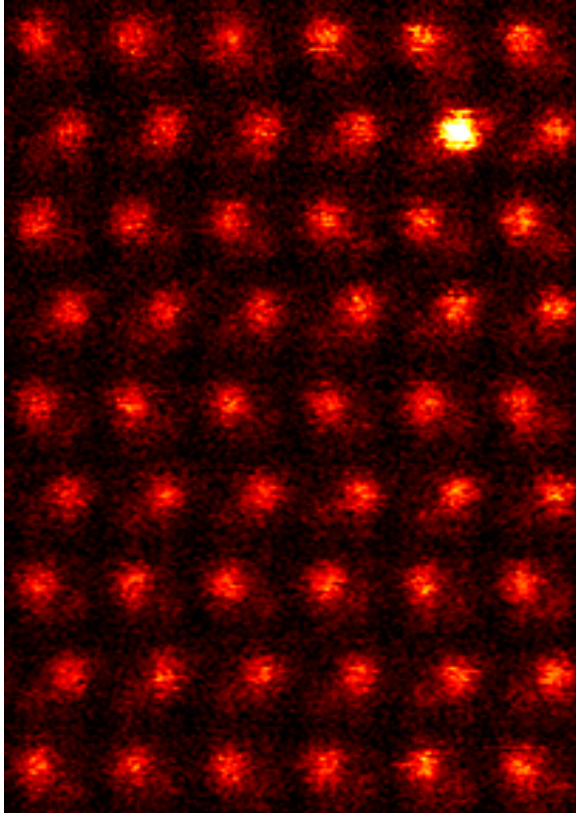
# We can see atomic diffusion...

