



# Progress towards Low Vacuum Critical Dimension Metrology

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**Frontiers of Characterization**  
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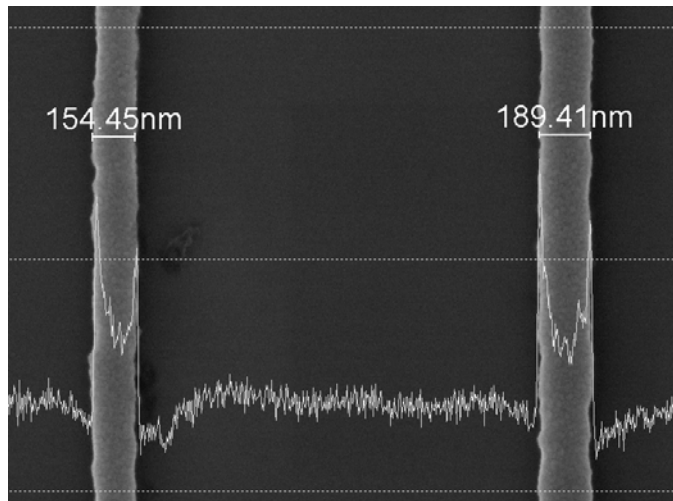
- **Introduction**
  - Critical Dimension Metrology (CD SEM)
  - Necessity for Modeling
- **Low Vacuum CD Metrology Approach**
  - Charge Control
  - Contamination Control
- **Principles of Low Vacuum SEM**
- **Analytical Modeling of Amplification and Noise Characteristics**
  - ESD (Environmental Secondary Electron Detector)
  - Helix (Magnetic Immersion Lens Detector)
- **Summary & Pending Work**



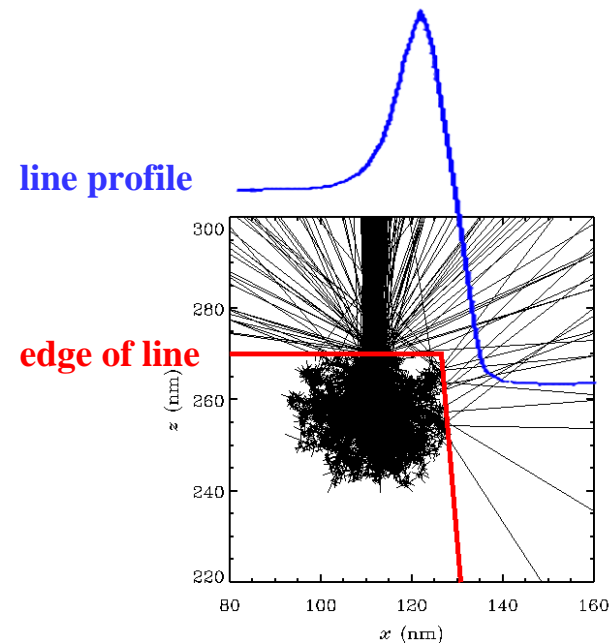
- **Scanning electron microscope methods are some of the most prevalent technologies for critical dimension (CD) measurement.**
  - High resolution imaging (spatial res.  $\sim 1\text{nm}$ , measurement res.  $\sim 1\text{nm}$ )
  - Acceptable throughput and speed (typically 1 wafer per  $\sim 5\text{min}$ , 17 points on wafer)
  - Site specific measurements
  - Well-established procedures for interpreting the results
- **CD's are rapidly approaching the limitations of conventional CD-SEM imposed by:**
  - Charging
  - Contamination
  - Specimen interaction volume
  - Secondary electron mean-free-path
- **Low Vacuum SEM technology can address charging and contamination without sacrificing performance**



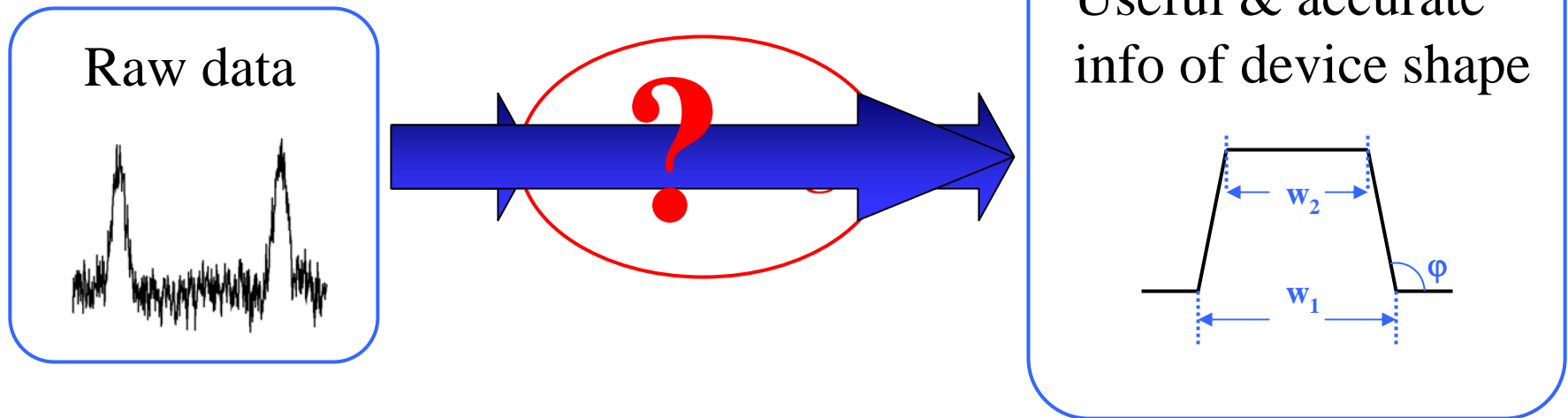
- SEM information is distilled into an intensity profile
- The peaks indicate the position of the edges
- Simulation codes correlate peak position and shape with actual feature shape
  - Edge “blooming” effect



chrome-on-quartz photolithographic mask



MC simulation of electron trajectories near the edge of a line and corresponding line profile



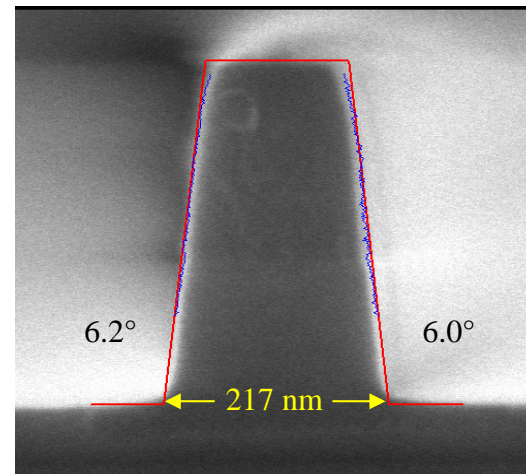
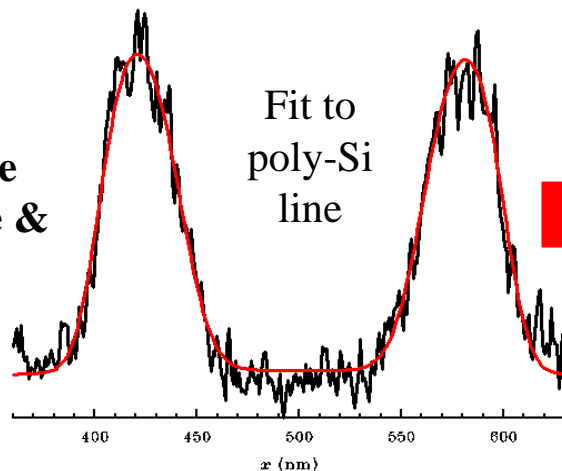
**ITRS – Need for standardization of CD process,**

i.e. “robust conversion of massive quantities of raw data to information useful for enhancing the yield of the semiconductor manufacturing process”



- **Analytical descriptions** are implemented for **all physical processes**
  - **Electron beam/specimen interactions**
- **Instrumentation effects** are included
- Software that can reach, and if needed interpolate between a library of precomputed line profiles for model geometries
- Non-linear least squares algorithm solves for the particular set of parameters that produce the best least squares match between measured and calculated images

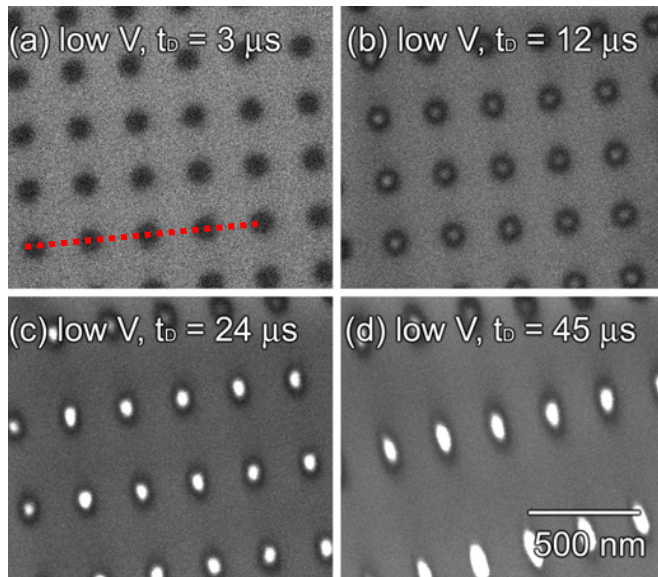
**Top down image  
intensity profile &  
Simulation fit**



**X-sectional image &  
X-sectional fitted  
shape**

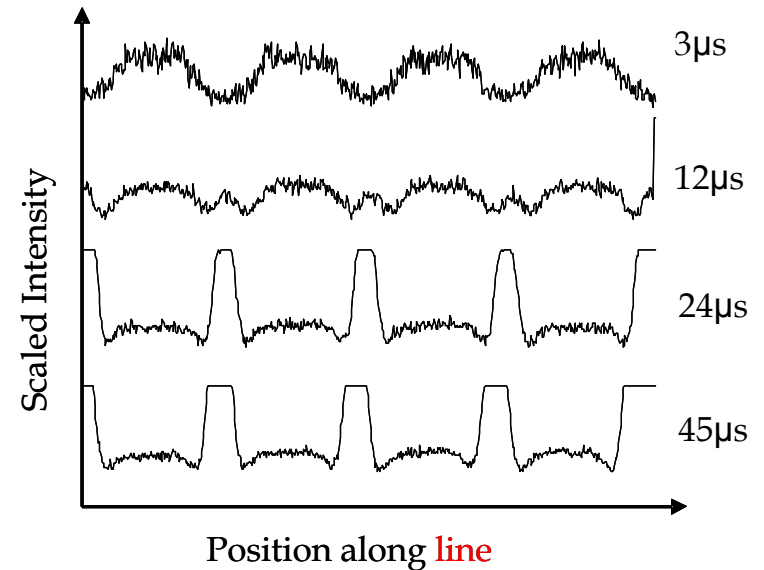


- **CD measurement of non-conductors is challenging because charging effects depend strongly on imaging conditions and cause:**
  - Image distortion
  - Contrast inversion