Ce  
in  
AlN

R. Ishikawa and  
A. R. Lupini  
UltraSTEM200

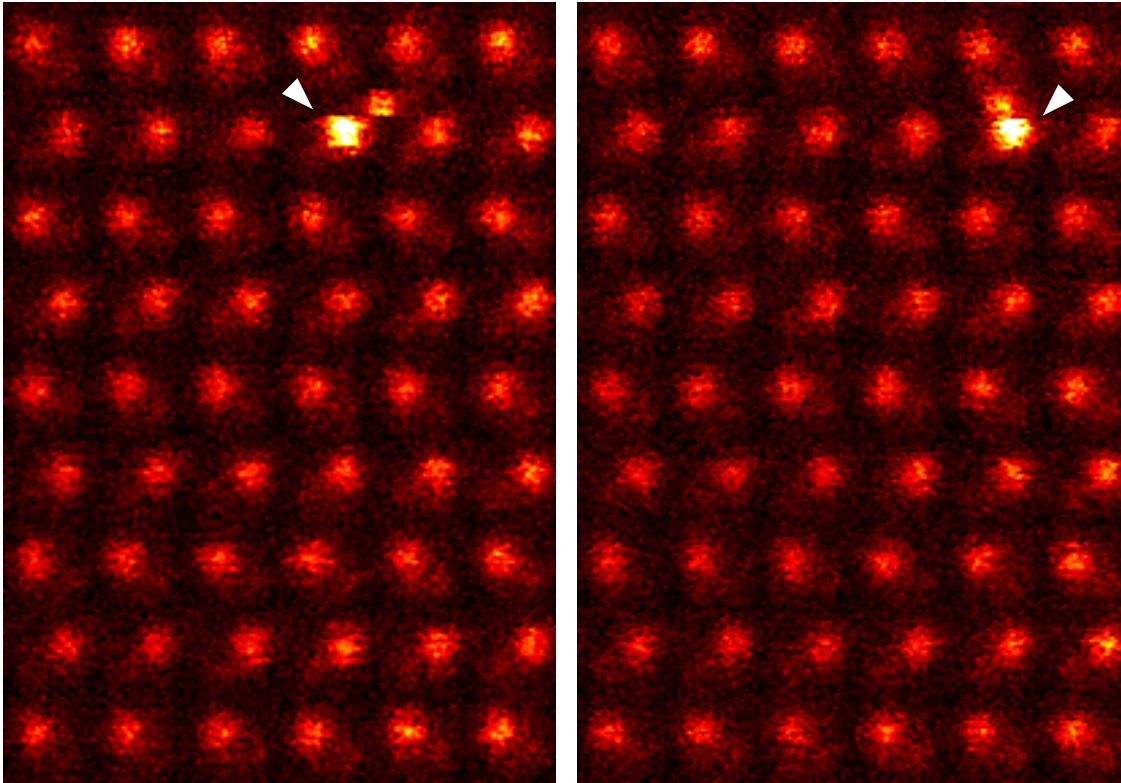


# Correlated vacancy-dopant motion





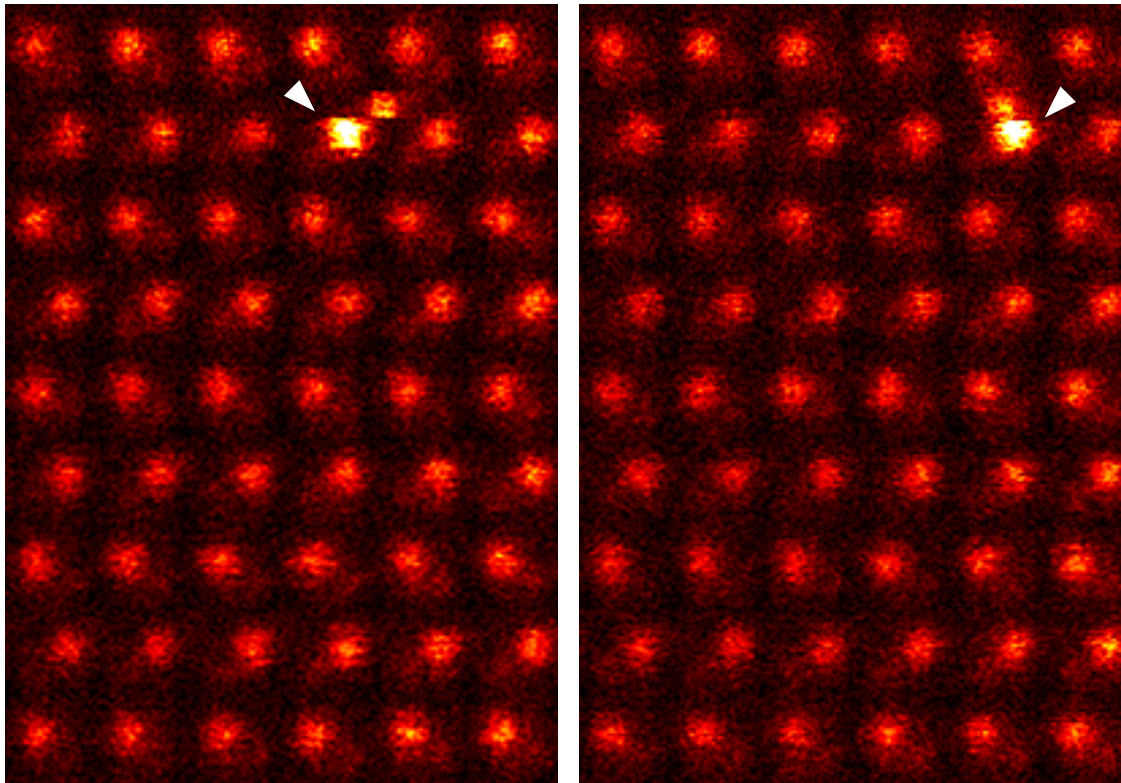
# Interstitial “Kick-out” mechanism



Ce atom at interstitial sites

Simulate higher-temperature  
diffusion events!

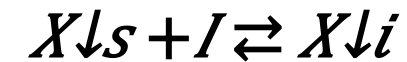
# Interstitial “Kick-out” mechanism



Ce atom at interstitial sites

Simulate higher-temperature diffusion events!