$E_0 = 600 \text{ eV}$   
 $I_0 = 44 \text{ pA}$

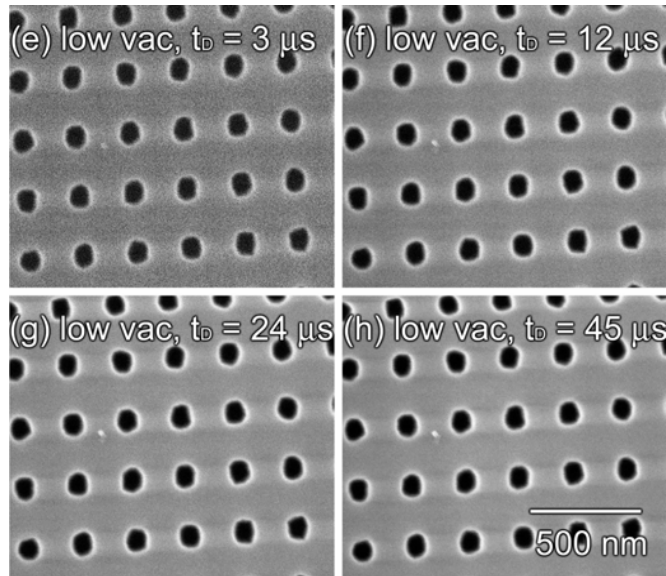
$\text{Si}_3\text{N}_4$  contact holes



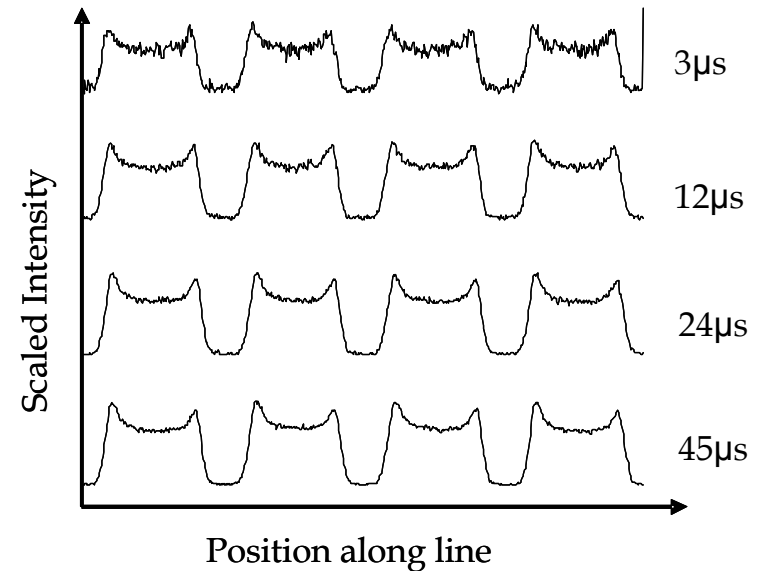




- **Good charge control independent of scan rate**
- **Good SNR at even lower electron fluence**
- **Contrast is preserved**



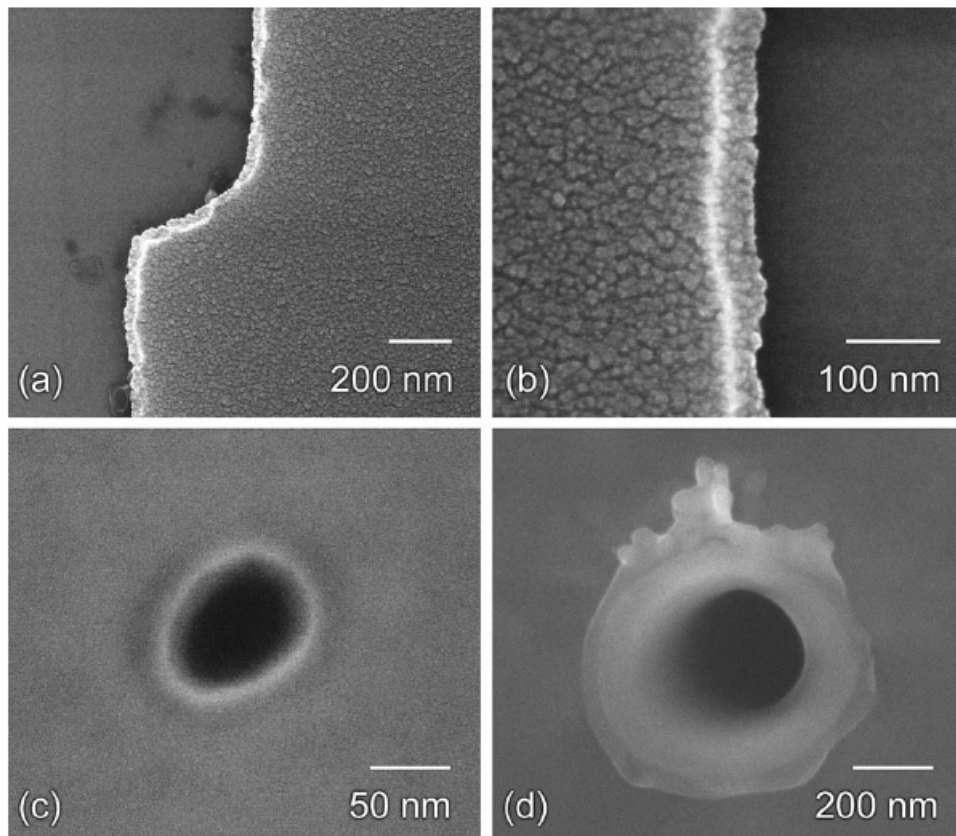
$E_0 = 5 \text{keV}$   
 $I_0 = 28 \text{pA}$



$\text{Si}_3\text{N}_4$  contact holes



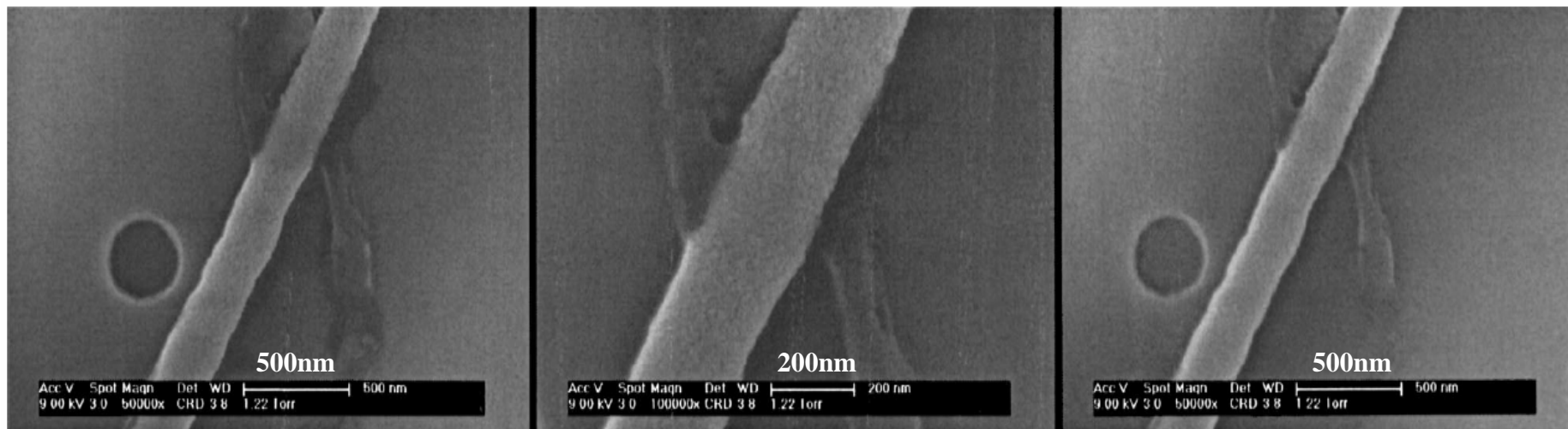
- **High resolution imaging with existing electron optics**



- (a) & (b) chrome-on-quartz  
photolithographic mask
- (c) contact hole in  $\text{Si}_3\text{N}_4$
- (d) electrostatic discharge pit in  
 $\text{SiO}_2$



- **Hydrocarbon contamination is reduced**



3min e-beam irradiation

(a) no contamination square

(b) shrinkage of the residue



ITRS challenges for both  $\leq$  &  $\geq$  22nm

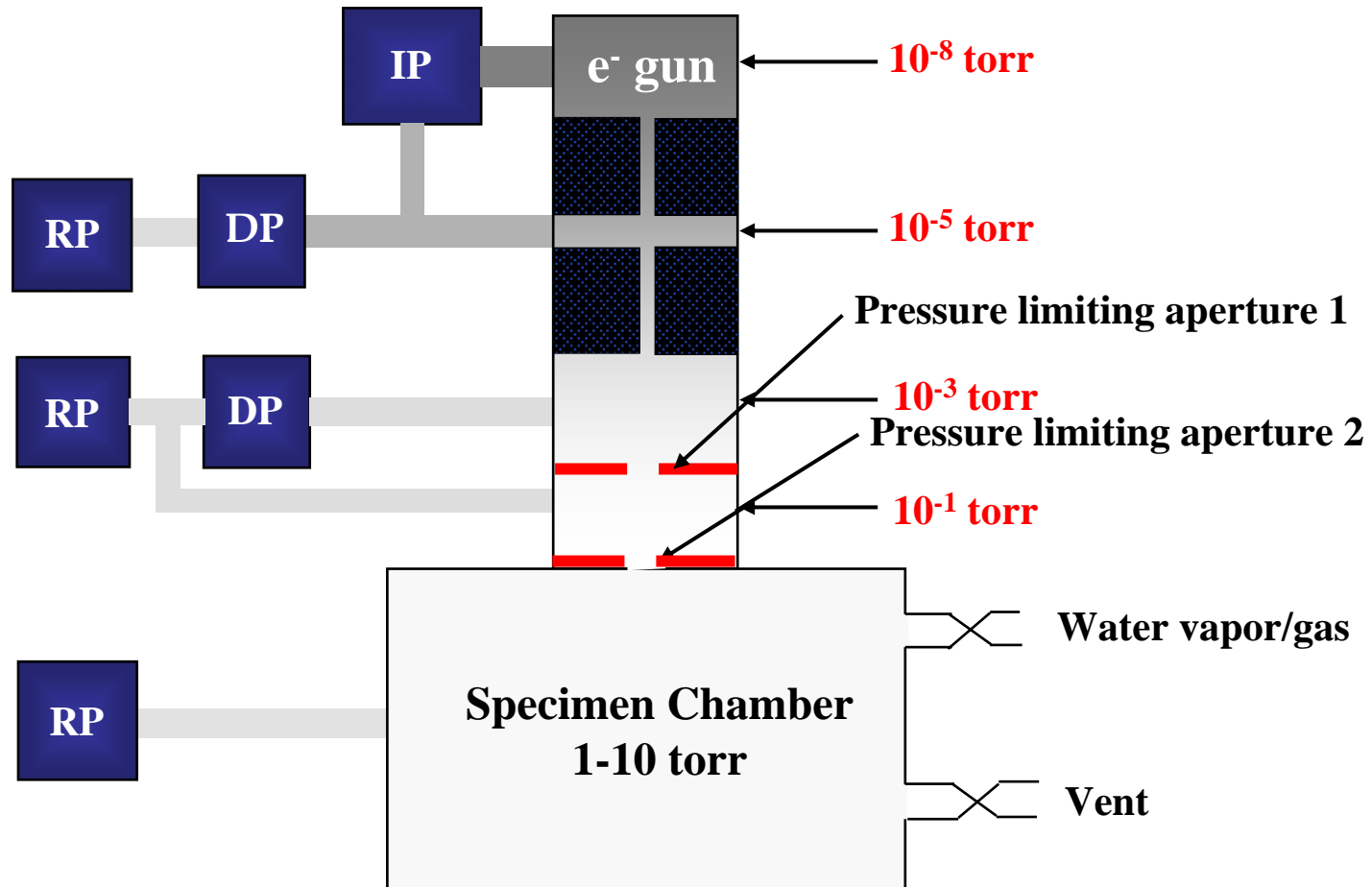
- ✓ Surface charging stabilization
- ✓ High resolution imaging
- ✓ Contamination control
- ✗ Method for interpreting the results