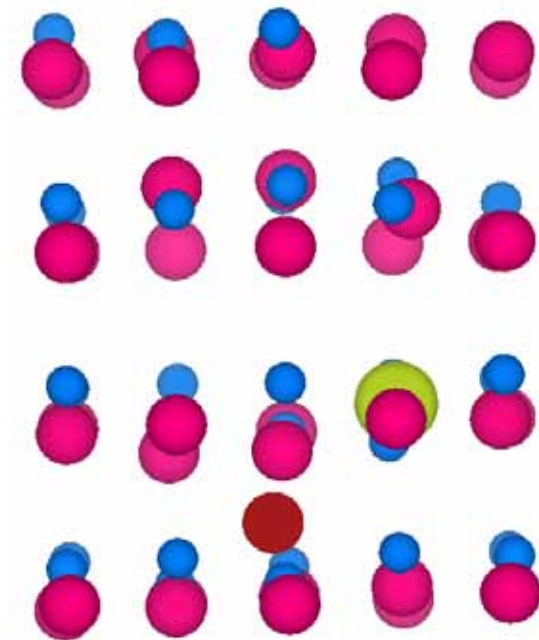
Kick-out mechanism



$X_s$  : Substitutional dopant

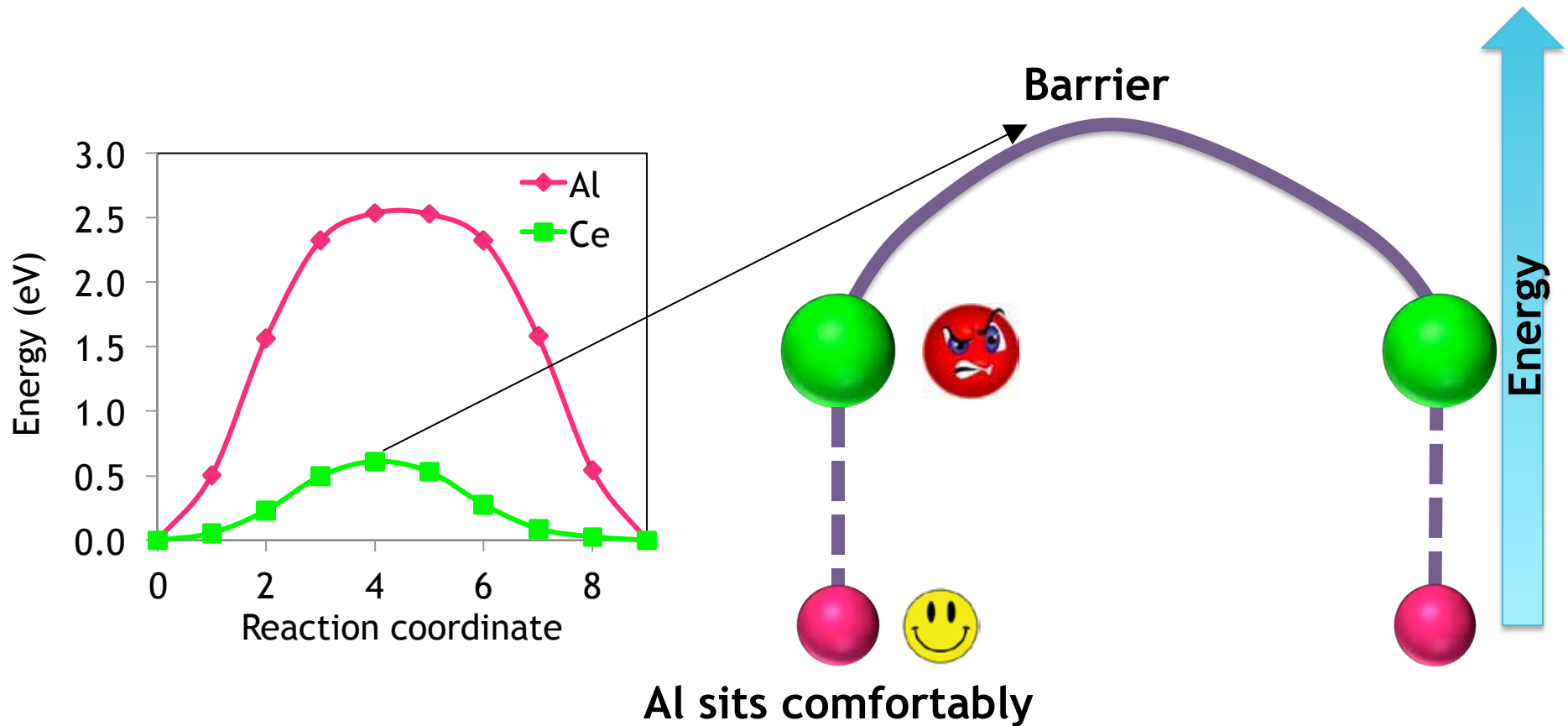
$X_i$  : Interstitial dopant

$I$  : Self-interstitial



Barrier : 3.7 eV

# Why is the Ce barrier so low?



Large Ce atom is highly strained in a small AlN lattice

# nature

International weekly journal of science

## Hasten high resolution

Build precision microscopes to map atoms, say  
Stephen J. Pennycook and Sergei V. Kalinin.

***Feynman's goal = 0.1Å!***  
***"just look at the thing!"***

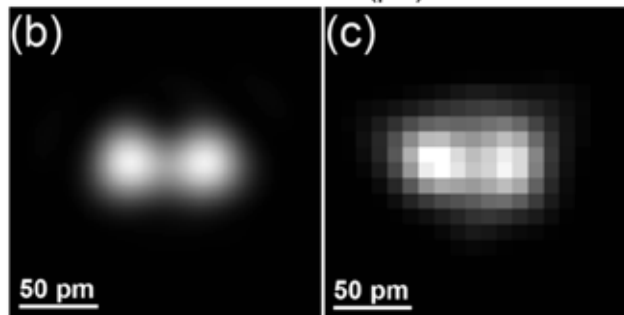
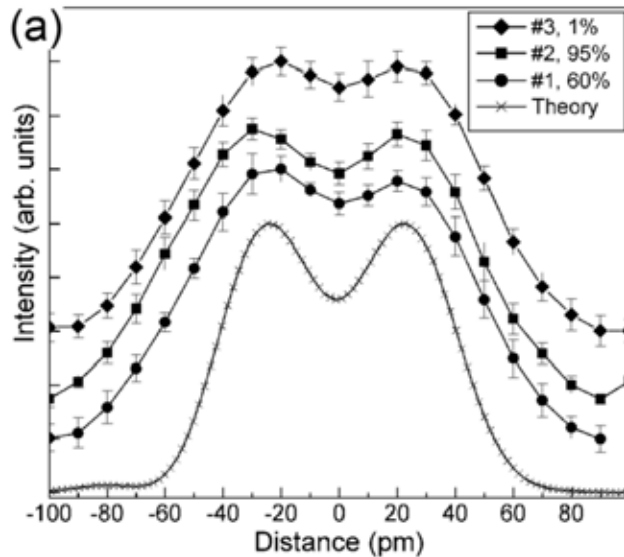


*"I would like to try and impress upon you while I am talking about all of these things on a small scale, the importance of improving the electron microscope by a hundred times. It is not impossible; it is not against the laws of diffraction of the electron."*

1. S. J. Pennycook and S. V. Kalinin, "Hasten high resolution," *Nature*, 515, 487–488 (2014).
2. S. J. Pennycook, "Fulfilling Feynman's dream: "Make the electron microscope 100 times better"—Are we there yet?" *MRS Bull*, 40, 71–78 (2015).

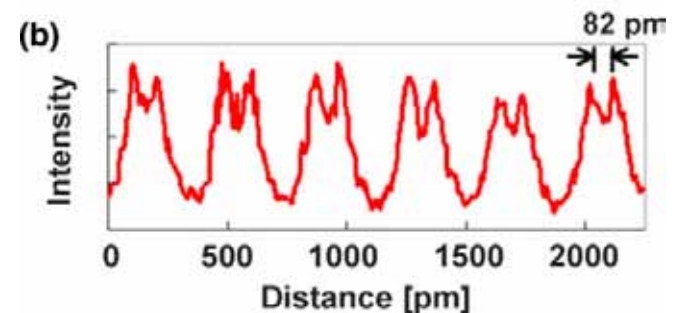
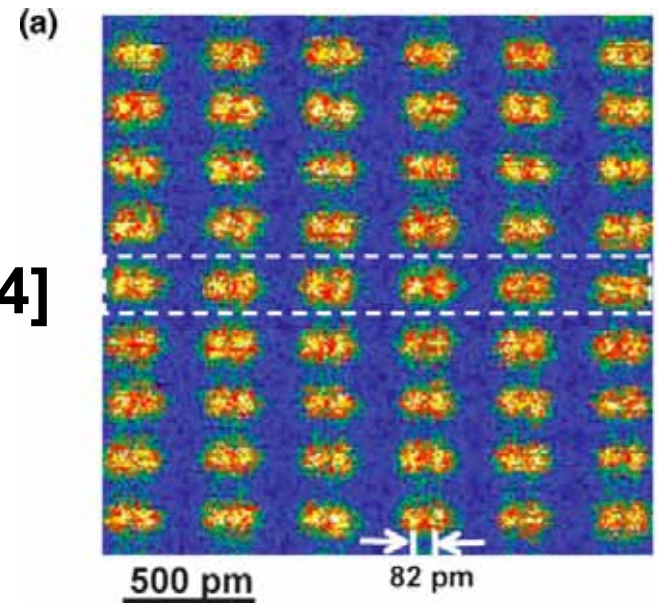
# STEM achieves 0.5 Å resolution

5<sup>th</sup> order aberration correction, 300 kV



Erni, R. et al., *Phys Rev Lett* 2009, 102, 96101;  
TEAM project

Ge [114]



Sawada, H. et al., *J Electron Microscop* 2009, 58, 357;  
CREST project

# But depth resolution is still on the **nm-scale**

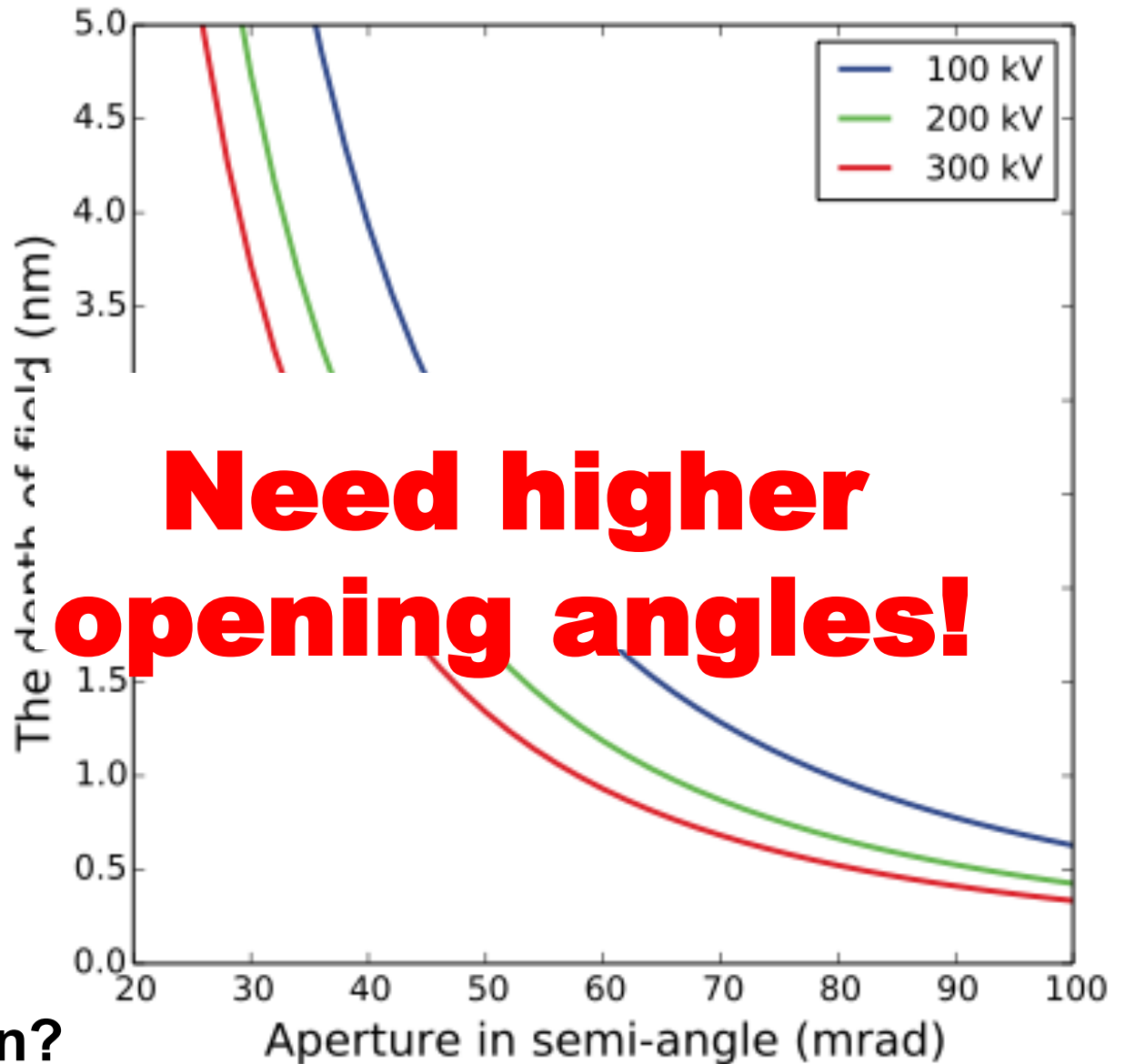
$$d_{x,y} = 0.61 \frac{\lambda}{\theta}$$

$$d_z = 2 \frac{\lambda}{\theta^2}$$

30 mrad:  
few nm @ 300 kV

100 mrad:  
4 Å @ 300 kV

Sub-unit cell resolution?