- **Low Vacuum SEM technology can address charging and contamination without sacrificing performance**
- **How are optimal conditions determined for alleviating charging artifacts?**
- **How does the gas affect the incident beam?**
- **How does the gas affect the secondary electron emission?**
- **What are the sources of signal, background, and noise collected by the detector?**



# **Low Vacuum SEM**





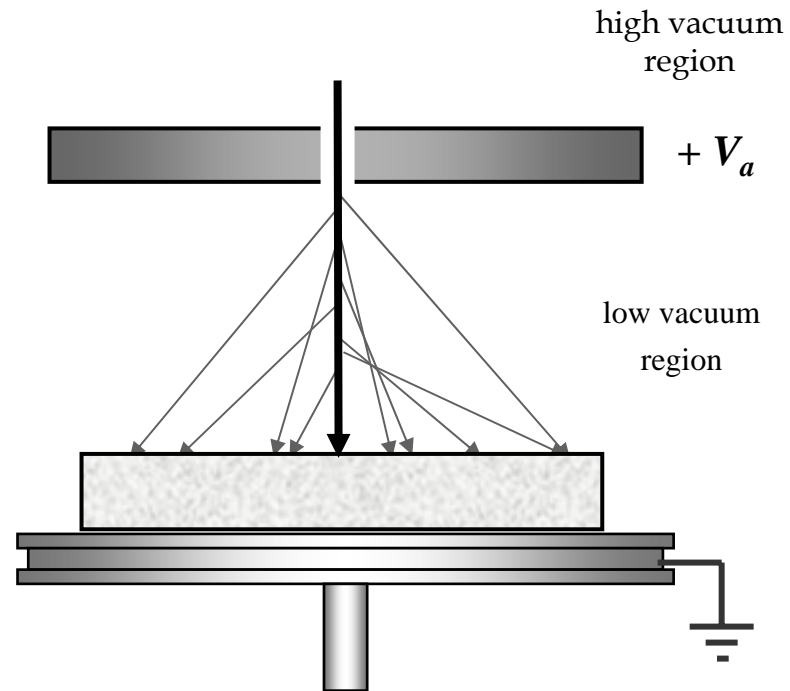
- Collisions of primary beam electrons with gas molecules result in a low current density “skirt” of scattered electrons surrounding the central beam

- It can extend over several hundred micrometers
- It does not degrade resolution
- It contributes to the background

- The unscattered fraction of the beam is given by:

$$f = \exp\left(-\frac{Q(E)Pl}{RT}\right)$$

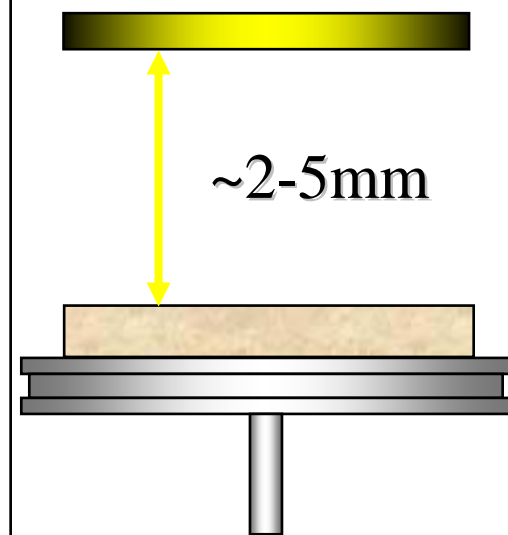
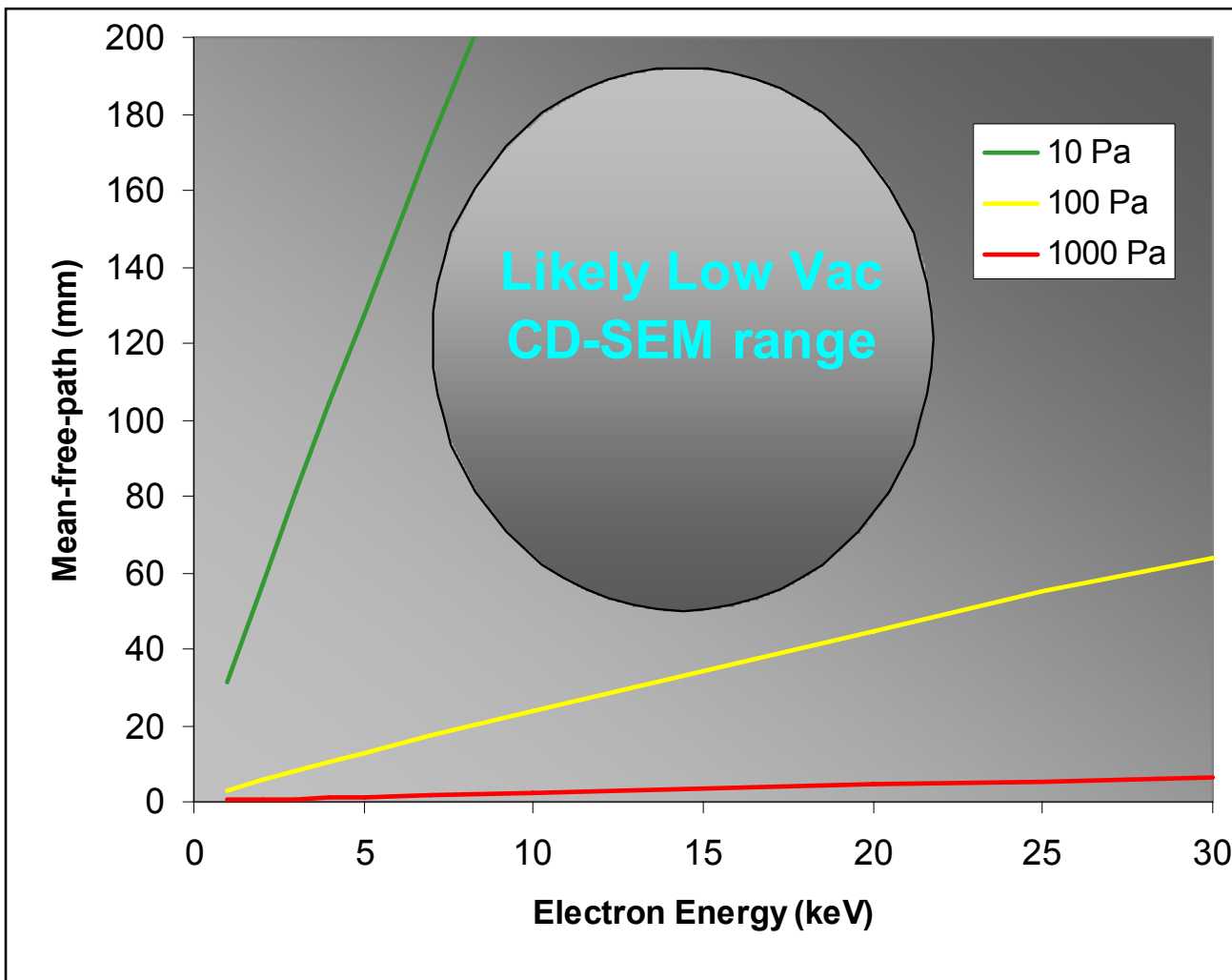
$Q(E)$ : total scattering cross-section,  $P$ : gas pressure,  
 $l$ : gas path length,  $R$ : gas constant,  $T$ : temperature







# Mean free path of Electrons in H<sub>2</sub>O Vapor





- The amplification of the emitted secondary electron (SE) signal takes place inside the chamber by ionizing collisions of the SEs with gas molecules
- A variety of electrodes and pole-pieces integrated into the final lens are used to create the electromagnetic field which drives the amplification cascade.

- Signal

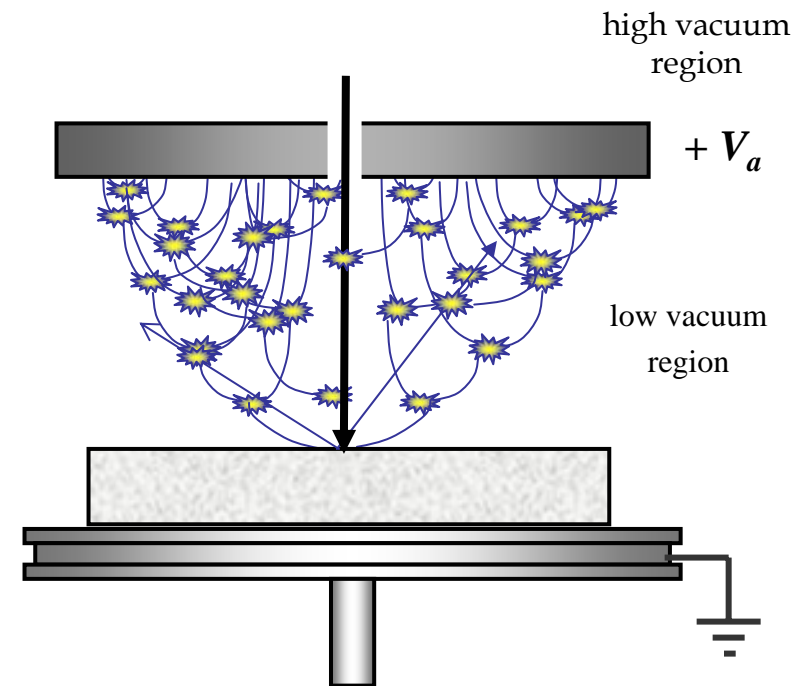
- Normally amplified SEs (gain  $\sim 10^3$ )

- Noise

- SEs amplification is a stochastic process

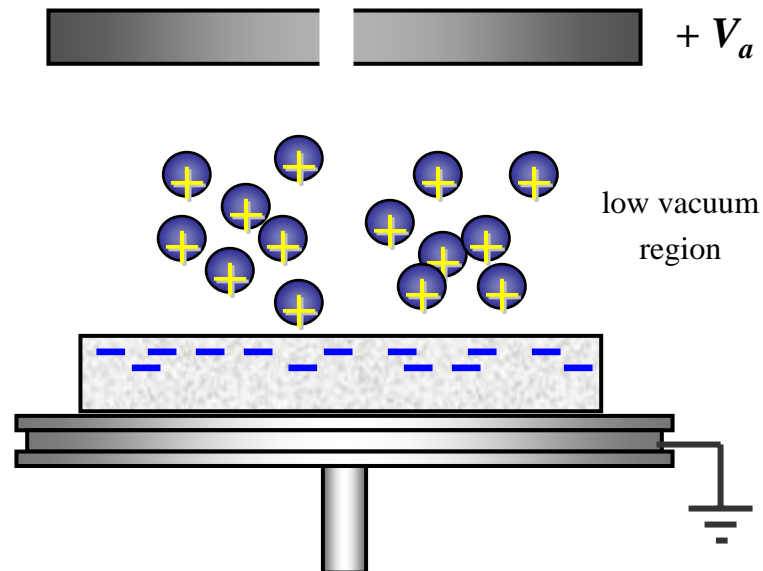
- Background

- Inevitable multiplication of BSEs and PEs





- Collisions of all emissions with gas molecules create positive gaseous ions
- Ions flow to regions of the specimen surface with a negative potential
- Ions recombine at surface, removing excess electronic charge