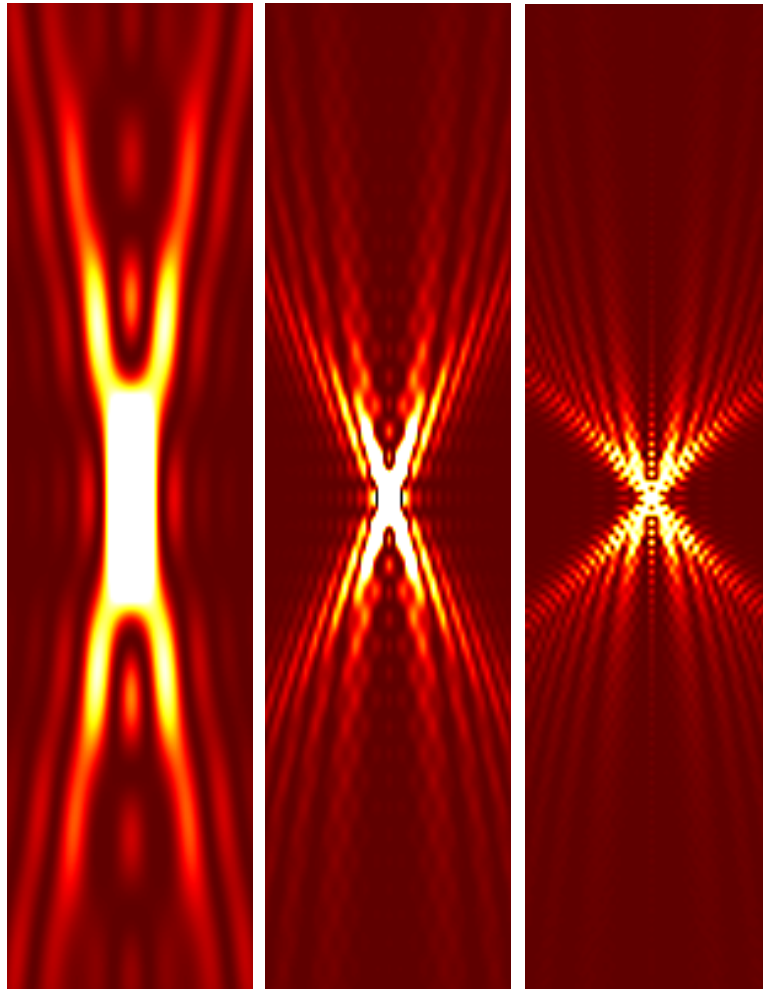


# Depth resolution at unit cell level...

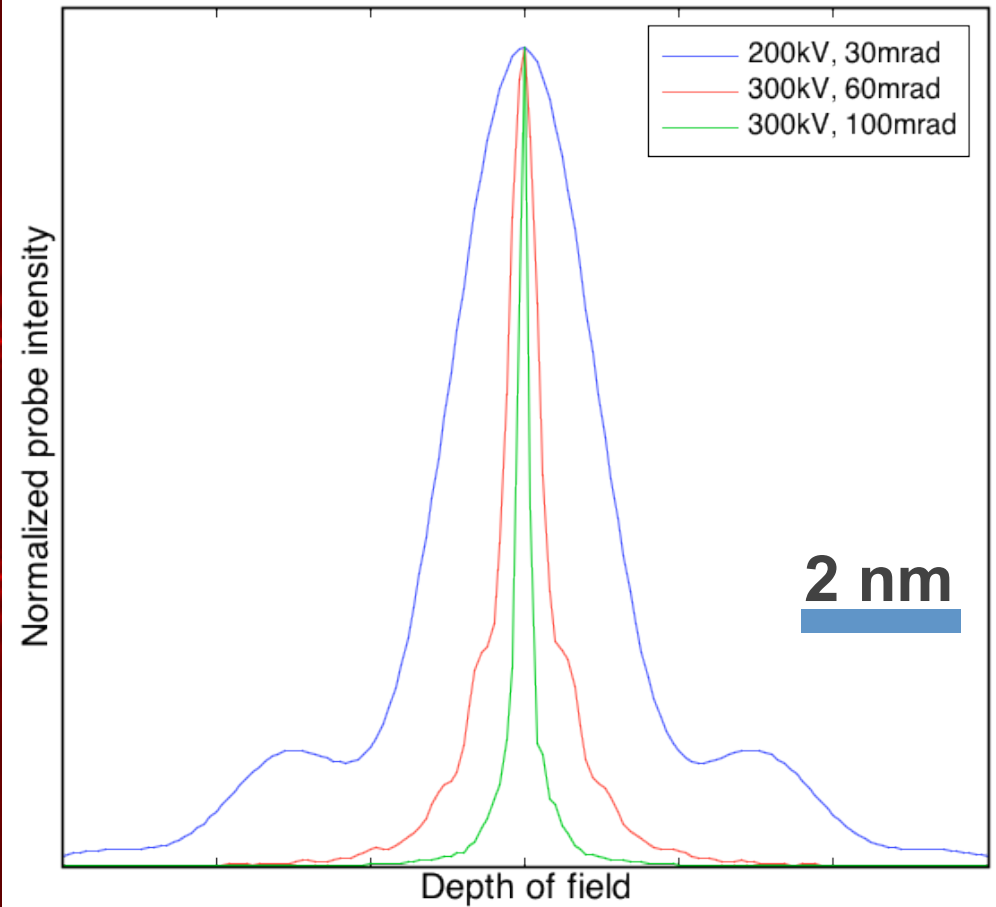
200 kV 300 kV 300 kV  
30 mrad 60 mrad 100 mrad

+10 nm

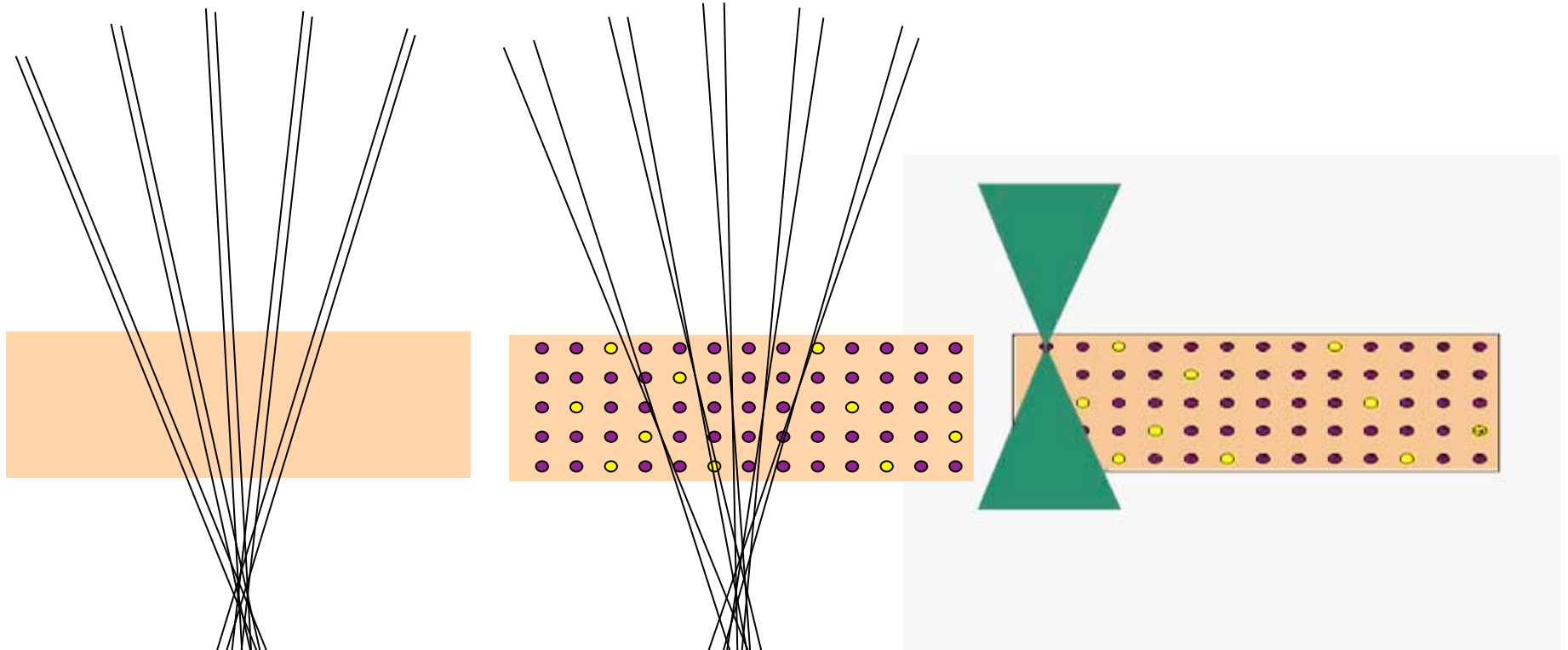
defocus



-10 nm



# Resolution in the 3<sup>rd</sup> dimension?



**Tomography  
requires direct  
projection eg HAADF  
at low resolution**

**At high resolution direct  
projection breaks down  
Stability?**

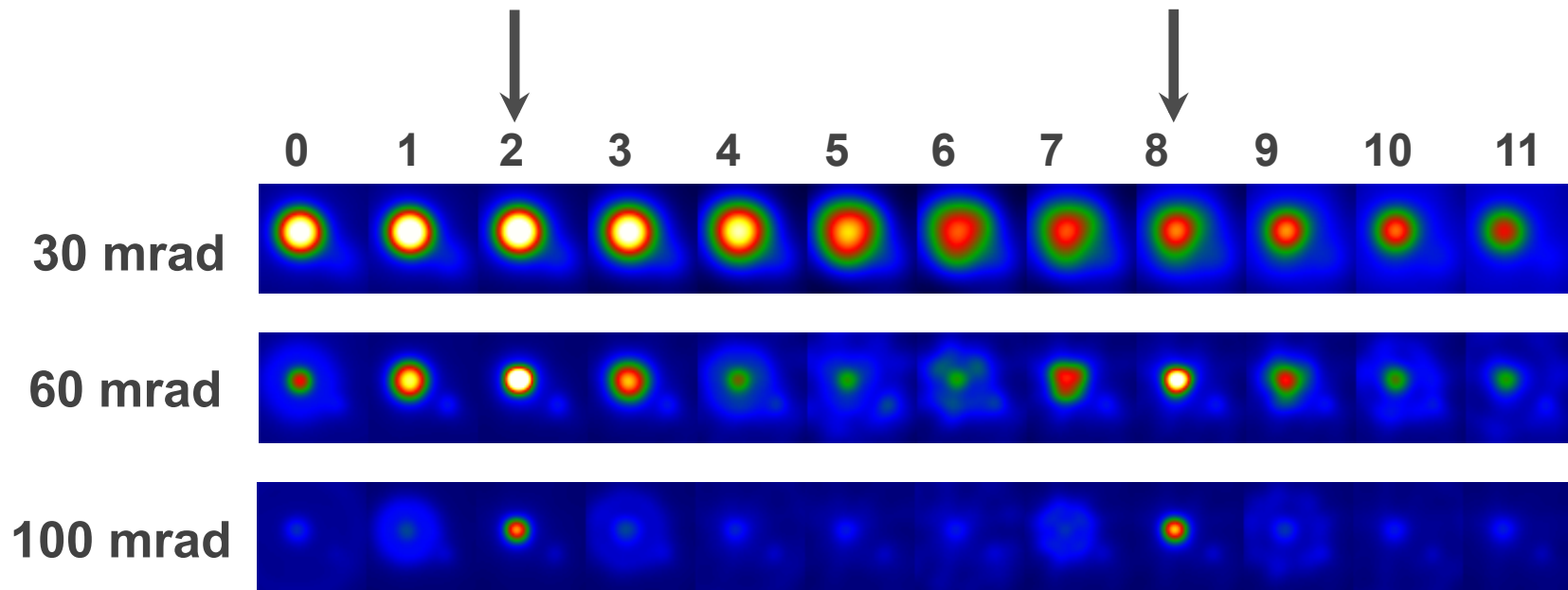
**Natural transition  
to depth  
sectioning**