- **Operating parameter space is much larger to optimize**
  - **High Vacuum**
    - Beam energy
    - Beam current
  - **Low Vacuum**
    - Beam energy
    - Beam current
- **Charge control is a strong function of operating conditions**
  - **High Vacuum**
    - Working distance
  - **Low Vacuum**
    - Working distance
    - Gas-path-length
    - Cascade distance
    - Gas type
    - Gas pressure
- **Signal, background, and noise production are strong functions of operating conditions**
  - SE signal amplification via gas ionization cascade
  - Background due to elastic-scattering of the primary beam
  - Background due to inelastic high-energy electron scattering
  - Excess noise due to stochastic nature of cascade amplification



# **Cascade Amplification in a Constant Field**



- **Allows study of fundamental electron-gas interactions in a simple geometry**
  - Uniform electric field
  - Ionization efficiency largely independent of position in cascade

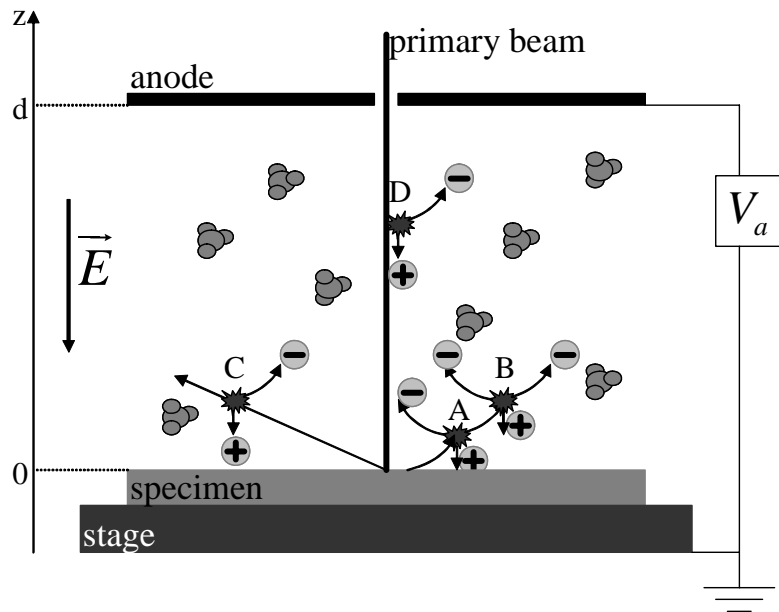
$$\alpha = AP \exp\left(-\frac{BPd}{V_a}\right)$$

$A, B$  : gas-specific parameters

$P$  : pressure

$d$  : cascade distance

$V_a$  : detector bias



A: ionizing events caused by secondary electrons (SEs) emitted from the specimen;

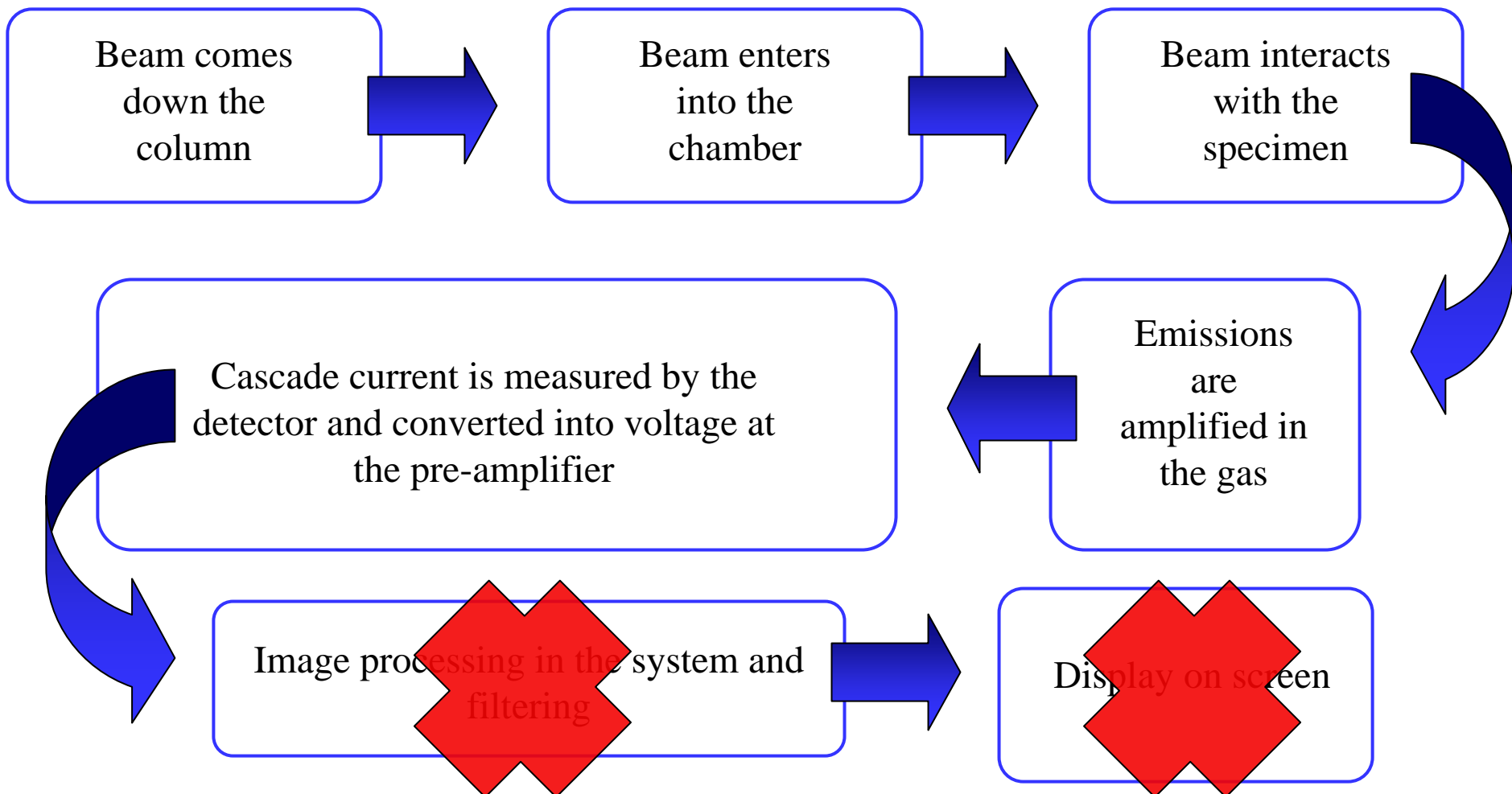
B: multiplication process;

C: ionizing events caused by backscattered electrons (BSEs);

D: corresponding effect caused by primary electrons (PEs).



# Flow Chart of the Signal Chain





- The secondary electron cascade current  $I_s$  arriving at the detector can be described by the following equation:

$$I_s = I_o \delta f \exp(\alpha d)$$

$I_o$  : primary electron beam current

$\delta$  : secondary electron yield

$\alpha$  : gas ionization efficiency

$d$  : cascade distance

$f$  : fraction of primary electrons that do not scatter into the skirt

- However, this needs to be modified in order to include the total measured signal in our experimental set-up
  - Addition of all SEs (included those generated by the skirt, BSEs, & PEs)

$$I = \left( \delta \exp(\alpha d) + \frac{\eta S_{BSE} P}{\alpha} \exp(\alpha d - 1) + \frac{S_{PE} P}{\alpha} \exp(\alpha d - 1) \right) I_o$$





- **Beam or shot noise:** Arising from Poisson distribution of emission of  $e^-$ .

$$N_{BE} = \sqrt{2eI_o b}$$

$e$  : electron charge

$I_o$ : primary beam current

$b$  : signal bandwidth

- **Gas cascade noise:** Arising from the formation of plasma in the chamber.

– Contribution to noise from:

- Secondary Electrons (SEs)
- Primary Electrons (PEs)
- Backscattered Electrons (BSEs)

- **Amplifier or Johnson's noise:** Arising from thermal and electronic noise in the detection system.

$$N_{AMP} = \sqrt{4kTRb}$$

$k$  : Boltzmann's constant

$T$  : temperature

$R$  : resistor





- The physical process of ionization of the gas in the presence of an electric field is similar to that of the production of  $e^- - h^+$  pairs in an avalanche photodiode
- The contribution to noise from the cascade is:

$$N = \sqrt{M^2 F_e}$$

$M$  : amplification factor (gain)

$F_e$  : excess noise factor (gain uncertainty)

- Application of the model for single carrier initiated/single carrier multiplication conditions (the lowest noise case for diodes) gives:

$$F_e = 2 - \frac{1}{M}$$

- Therefore the multiplication noise can be described by the equation:

$$N = \sqrt{M(2M - 1)}$$



- **Secondary electrons**

$$N_{SE} = \sqrt{2eI_o b(\delta(\delta + 1)) \exp(\alpha d) (2 \exp(\alpha d) - 1)}$$

- **Primary electrons**

$$N_{PE} = \sqrt{2eI_o b S_{PE} P l \left( \frac{\exp(\alpha d) - 1}{\alpha d} \right) \left( 2 \frac{\exp(\alpha d) - 1}{\alpha d} - 1 \right)}$$

- **Backscattered electrons**

$$N_{BSE} = \sqrt{2eI_o b \eta S_{BSE} P l \left( \frac{\exp(\alpha d) - 1}{\alpha d} \right) \left( 2 \frac{\exp(\alpha d) - 1}{\alpha d} - 1 \right)}$$

$I_o$  : primary beam current

$e$  : electron charge

$b$  : bandwidth

$\alpha$ : gas ionization efficiency

$\delta$ : SE yield

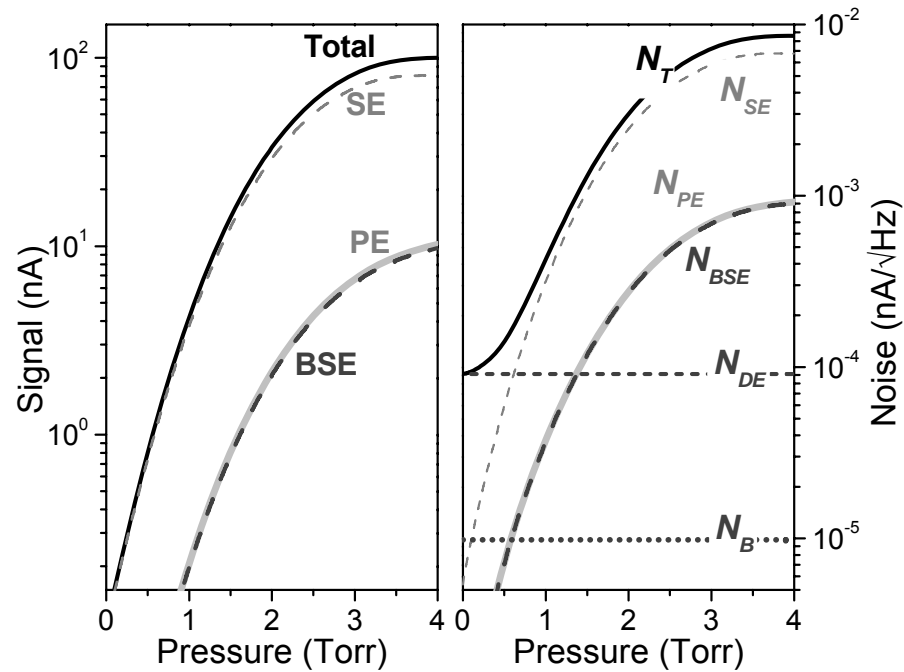
$S_{PE}$  &  $S_{BSE}$ : PEs & BSEs ionization efficiency