# Substitutional Ce dopants in AlN

Ce @ 2.2 & 7.8 nm depth

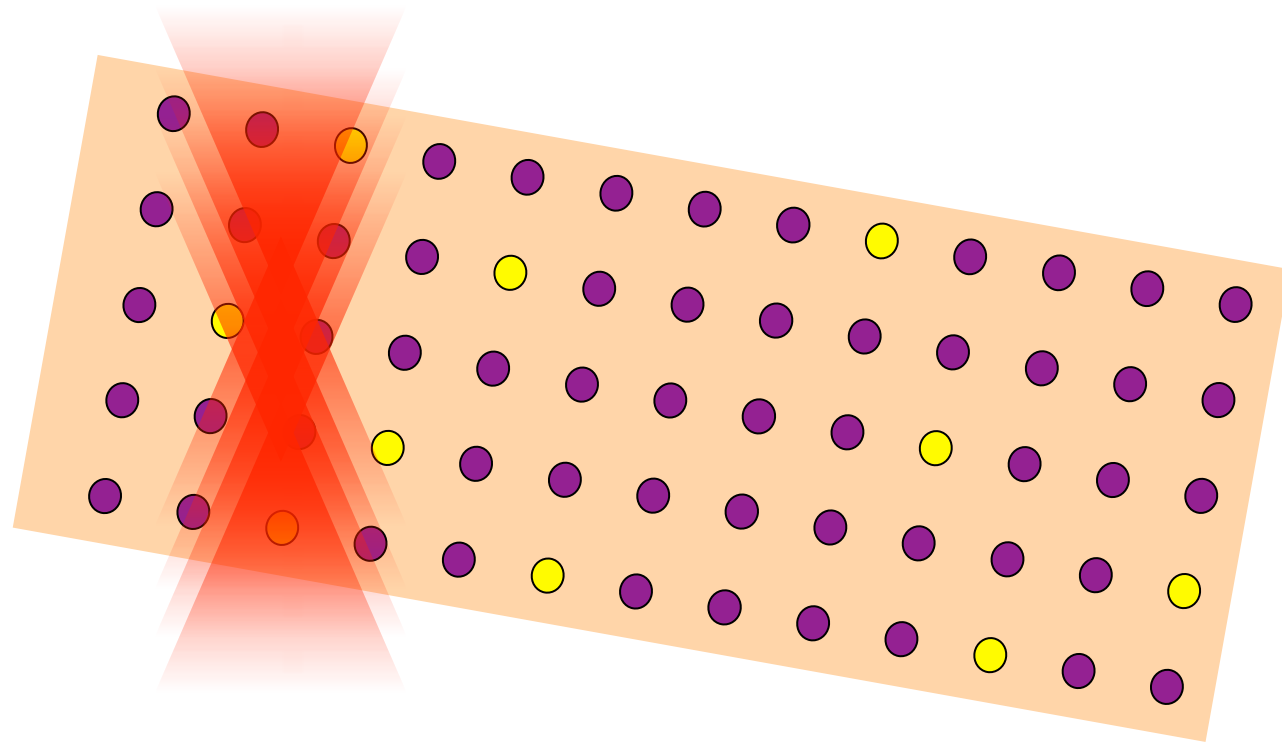
defocus (nm)