$P$  : pressure

$l$  : gas path length

$d$  : cascade distance

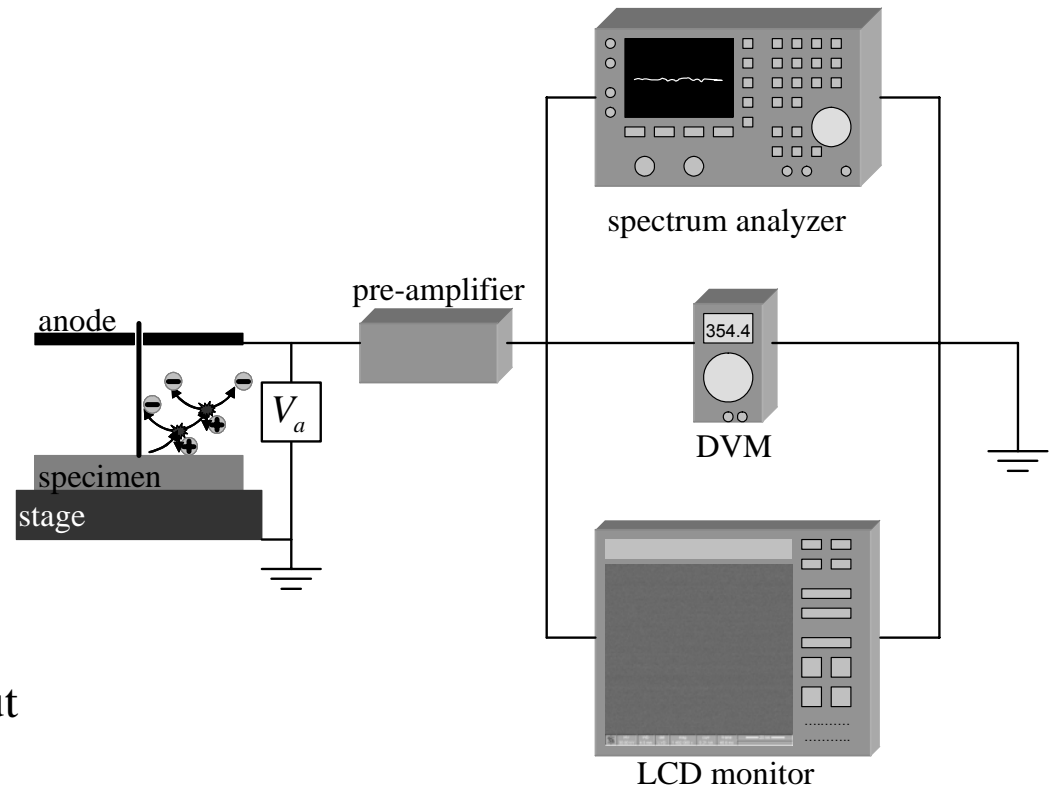
$\eta$  : BSE yield



- Total signal driven by the SE contribution
- Total noise dominated by:
  - Amplifier noise at low pressures
  - SE cascade component at high pressures

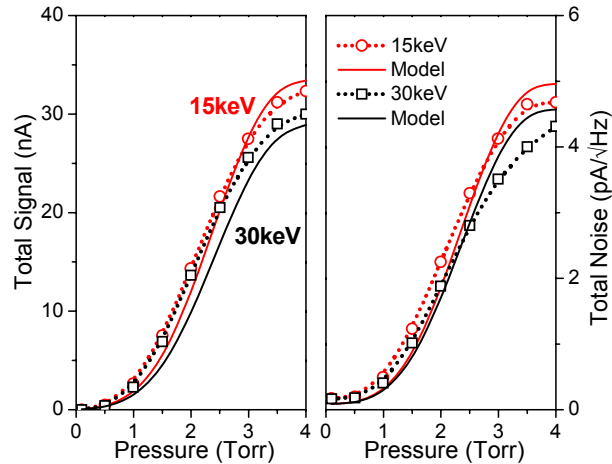


- **SEM data**
  - FEI Nova NanoSEM
- **Signal data**
  - Digital volt meter (DVM) connected to the output of the pre-amplifier
- **Noise data**
  - Agilent E4401B Spectrum Analyzer connected to the output of the pre-amplifier
    - Noise marker function

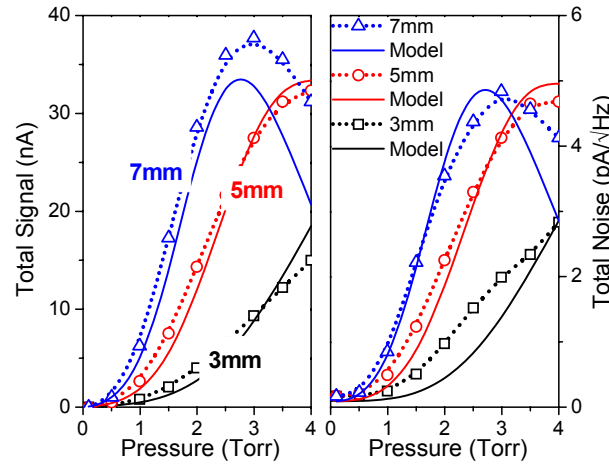




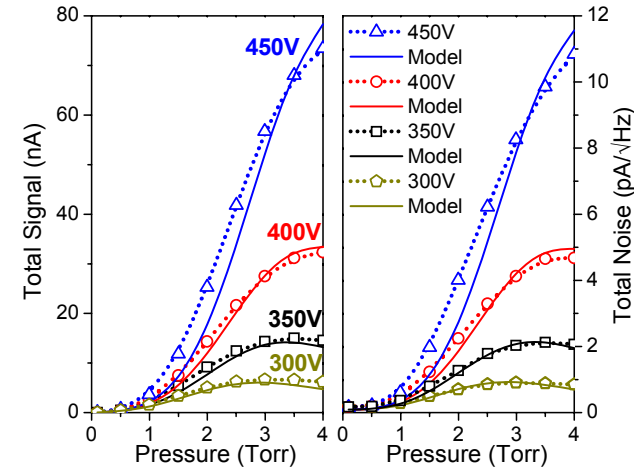
beam energy  $E_0$



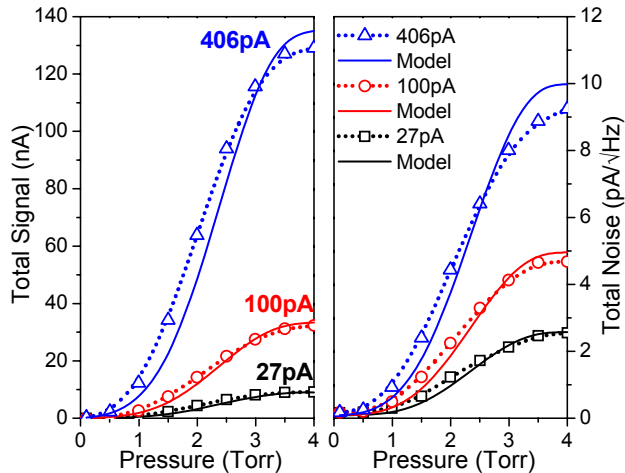
anode/specimen separation  $d$



detector bias  $V_a$



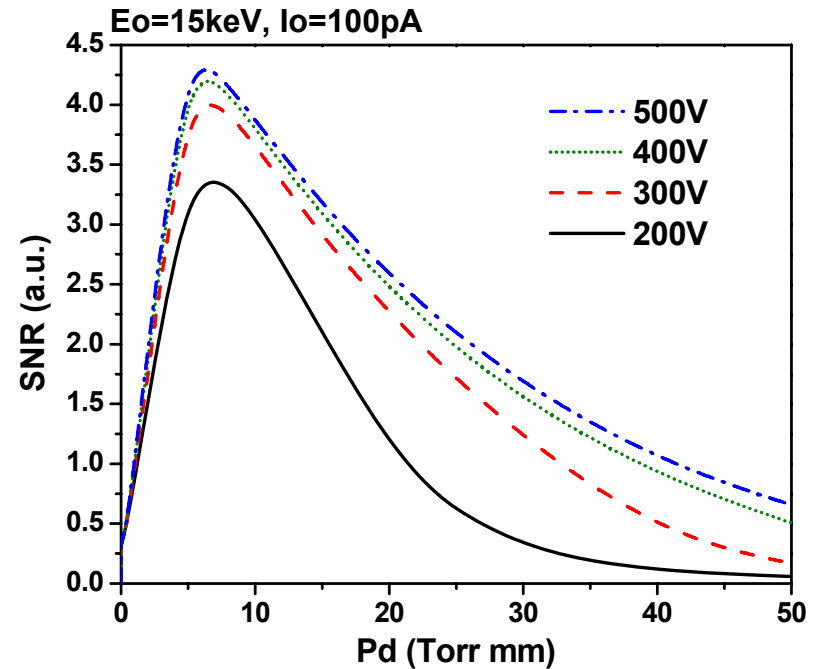
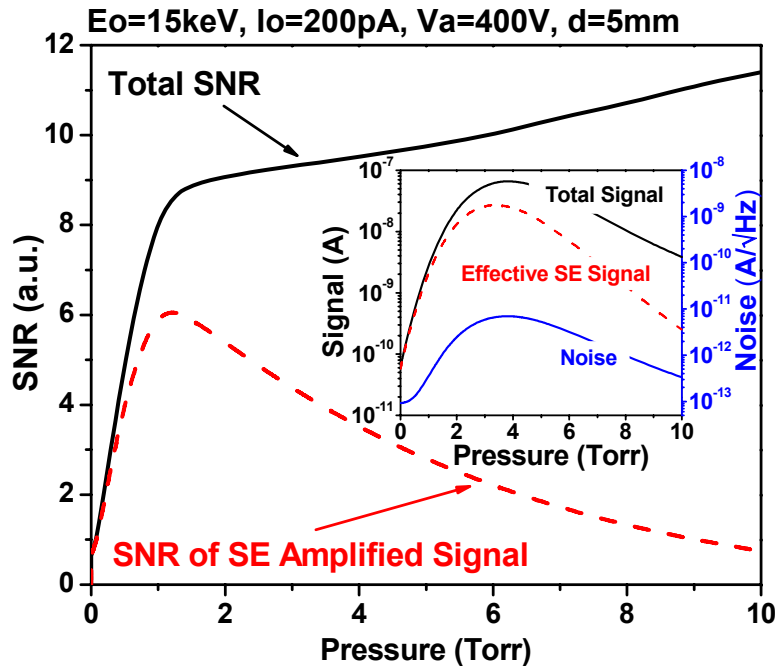
beam current  $I_0$



- Total signal increases as the primary beam energy is lowered
- The pressure at which maximum gain occurs is proportional to the electric field strength
- Maximum gain is an exponential function of the anode bias
- Total signal is proportional to the incident current



# Signal-to-Noise Ratio (SNR)

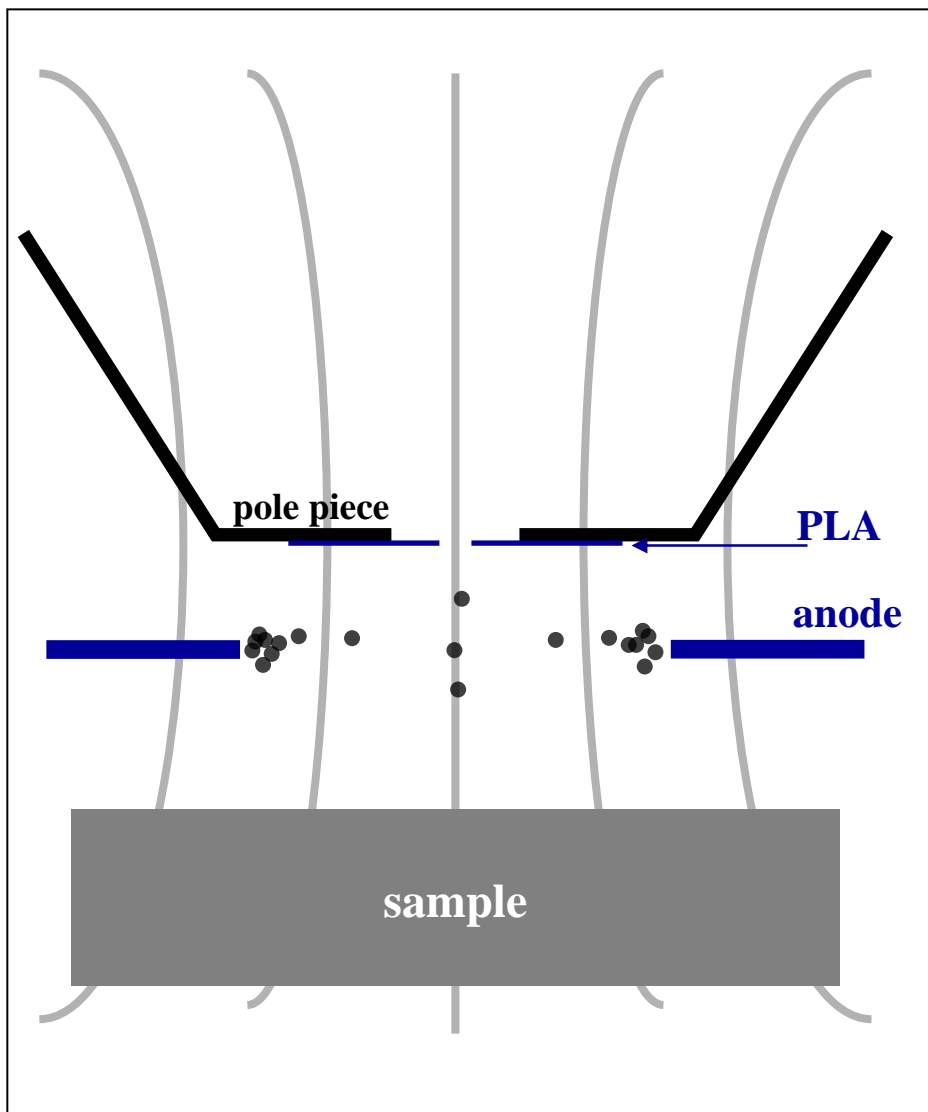


- Total cascade current reaching the detector contains a significant background component
- Optimization through the master curve approach

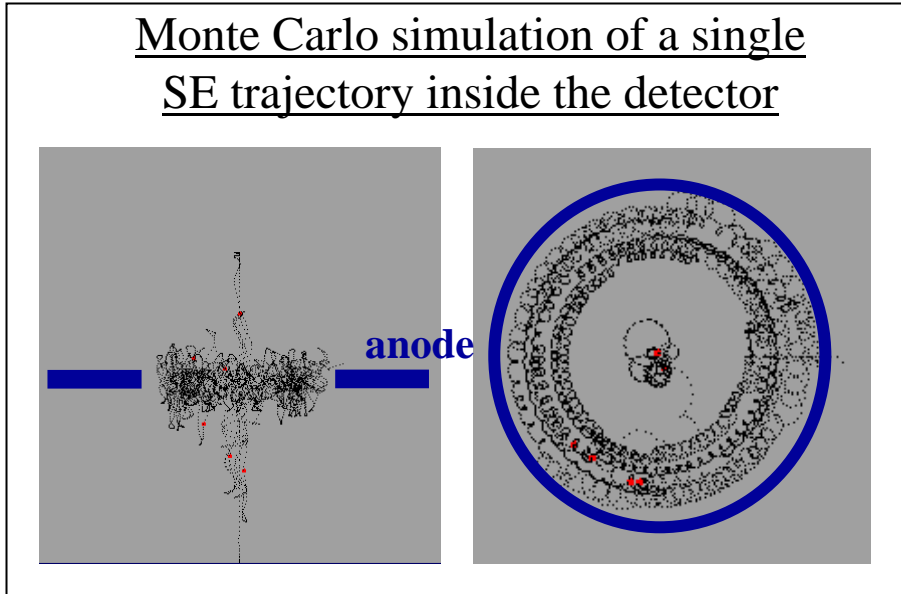


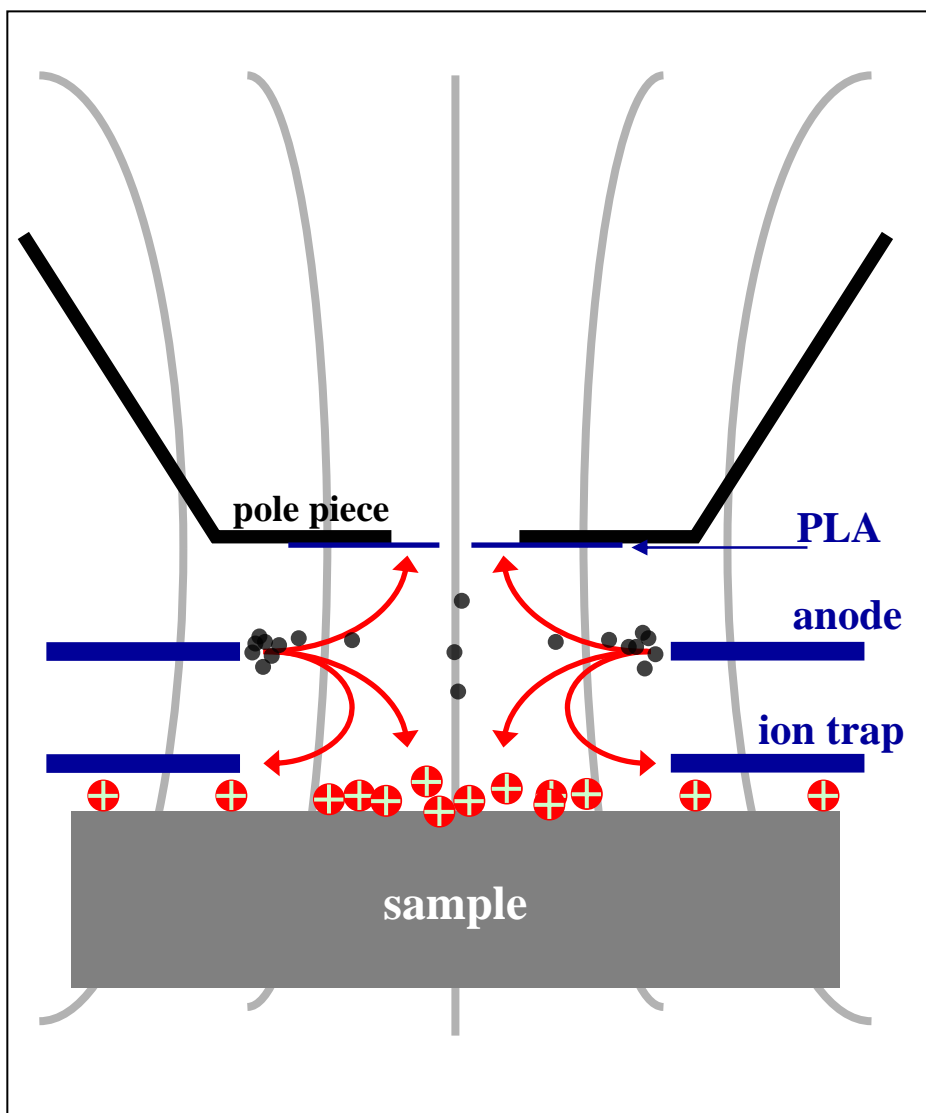


**Low Vacuum Immersion  
Lens Detector – Helix™**



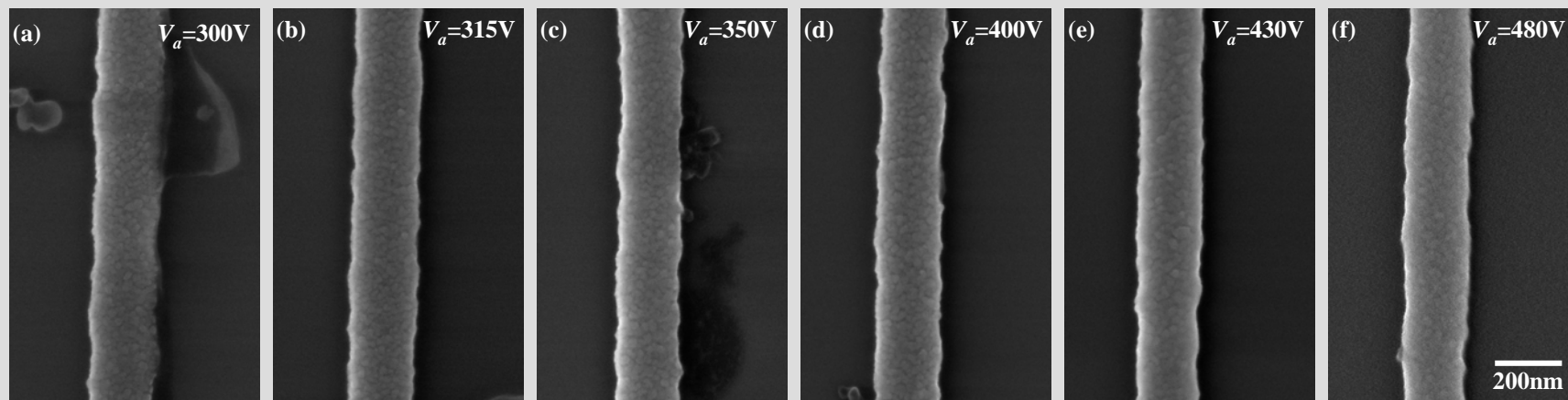
- SEs are attracted by an electric field, confined by a magnetic field
- Most amplification takes place along the anode periphery