thickness = 12.4 nm

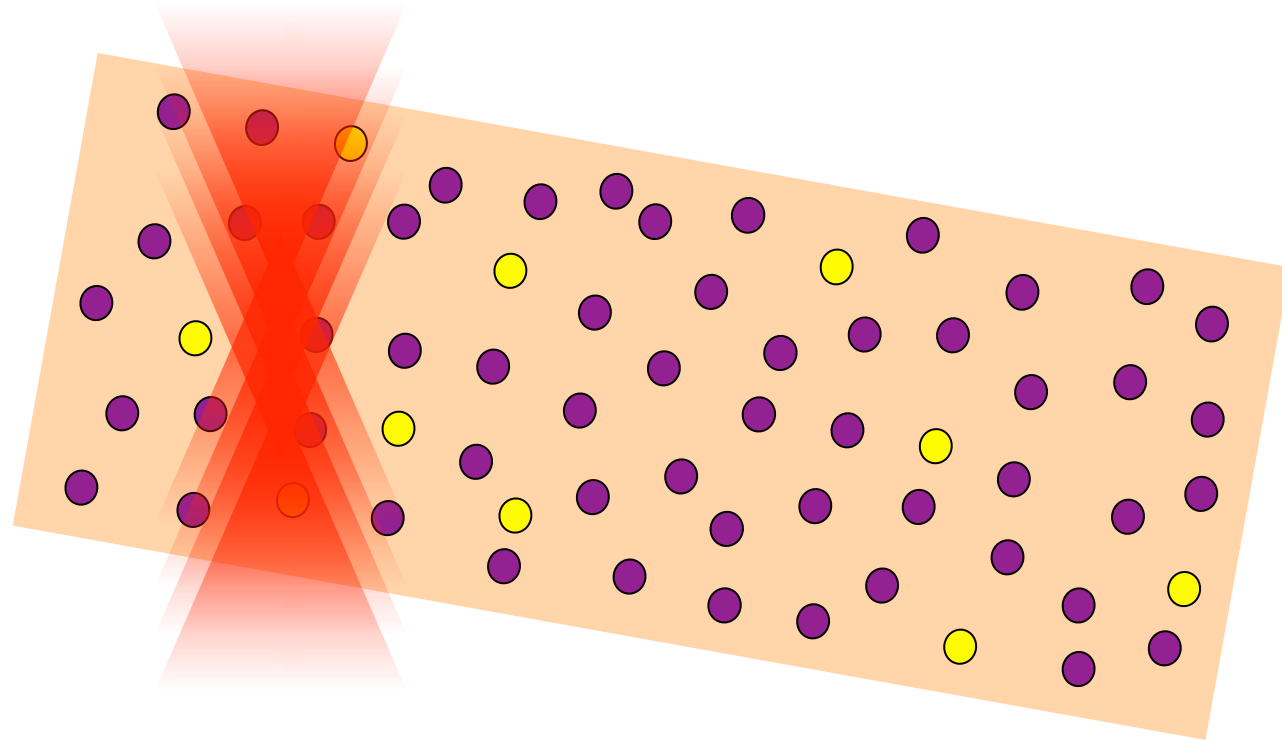
Simulation by Ryo Ishikawa

**No need to align the sample!**



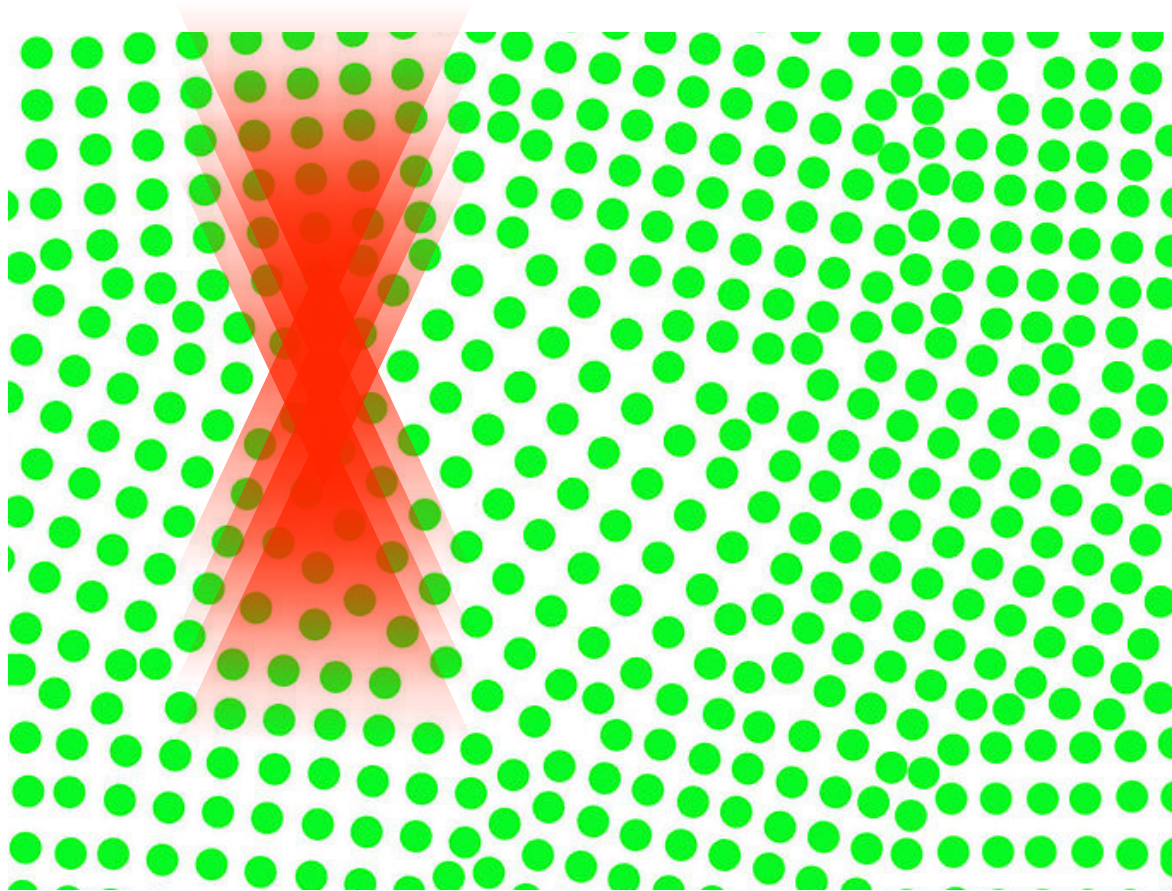
**3D atomic structure!**

**We could look at glasses!**



**3D atomic structure!**

**We could see grain boundaries in polycrystalline materials!**



**3D atomic structure!**

# There's still plenty to see at the bottom!

- Next generation corrector
  - 300 kV, CFEG, 60-100 mrad probe angle
  - ~0.2-0.12 Å diffraction limit
- Improved signal to noise ratio
- Optical properties with atomic sensitivity
- Improved precision of atomic positions
- 3D atomic resolution