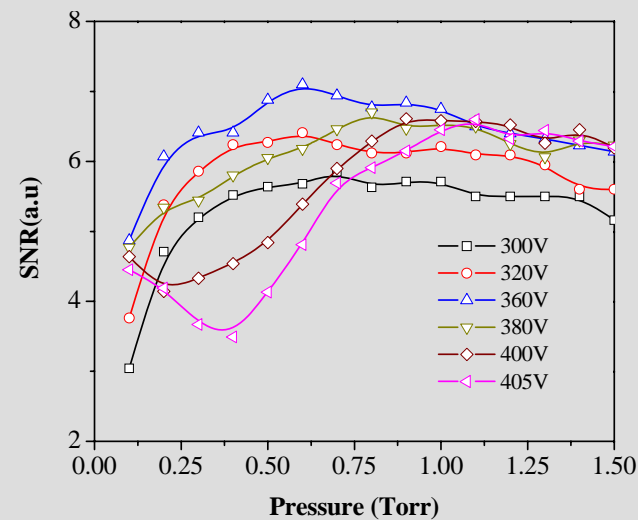
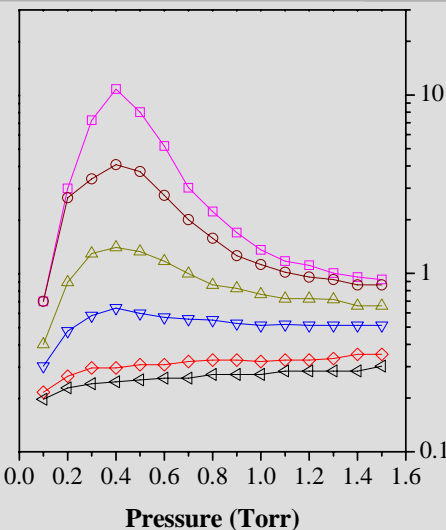
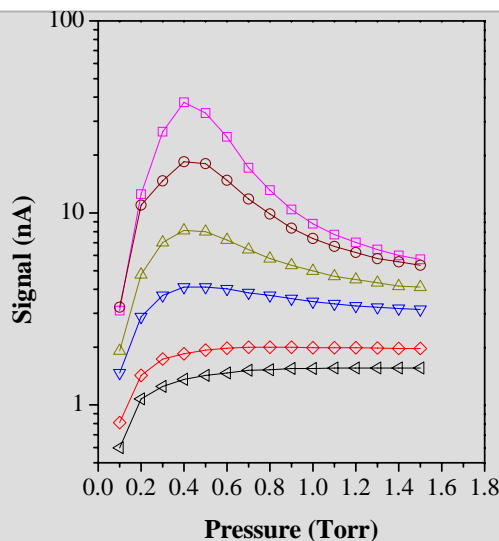
### Ions

- Generated along the anode and the beam axis
- $\sim 1/2$  drift up to the PLA, and  $1/2$  down to the sample
- Excess ions charge up the sample surface, reduce the detector field & give rise to imaging artifacts (gain  $10^3$ )
- Ion trap helps in self-regulation of ion current reaching the sample



Operating parameters:  $E_o=10kV$ ,  $WD=5mm$ ,  $I_o \sim 100pA$ ,  $B_f=200mT$

Operating parameters:  $E_o=5kV$ ,  $WD=4mm$ ,  $I_o \sim 100pA$ ,  $B_f=200mT$ , Pt



**Good charge control and SNR achievable but not intuitive**



- Low Vacuum SEM can be a viable approach for high resolution CD Metrology if appropriate modeling of the instrumental effects is developed
- Charge control and contamination reduction have been demonstrated
- Analytical signal and noise performance for the constant field cascade were modeled successfully



- Identification of the optimal gas composition for CD measurements
- Generalize analytical expressions for the gain and noise processes for complex detector configurations
- Develop analytical expressions for non-linear cascade processes
  - Breakdown (secondary ionization effects)
  - Scavenging (recombination of SE's with ions)
- Establish the limits of charge control under low vacuum conditions
- Evaluate the model for incorporation in the CD simulation code MONSEL for the realization of Low Vacuum Critical Dimension Metrology

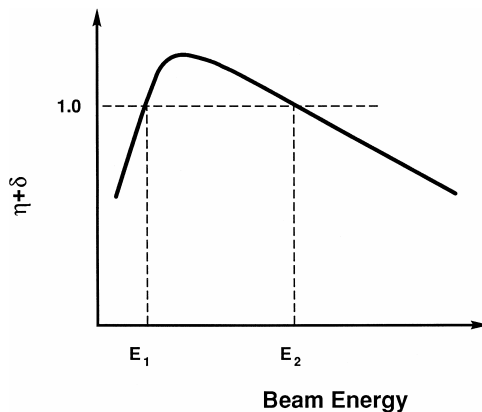


# Back-up slides



Useful SEM imaging of dielectrics, i.e. no charging artifacts, is possible by:

## Low voltage

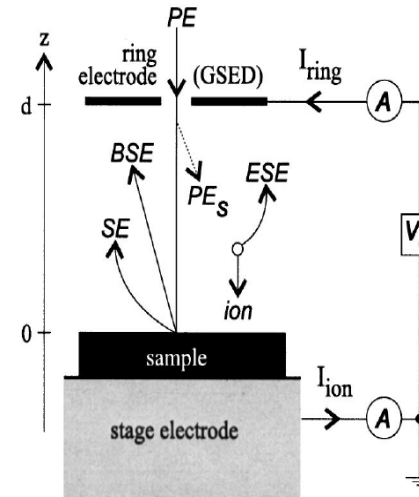


Charging is stabilized by fine-tuning of the electron beam energy to a critical value  $E_2$

But  $E_2$ :

- Is specimen dependent
- Can vary within the imaged area
- Is a function of electron fluence

## Low vacuum



Charging is stabilized by a weakly ionized gas inside the specimen chamber

Effective because it does not rely on:

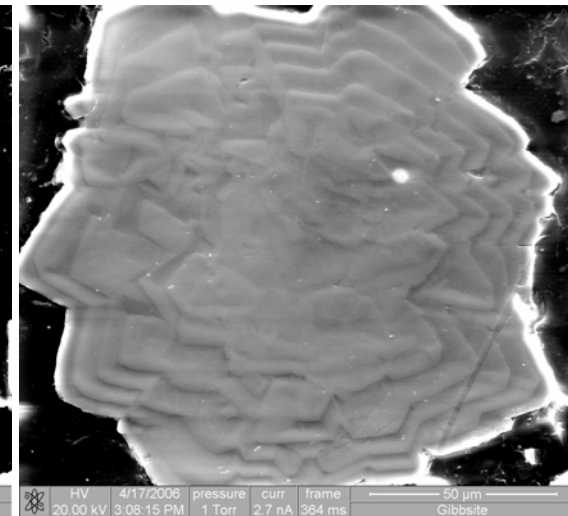
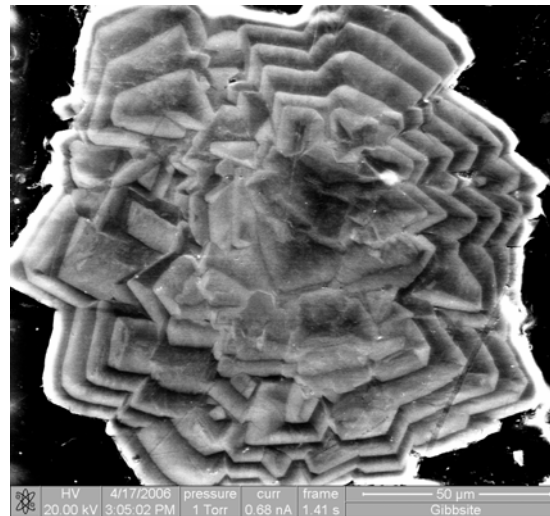
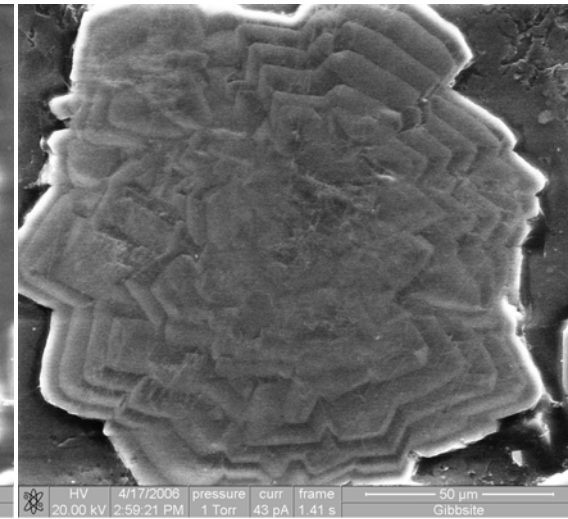
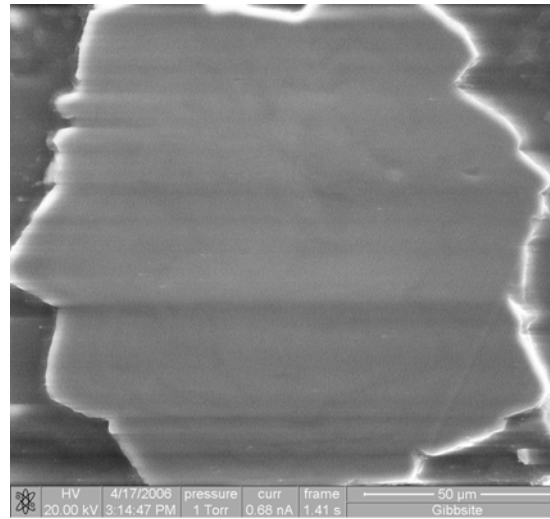
- Landing energy modulation
- SE recollection by the sample

Both suppress imaging artifacts but neither actually eliminates charging





- It stabilizes the surface potential
- It removes restriction on beam energy allowing it to be chosen on the bases of different criteria
- It eliminates image distortion
  - Image fidelity on metrology does not depend on imaging conditions
- It allows true SE imaging
  - Low voltage gives a signal complimentary to the backscattered
- It is potentially useful for acquiring information on buried structures
  - SE contrast reflects the local dielectric properties



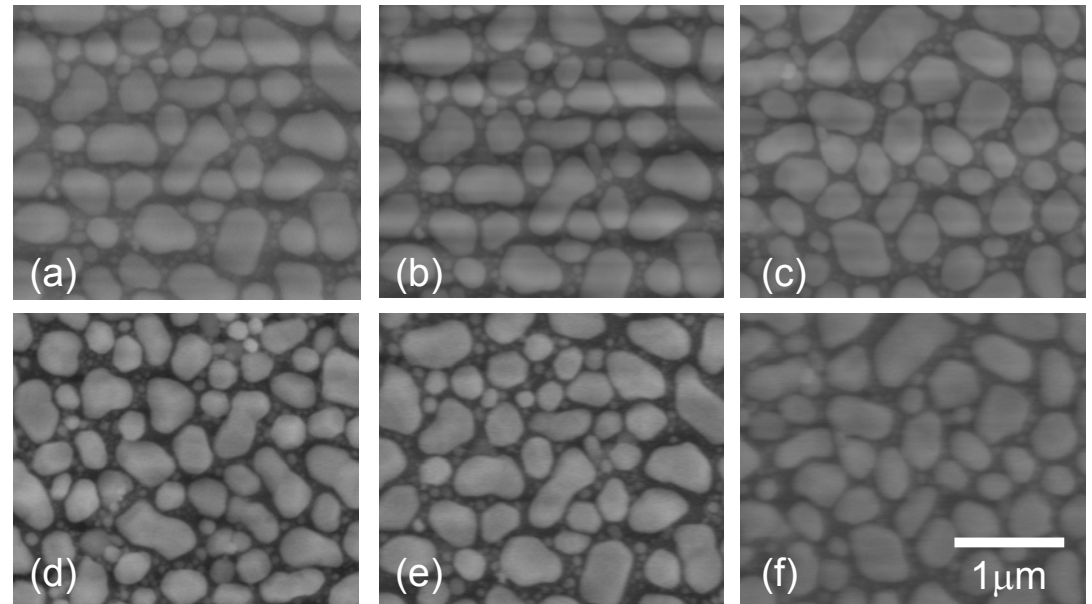
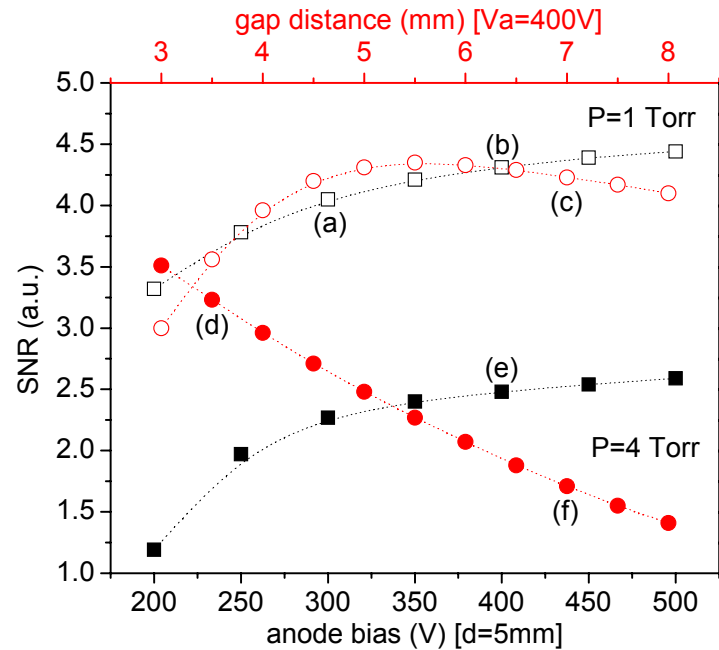


## Uncorrelated events add up in quadrature, correlated sum linearly

- All processes that take place in the cascade are correlated events
  - After ionization collisions, PEs and BSEs generate SE cascade
- Amplifier noise is independent of cascade processes

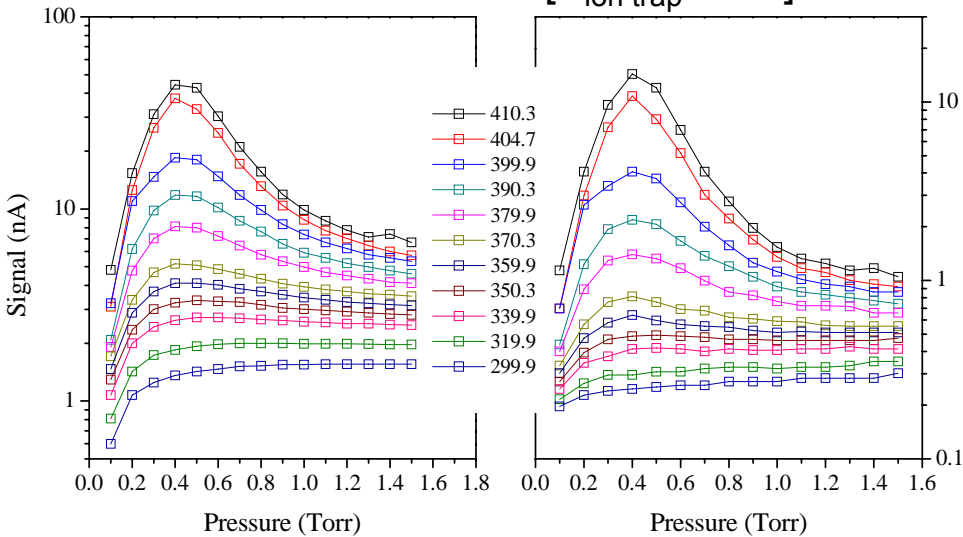
Therefore:

$$N_T = \sqrt{(N_{SE} + N_{BSE} + N_{PE})^2 + N_{AMP}^2}$$

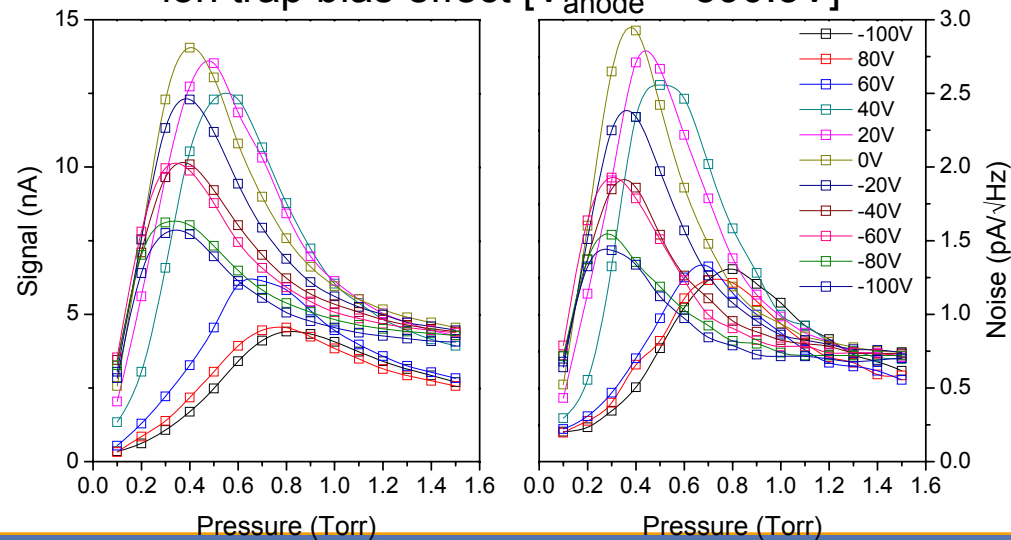


- Model predictions (plot) are in accordance with corresponding ESD images
- FFT image analysis is required to confirm the model predictions

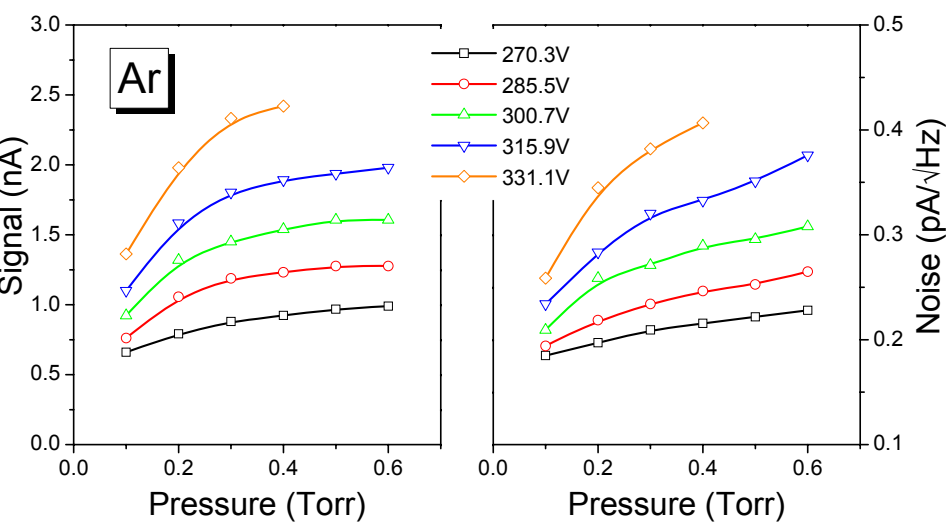
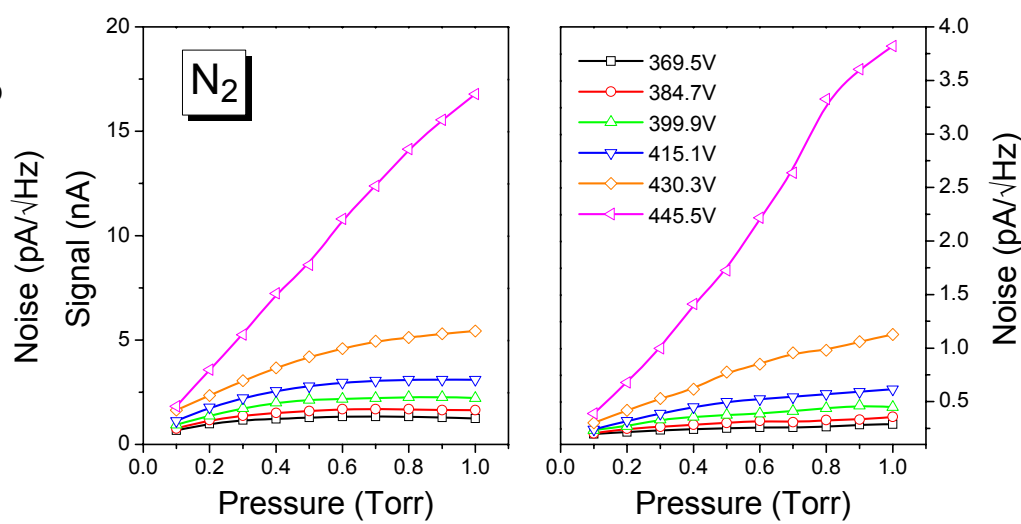
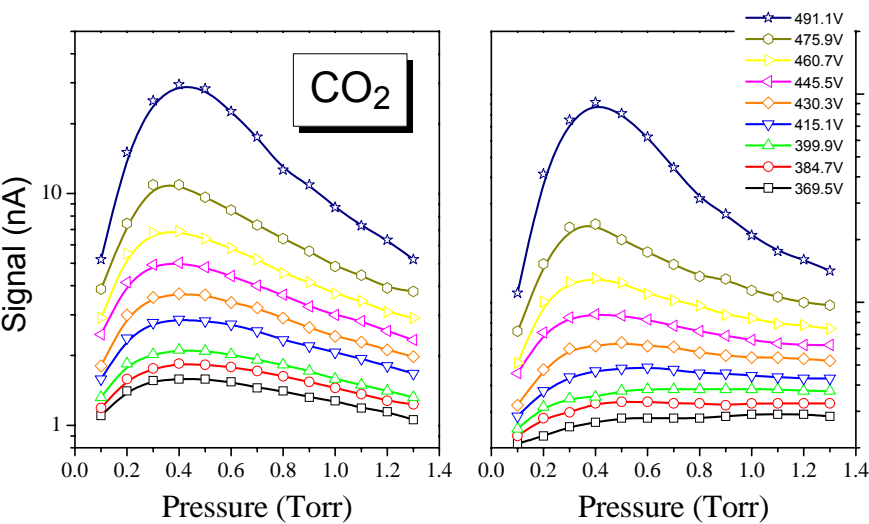
## anode bias effect [ $V_{\text{ion trap}} = 0V$ ]



## ion trap bias effect [ $V_{\text{anode}} = 390.3V$ ]



- Signal and Noise exhibit the same characteristics
- Low pressure peak (LPP) centered at ~0.4Torr
- High gains introduce instability in the system
- LPP changes position
  - Positive bias moves the peak at lower pressures
  - Negative bias moves the peak at higher pressures
- Grounded ion trap provides maximum gain (for the conditions investigated)



- Noise follows signal behavior for all gases
- CO<sub>2</sub> shows similar characteristics with H<sub>2</sub>O
- N<sub>2</sub> and Ar do not exhibit LPP
- Pressure ranges are different for each gas due to lower column vacuum breakdown